

# Secondary $\gamma$ Transitions in $^{159}\text{Gd}$ After Neutron Capture at Isolated Resonances

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## S. Pospišil

Department of Physics, Faculty of Nuclear Sciences and Physical Engineering  
Czech Technical University, Břehova 7,  
115 19 Prague 1, Czech Republic

## F. Bečvář

Department of Low Temperature Physics, Faculty of Mathematics and Physics, Charles University,  
V Holesovickách 2, 180 00 Prague 8,  
Czech Republic

## C. Granja Bustamante and J. Kubašta

Department of Physics, Faculty of Nuclear Sciences and Physical Engineering  
Czech Technical University, Břehova 7,  
115 19 Prague 1, Czech Republic

and

## S. A. Telezhnikov

Frank Laboratory of Neutron Physics,  
Joint Institute for Nuclear Research,  
141 980 Dubna, Russia

The  $^{158}\text{Gd}(n,\gamma)^{159}\text{Gd}$  reaction was studied at 12 isolated neutron resonances by the TOF method at the IBR-30 Fast Pulse Reactor at JINR Dubna. Totally 15 secondary  $\gamma$  transitions in  $^{159}\text{Gd}$  were recorded in the range from 450 keV to 750 keV. Of these, six previously unseen transitions were placed on the established  $^{159}\text{Gd}$  level scheme. The depopulation of strongly populated levels at 507.7 keV and 558.2 keV (the head and the first excited members of band  $1/2^- [521]$ ) was observed for the first time. It was shown that the observed 507.7 keV  $\gamma$  line, masked by the annihilation peak, originates from an unresolved doublet of transitions from the 507.7 keV level to the ground state and

from the 558.2 keV level to the level at 50.7 keV. The 507.7 keV level decays exclusively to the ground state, while the 558.2 keV level decays via two transitions with a branching ratio that agrees well with the prediction according to Alaga's rule.

**Key words:** branching ratio; capture; energy-level transitions; energy levels; gadolinium 159; gadolinium 158 target; gamma cascades; neutron reactions; resonance neutrons.

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## 1. Introduction

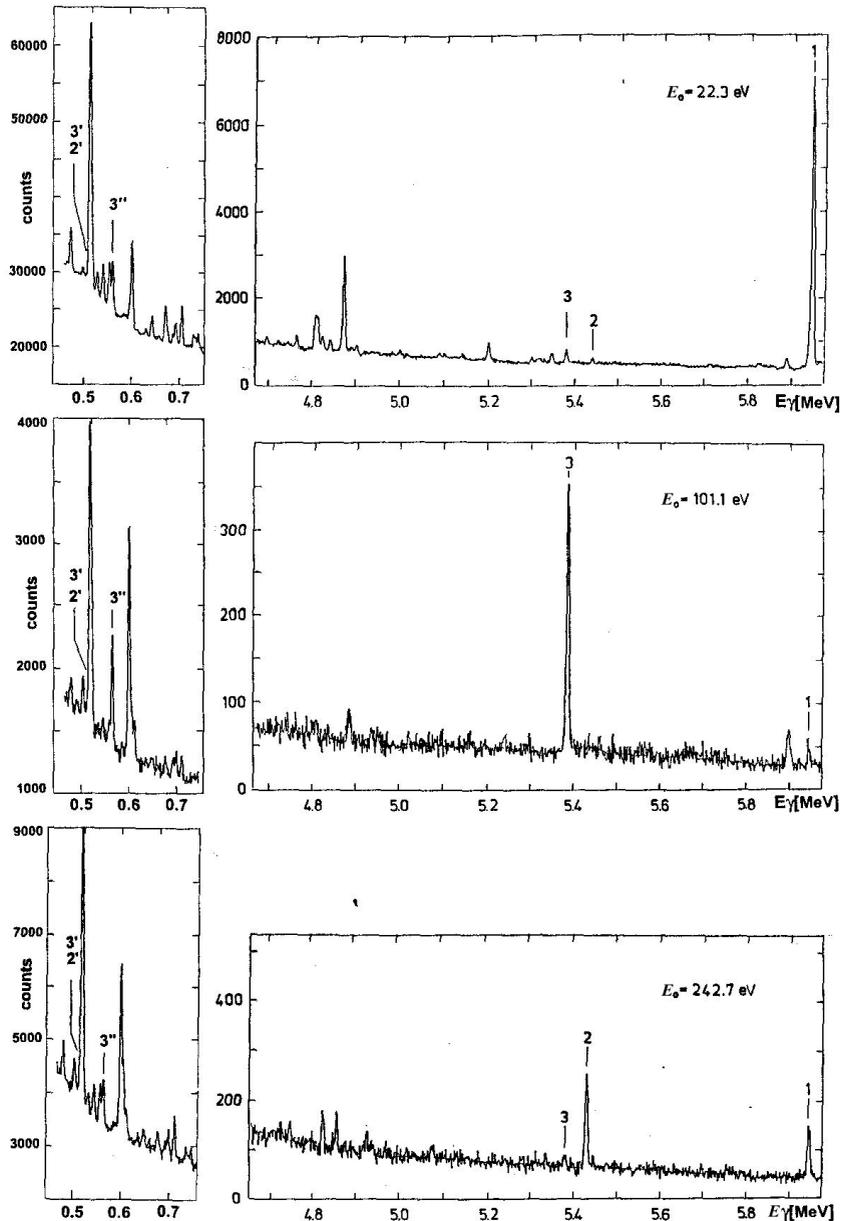
The well deformed  $^{159}\text{Gd}$  nucleus is an ideal ground to test the couplings of rotational and vibrational motion to single particle modes of excitation where the unpaired nucleon (neutron) can be thought of revolving around the deformed nuclear core. While the primary  $\gamma$  transitions in  $^{159}\text{Gd}$  have been studied extensively from radiative neutron capture [1,2,3,4], scarce information about secondary  $\gamma$  transitions is known only from  $\beta^-$ -decay of

$^{159}\text{Eu}$ . No  $(n,\gamma)$  experiment devoted to the study of secondary transitions in  $^{159}\text{Gd}$  nucleus has been done so far. Such reaction is hindered by the extremely high thermal cross sections of the neighboring isotopes, namely  $^{155,157}\text{Gd}$ . This work presents the data on secondary  $\gamma$  transitions in  $^{159}\text{Gd}$  observed following neutron capture at individual resonances of the  $^{158}\text{Gd}$  target.

## 2. Experimental

The experiment was carried out at the JINR Pulsed Fast Reactor IBR-30 which worked as a booster in conjunction with the 40 MeV electron linac LUE-40. The sample consisted of 48.56 g of  $\text{Gd}_2\text{O}_3$  enriched in  $^{158}\text{Gd}$  to 97.7 %. The Time-Of-Flight (TOF) resolution of 70 ns/m (at 50 m flight path) enabled us to accumulate  $\gamma$ -ray spectra for 12 isolated neutron  $s$ -wave resonances with  $J^\pi = 1/2^+$  at energies of (22.3, 101.1, 242.7, 277.2,

344.8, 409.1, 503.3, 588.5, 692.9, 847.3, 917.1, and 1068.0) eV [5]. The resonance at energy 847.3 eV contains a small contribution of the next weak unresolved resonance at energy 869.3 eV. The  $\gamma$ -ray spectra, recorded by means of a Ge(Li) spectrometer, ranged from 450 keV to 750 keV and from 3.5 MeV to 6.0 MeV (neutron separation energy  $B_n = 5943.3$  keV for  $^{159}\text{Gd}$  [1,3,4]). The spectra for the first three neutron resonances are shown in Fig. 1. Relative intensities of  $^{159}\text{Gd}$   $\gamma$  transitions were determined in individual resonances,



**Fig. 1.** Gamma-ray spectra of  $^{159}\text{Gd}$  following neutron capture at 22.3 eV, 101.1 eV, and 242.7 eV resonances. Peaks of primary transitions to the ground state, 507.7 keV and 558.2 keV levels are denoted as 1, 2, and 3, respectively. Peaks of transitions 507.7 keV  $\rightarrow$  g.s., 558.2 keV  $\rightarrow$  50.7 keV and 558.2 keV  $\rightarrow$  g.s. are labelled as 2', 3', and 3'', respectively. Note that intensities of lines 3 and 3'' are correlated.

the total area under seven low-energy  $\gamma$ -lines at (467.2, 524.5, 537.1, 551.0, 601.8, 677.4 and 715.2) keV being adopted as a relative measure of number of neutrons captured at each resonance. A separate run was undertaken using a composite sample consisting of a layer of enriched Gd covered on the back side by a layer of natural boron. Comparing the yields of 477.7 keV  $\gamma$  rays, resulting from the  $^{10}\text{B}(n,\alpha\gamma)$  reaction, and the 601.8 keV  $\gamma$  rays, following neutron capture in  $^{158}\text{Gd}$  at the 22.3 eV resonance, the absolute intensities of  $\gamma$  transitions in  $^{159}\text{Gd}$  were established. This was done by the method described in [6] omitting the role of multiple neutron scattering as the 22.3 eV resonance has a total radiative width  $I_\gamma$  significantly large compared to its neutron width  $I_n$ . The overall uncertainty of the absolute intensity calibration in the low energy region of the spectra was estimated at 18 %.

### 3. Results and Discussion

Fifteen secondary transitions were observed, of which 10 were located in the known level scheme, see Ref. [1,3,4]. Results are summarized in Table 1 where intensities are given for the strongest resonance at 22.3 eV (column 2) together with intensities averaged over

the full set of 12 resonances (column 3). The unresolved  $\gamma$ -ray doublet at 507.7 keV, masked additionally by the neighboring 511 keV annihilation peak, was decomposed into its components belonging to the transitions 558.2 keV  $\rightarrow$  50.7 keV and 507.7 keV  $\rightarrow$  g.s. This was possible thanks to correlations observed among resonances of the intensities of the 5384.7 keV primary transition, which populates the 558.2 keV level, with the intensities of the 507.7 keV  $\gamma$ -ray line (Fig. 2a) as well as between the intensities of the 5434.7 keV primary transition, populating the 507.7 keV level, with the intensities of the same 507.7 keV  $\gamma$ -ray line (Fig. 2b). The strong correlation observed between intensities of the 5384.7 keV primary transition with the intensities of the 558.1 keV secondary transition are shown in Fig. 2c. By analysing these correlations we established that the 558.2 keV level with  $J^\pi = 3/2^-$  decays via two transitions to the 50.7 keV level (51 %) and to the ground state (49 %). We also established that the 507.7 keV level with  $J^\pi = 1/2^-$  decays exclusively to the ground state. By examining the depopulation of the 558.2 keV level (Fig. 2a and 2c) the  $\gamma$ -branching ratio of the 558.2 keV level was determined to be equal  $1.38 \pm 0.26$ . This agrees with the branching ratio of 1.5 following from Alaga's rule [7].

**Table 1.** Secondary  $\gamma$  transitions in  $^{159}\text{Gd}$  observed at isolated neutron resonances

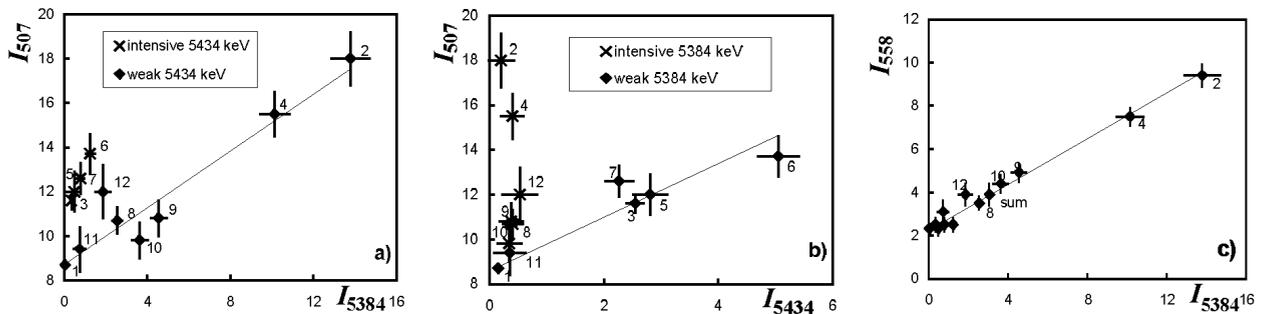
$E_\gamma$ [keV]	$I_\gamma$ [ $\gamma$ per 100 n]	$\langle I_\gamma \rangle$ [ $\gamma$ per 100 n]	$E_i \rightarrow E_f$ [keV]	$E_f$ [keV] <sup>d</sup>
467.2 <sup>a</sup> $\pm$ 0.2	1.58 $\pm$ 0.08	1.70 $\pm$ 0.06	973.7 $\rightarrow$	507.7 <sup>c</sup>
507.7 <sup>a</sup> $\pm$ 0.4	5.59 $\pm$ 0.40	6.15 $\pm$ 0.30	507.7 $\rightarrow$	0.0 <sup>b</sup>
507.7 <sup>a</sup> $\pm$ 0.4	2.36 $\pm$ 0.20	3.61 $\pm$ 0.20	558.2 $\rightarrow$	50.7 <sup>b</sup>
524.5 $\pm$ 0.3	0.95 $\pm$ 0.06	0.87 $\pm$ 0.04		
537.1 $\pm$ 0.2	1.62 $\pm$ 0.07	1.62 $\pm$ 0.03		
551.0 $\pm$ 0.2	2.05 $\pm$ 0.07	1.99 $\pm$ 0.03	602.1 $\rightarrow$	50.7 <sup>c</sup>
558.1 <sup>a</sup> $\pm$ 0.2	2.28 $\pm$ 0.07	3.47 $\pm$ 0.09	558.2 $\rightarrow$	0.0 <sup>b</sup>
601.8 $\pm$ 0.1	4.30 $\pm$ 0.05	4.34 $\pm$ 0.04	602.1 $\rightarrow$	0.0 <sup>c</sup>
646.9 <sup>a</sup> $\pm$ 0.3	1.00 $\pm$ 0.10	0.50 $\pm$ 0.20	647.2 $\rightarrow$	0.0 <sup>c</sup>
677.4 $\pm$ 0.3	1.88 $\pm$ 0.70	1.77 $\pm$ 0.06	744.4 $\rightarrow$	67.8 <sup>c</sup>
682.2 $\pm$ 0.4	0.60 $\pm$ 0.10	0.40 $\pm$ 0.10	732.6 $\rightarrow$	50.7 <sup>c</sup>
700.3 $\pm$ 0.4	0.60 $\pm$ 0.10	0.90 $\pm$ 0.20		
715.2 <sup>a</sup> $\pm$ 0.2	2.08 $\pm$ 0.08	2.33 $\pm$ 0.10	781.8 $\rightarrow$	67.8 <sup>c</sup>
742.4 $\pm$ 0.5	0.70 $\pm$ 0.10	0.60 $\pm$ 0.10		
747.5 $\pm$ 0.6	0.50 $\pm$ 0.10	0.20 $\pm$ 0.10		

<sup>a</sup> Newly observed  $\gamma$ -transitions.

<sup>b</sup> Placements based on correlations between primary and secondary  $\gamma$ -transitions.

<sup>c</sup> Placements based on Level-Fit calculations.

<sup>d</sup> Taken from [1,4].



**Fig. 2.** Correlations between the intensities of secondary and primary transitions. Data for the unresolved 507.7 keV doublet are shown for the primary transitions at 5384.7 keV in a) and at 5434.7 keV in b). Data points for each resonance, numbered from 1 to 12 in ascending order, are marked in a) either by “x” or “♦” when the 5434.7 keV transition is intense or weak, respectively. Analogous marking is used in b) to illustrate the strength of the 5384.7 keV transition. Data for the 558.1 keV and 5384.7 keV transitions are shown in c). Regression lines, calculated by least squares fit, refer to points marked by “♦”. Intensities are given in number of  $\gamma$  per 100 captured neutrons.

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*About the authors:* S. Pospíšil is a senior physicist at the Czech Technical University in Prague. F. Bečvář is a senior physicist at Charles University in Prague. C. Granja Bustamante and J. Kubašta are physicists at the Czech Technical University in Prague. S. A. Telezhnikov is a senior physicist at the Frank Laboratory of Neutron Physics at Dubna.

### 4. References

- [1] R. G. Helmer, <sup>159</sup>Gd Adopted Levels, Gammas. NDS **72**, 83 (1994).
- [2] S. Pospíšil, J. Kubašta, F. Bečvar, H. T. Hiep, and S. A. Telezhnikov, Photon Strength Functions of <sup>159</sup>Gd studied from Neutron Capture at Isolated Resonances in <sup>158</sup>Gd. Acta Polytechnica, Czech Technical University in Prague, Vol. 35, No. 4 (1995) pp. 128-131.
- [3] S. Pospíšil, E. Havrankova, C. Granja Bustamante, J. Kubašta, and S. A. Telezhnikov, Primary Transitions in <sup>159</sup>Gd Studied at Isolated Neutron Resonances of <sup>158</sup>Gd. Proceedings of the 9th International Symposium on Capture Gamma Ray Spectroscopy and Related Compounds, G. L. Molnar, T. Belgya, and Zs. Revay, eds., Vol. 1, Springer Hungarica, Budapest, Hungary (1997) pp. 432-433.
- [4] C. Granja Bustamante, Energy Levels of <sup>159</sup>Gd and <sup>174</sup>Yb after Radiative Neutron Capture. Diploma thesis, Czech Technical University, Prague (1997).
- [5] S. F. Mughabghab, M. Divadeenam, and N. E. Holden, Neutron Cross Sections, Neutron Resonance Parameters and Thermal Cross Sections. Vol. 1, part B: Z=61-100, BNL-325 Tables, Academic Press (1981).
- [6] F. Bečvar, M. E. Montero-Cabrera, S. Pospíšil, and S. A. Telezhnikov, Yad. Fiz. **44**, 1 (7) (1986) pp. 3-9.
- [7] G. Alaga, K. Alder, A. Bohr, and B. Mottelson, Mat. Fys. Medd. Dan. Vid. Selsk. **29**, No. 9 (1955).