

# Theory of Wagner Ground Balance for Alternating-Current Bridges

By Richard K. Cook

A method is described for getting the Wagner ground balance of an alternating-current bridge when used with any three-terminal source. The basic idea is to insert impedances between the two ungrounded terminals of the source, and the corresponding terminals of the Wagner ground arm and the bridge, so as to "balance" approximately the currents from the source. The method is applicable to high-voltage type Schering bridges. Such a bridge, incorporating the principles of the new method for securing the Wagner ground balance, has been designed and built principally for measurement of the small capacitances, about 50  $\mu\mu\text{f}$ , of some types of condenser microphones.

## I. Introduction

A Wagner ground is used sometimes with alternating current bridges in order to eliminate the effects of capacity to ground. The same end can be achieved by a system of electrostatic shielding, but the required shielding can become very complicated. With a Wagner ground, it is possible to secure high precision with simple shielding.

The object of this paper is to show how almost any available three-terminal source can be used with a Wagner-grounded bridge. The simple theory that is usually given assumes that the source has only two accessible output terminals, which are connected to the two ungrounded input terminals of the bridge [1]<sup>1</sup>. The Wagner ground balance is readily found theoretically, but sometimes it is difficult or impossible to find experimentally, particularly in the case of high-voltage type Schering bridges used to measure small capacitances.

Ogawa [2] has pointed out that the experimental difficulty is due to the fact that available sources really have three output terminals, and the simple theory of a two-terminal source does not apply. He has given a comprehensive treatment of the general theory of Wagner-grounded alternating-current bridges and has solved the difficulty of securing balance by means of a specially con-

structed transformer satisfying an electrical condition which simplifies the Wagner ground balance equation. The theory has been applied by Astin [3] to a Schering bridge used for precision measurement of the power factors of air capacitors.

The theory given below shows how almost any available three-terminal source can be used, thus obviating the need for a special transformer. The theory is applied to the design of a high-voltage type Schering bridge for measurement of the small capacitances of condenser microphones at frequencies between 500 and 10,000 c/s. Such capacitance measurements are required when making absolute calibrations of condenser microphones by the reciprocity technique.

It should be pointed out that there are other ways of securing Wagner ground balance. For example, in a high-voltage type Schering bridge for capacitance measurements at power-supply frequencies, which has recently become commercially available [4], the Wagner ground balance is obtained by adjustment of the amplitude and phase (relative to ground) of the voltage applied to one of the input terminals of the bridge.

## II. Theory of a Wagner-grounded Bridge

The basic idea of the new way for getting Wagner ground balance is as follows: In the circuit shown in figure 1, suppose the bridge to be balanced, so that points  $B_1$  and  $B_2$  are at the

<sup>1</sup> Figures in brackets indicate the literature references at the end of this paper.

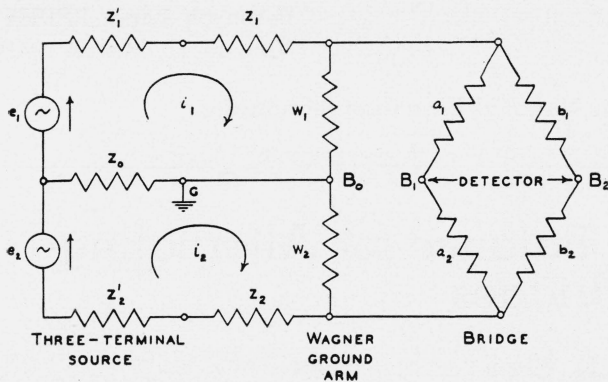


FIGURE 1.—Circuit diagram of Wagner-grounded alternating-current bridge.

same potential. Hence  $a_1 b_2 = a_2 b_1$ . Suppose also that  $i_1 = i_2$ ; the wire  $B_0 G$  will then carry no current. This will be referred to as the balanced ground-current condition. It is clear that this assumption eliminates the effect of  $z_0$ , so that the simple bridge theory will apply. The third assumption is that  $a_1 w_2 = a_2 w_1$ . Then it follows from these three assumptions that the Wagner ground is balanced, or that points  $B_0$  and  $B_1$  are at the same potential. If  $e_1 = e_2$ , the balanced ground-current condition can be realized by adjusting  $z_1$  and  $z_2$  so that  $z_1 + z_1'$  plus the parallel combination of  $a_1$ ,  $b_1$ , and  $w_1$ , is equal to  $z_2 + z_2'$  plus the parallel combination of  $a_2$ ,  $b_2$ , and  $w_2$ .

The mathematics needed for proof of the foregoing statements can be obtained from Ogawa's article, but a simplified derivation might be of interest. We assume a more general condition in which the ground-current might not be balanced, so that  $i_1$  is not necessarily equal to  $i_2$ , but in which the potentials at  $B_0$ ,  $B_1$ , and  $B_2$  are the same (Wagner ground and bridge balanced). Also, we assume  $e_1$  is not necessarily equal to  $e_2$ .

Before proceeding, a few words about three-terminal sources are in order. Ogawa has shown by a generalization of Thévenin's theorem that any three-terminal source can be represented by the network shown in figure 1. Subsequently a more general proof was given by Starr [5]. An equivalent network of this type applies whenever the potentials of the terminals are linear functions of the three currents through the terminals. Whether or not the network can be applied to a source that is available in the laboratory is a question that can, of course, be answered only by

experiment. The simple bridge theory which is usually given tacitly assumes that  $z_0 = \infty$ .

By virtue of the bridge balance, the circuit of figure 1 can be replaced by the simpler circuit of figure 2, in which the impedance  $s_1$  is the parallel combination of  $a_1$  and  $b_1$ , and  $s_2$  is the parallel combination of  $a_2$  and  $b_2$ . Also,  $Z_1 = z_1 + z_1'$ ,  $Z_2 = z_2 + z_2'$ . As in figure 1, the detector impedance is not indicated, as it has no current at balance. Application of Maxwell's loop equations to the network yields the following equations:

$$\left. \begin{aligned} (w_1 + Z_1 + z_0)i_1 - z_0 i_2 - w_1 i_3 - e_1 &= 0 \\ -z_0 i_1 + (w_2 + Z_2 + z_0)i_2 - w_2 i_3 - e_2 &= 0 \\ w_1 i_1 + 0 - (w_1 + s_1)i_3 + 0 &= 0 \\ 0 + w_2 i_2 - (w_2 + s_2)i_3 + 0 &= 0. \end{aligned} \right\} \quad (1)$$

As these equations are linear and homogeneous in  $(i_1, i_2, i_3, 1)$ , the necessary and sufficient condition for a nonzero solution for  $(i_1, i_2, i_3, 1)$  is the vanishing of the determinant of the coefficients:

$$\begin{vmatrix} w_1 + Z_1 + z_0 & -z_0 & -w_1 & -e_1 \\ -z_0 & w_2 + Z_2 + z_0 & -w_2 & -e_2 \\ w_1 & 0 & -(w_1 + s_1) & 0 \\ 0 & w_2 & -(w_2 + s_2) & 0 \end{vmatrix} = 0. \quad (2)$$

Expansion of the determinant yields the fundamental balance equation,

$$\begin{aligned} e_1 \left[ s_2 + Z_2 + \frac{s_2 Z_2}{w_2} + \left( \frac{s_2}{w_2} - \frac{s_1}{w_1} \right) z_0 \right] = \\ e_2 \left[ s_1 + Z_1 + \frac{s_1 Z_1}{w_1} - \left( \frac{s_2}{w_2} - \frac{s_1}{w_1} \right) z_0 \right], \end{aligned} \quad (3)$$

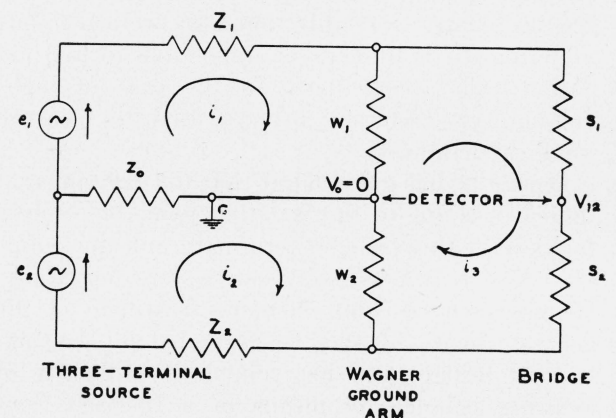


FIGURE 2.—Simplified circuit diagram of Wagner-grounded bridge.

which is basically the same as Ogawa's eq 62.

If  $i_1 = i_2$ , then eq 1 show it is necessary that

$$\frac{s_2}{w_2} = \frac{s_1}{w_1} \quad (4)$$

If in addition,  $e_1 = e_2$ , then eq 3 reduces to

$$Z_2 + \frac{s_2 w_2}{s_2 + w_2} = Z_1 + \frac{s_1 w_1}{s_1 + w_1} \quad (5)$$

This is the algebraic form of the statement made in the first paragraph of this section.

Suppose that  $z_0 \rightarrow \infty$ . Then it is found that eq 3 reduces to eq 4, provided  $e_1 + e_2 \neq 0$ . This is the usual Wagner ground balance condition obtained by application of the simple theory.

We define the unbalance voltage to be the amount by which  $V_{12}$  exceeds ground potential when the Wagner ground is not balanced. This quantity is needed for studying the convergence of the sequence of balancing operations. It is justifiable to ignore the impedance of the amplifier-detector system which might be employed to observe  $V_{12}$ . The reason for this is that the impedance will affect only the magnitude and phase of  $V_{12}$ , and will not affect the convergence of the sequence of balancing operations when the unbalance current is small in comparison with the main bridge currents. We find

$$V_{12} = \frac{1}{\Delta} \left\{ e_1 \left[ s_2 + Z_2 + \frac{s_2 Z_2}{w_2} + \left( \frac{s_2}{w_2} - \frac{s_1}{w_1} \right) z_0 \right] - e_2 \left[ s_1 + Z_1 + \frac{s_1 Z_1}{w_1} + \left( \frac{s_1}{w_1} - \frac{s_2}{w_2} \right) z_0 \right] \right\} \quad (6)$$

where

$$\Delta = \left( Z_1 + s_1 + \frac{Z_1 s_1}{w_1} \right) \left( 1 + \frac{Z_2}{w_2} + \frac{z_0}{w_2} \right) + \left( Z_2 + s_2 + \frac{Z_2 s_2}{w_2} \right) \left( 1 + \frac{Z_1}{w_1} + \frac{z_0}{w_1} \right) - z_0 \left( \frac{Z_1}{w_1} + \frac{s_2}{w_2} + \frac{Z_1 s_2}{w_1 w_2} + \frac{Z_2}{w_2} + \frac{s_1}{w_1} + \frac{Z_2 s_1}{w_2 w_1} \right)$$

The way in which eq 3 was derived shows it to be a necessary condition for Wagner ground balance. Equation 6 shows that eq 3 is also sufficient, if  $\Delta \neq 0$ .

### Theory of Wagner Ground Balance

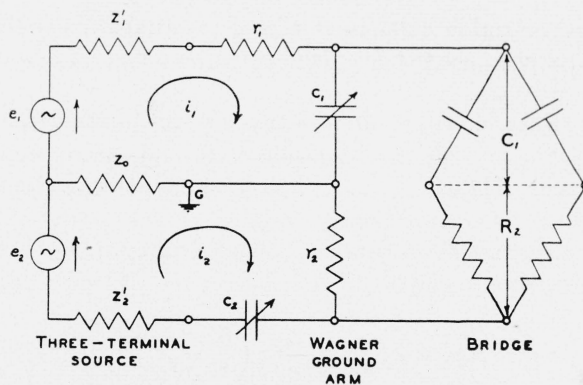


FIGURE 3.—High-voltage type Schering bridge with Wagner ground.

### III. Application to a Schering Bridge

The application of the Wagner-grounding technique described above to a high-voltage type Schering bridge, which is useful for measuring the capacitance of condenser microphones, is easily accomplished. A circuit diagram for such a bridge is shown in figure 3. Here  $C_1$  is the parallel combination of the two capacitances in the high-voltage arms, and  $R_2$  is the parallel combination of the two resistances in the low-voltage arms.

It is important to study how the Wagner ground balance is obtained when the circuit parameters are varied. A practical advantage of the Schering bridge is that balance can be obtained by variation of capacitance only. The Wagner ground balance of the Schering bridge circuit shown in figure 3 can also be obtained by variation of capacitance only, as the following analysis shows.

As before, we shall suppose that both the bridge and the Wagner ground are balanced. Also, we make the following simplifying assumptions, which can be readily fulfilled experimentally. (1) The three-terminal source is the secondary of a center-tapped transformer. (2) The transformer is symmetrical enough so that  $e_2 = (1 + \alpha + j\beta)e_1$ , where  $\alpha \ll 1$  and  $\beta \ll 1$ . (3) The terminals of the transformer secondary are externally connected to the center tap by equal low resistances, so that

$$z_1' \approx z_2', \text{ and } z_0 \approx \frac{-z_1'}{2}. \quad z_1' \text{ and } z_2' \text{ are positive}$$

resistances, but  $z_0'$  is a negative resistance. In what follows there will be no loss of generality in supposing that  $z_1'$  is included in  $r_1$ , and  $z_2'$  is included in the residual of  $e_2$ . (4)  $e_1$ ,  $e_2$ , and  $C_1$

are variable capacitors whose residuals are small, and  $r_1$ ,  $r_2$ , and  $R_2$  are fixed resistors whose residuals are small.

First we shall suppose that all the small quantities  $\alpha$ ,  $\beta$ , and the residuals of the impedances are zero. Hence  $e_1 = e_2$ . The object of this is to find the approximate values of the adjustable capacitances  $c_1$  and  $c_2$ . Application of eq 3 to the circuit yields

$$\left. \begin{aligned} R_2 &= r_1 + \left(\frac{c_1}{C_1}\right)r_1 - 2z_0\left(\frac{R_2 - c_1}{r_2 - C_1}\right); \\ \frac{1}{\omega c_2}\left(1 + \frac{R_2}{r_2}\right) &= \frac{1}{\omega C_1}. \end{aligned} \right\} \quad (7)$$

The first of these corresponds to the real part of eq 3, and the second corresponds to the imaginary part. Solutions of eq 7 for  $c_1$  and  $c_2$  are as follows;

$$\left. \begin{aligned} c_1 &= \left(\frac{R_2 - r_1 + \frac{2z_0 R_2}{r_2}}{r_1 + 2z_0}\right)C_1; \\ c_2 &= \left(\frac{R_2}{r_2} + 1\right)C_1. \end{aligned} \right\} \quad (8)$$

These equations indicate that if

$$r_1 + 2z_0 > 0, \text{ and } R_2 - r_1 + \frac{2z_0 R_2}{r_2} > 0, \quad (9)$$

the Wagner ground balance can be obtained by variation of capacitance only. The balance is independent of frequency.

It is to be noted that  $c_1$  can be adjusted to satisfy the real part of eq 3, and  $c_2$  can be separately adjusted to satisfy the imaginary part. This, together with eq 6, shows that the Wagner ground unbalance voltage, which is a vector in the complex plane, can be reduced to zero by a series of steps, successive steps being mutually perpendicular. This is a highly desirable property of a sequence of balancing operations for an alternating-current bridge.

The unbalance voltage can be reduced to zero by adjustment of other pairs of parameters besides  $c_1$  and  $c_2$ . Equations 7 show that variation of  $r_1$  and  $c_2$  will also yield successive voltage steps that are mutually perpendicular. However, all other pairs will yield successive steps, which in general are not mutually perpendicular, an example of which is the combination  $r_2$  and  $c_1$ .

The pair  $r_1$ ,  $c_1$  cannot be used at all to get balance, as the second eq 7 contains neither of these parameters.

A more detailed analysis in which  $\alpha$ ,  $\beta$ , and the residuals are not assumed to be zero shows that the Wagner ground balance can still be secured by variation of the capacitances  $c_1$  and  $c_2$  only. However, the successive steps in the unbalance voltage, obtained by varying  $c_1$  and  $c_2$  successively, are vectors which are not quite perpendicular to one another.

It might be inconvenient to select the circuit parameters so that the inequalities (9) hold if a center-tapped transformer is used as a source, as for such a transformer,  $z_0$  is a negative resistance. An examination of the fundamental balance eq 3 shows that the effect of  $z_0$  disappears if the condition given by eq 4 holds. When this condition is applied to the Schering bridge circuit discussed above, we find

$$c_1 = \frac{R_2}{r_2}C_1, \quad r_1 = \frac{r_2 R_2}{r_2 + R_2}, \quad c_2 = c_1 + C_1. \quad (10)$$

It is understood that  $r_1$  includes the resistive impedance  $z_1'$ . The ground current is balanced in this case.

Another possibility is to introduce a positive external resistor, whose magnitude is  $|z_0|$ , at the center tap of the transformer. This would, in effect, reduce  $z_0$  to zero, and the terms in eq 3 containing  $z_0$  would vanish.

Still another possibility is to use two similar transformers driven by the same oscillator. The transformers would have to be arranged so that there is no mutual induction between them. The use of buffer amplifiers might also be required between the oscillator and transformers. The whole arrangement can be converted into a three-terminal source having  $z_0 = 0$  by connecting a terminal on the secondary of one transformer to a terminal on the other.

#### IV. Description of a Bridge

A Wagner-grounded Schering bridge, based on the principles outlined above, has been constructed in our Sound Laboratory. The bridge is the high-voltage type and has been used mainly to measure condenser-microphone capacitances, which are about 50  $\mu\mu\text{f}$ , at frequencies ranging from 500 to 10,000 c/s. Accuracy of the order of 0.1 per-

cent can be obtained, but the microphones must of course be converted into three-terminal impedances by means of a grounded shield in order to realize this accuracy. The source is the center-tapped output transformer of a commercially available beat-frequency oscillator. The circuit of the high-voltage type bridge easily allows the introduction of the direct-current polarizing voltage needed for condenser microphones.

The resistance arms of the bridge are fixed resistors, each of about 50,000 ohms. The capacitance arms have the unknown condenser, usually about 50  $\mu\mu\text{f}$ , and a standard condenser, variable over the range 40 to 60  $\mu\mu\text{f}$ . Hence  $C_1$  (see fig. 3) is about 100  $\mu\mu\text{f}$ , and  $R_2$  is about 25,000 ohms. The condensers  $c_1$  and  $c_2$  are variable up to about 1,000  $\mu\mu\text{f}$ , and  $r_1$  and  $r_2$  are each approximately 3,000 ohms.

The three-terminal source has  $z_1'$  and  $z_2'$ , which

are mainly resistive, both approximately 800 ohms. The impedance  $z_0$  is also mainly resistive, and is approximately -400 ohms. The voltages  $e_1$  and  $e_2$  are closely equal, both in magnitude and phase, and are about 8 volts for frequencies between 500 and 10,000 c/s.

## V. References

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