

On Cauchy-Riemann Equations in Higher Dimensions¹E. Stiefel²The n linear partial differential equations with constant complex coefficients

$$l_j = \sum_{k=1}^n a'_{jk} \frac{\partial u_k}{\partial x_j} = 0,$$

($j=1, \dots, n$) are said to form a system of generalized Cauchy-Riemann equations, if there exist constants f'_k such that

$$\Delta u_j = \sum_{k=1}^n f'_k \frac{\partial l_k}{\partial x_j}.$$

It is proved that such systems exist for $n=1, 2, 4, 8$ only. In the cases $n=2, 4$ there are three essentially inequivalent systems; $n=8$, only two. If the coefficients are required to be real, there exist only the classic system of two equations, the two systems of Dirac-Fueter equations, and two systems of eight equations.

If two real functions u_1, u_2 of the real variables x_1, x_2 satisfy the Cauchy-Riemann equations

$$\frac{\partial u_1}{\partial x_1} - \frac{\partial u_2}{\partial x_2} = 0, \quad \frac{\partial u_2}{\partial x_1} + \frac{\partial u_1}{\partial x_2} = 0, \quad (1)$$

they are harmonic, that is, from (1) follows Laplace's equation by differentiation:

$$\Delta u_1 = \frac{\partial^2 u_1}{\partial x_1^2} + \frac{\partial^2 u_1}{\partial x_2^2} = 0, \quad \Delta u_2 = \frac{\partial^2 u_2}{\partial x_1^2} + \frac{\partial^2 u_2}{\partial x_2^2} = 0. \quad (2)$$

Introducing the left sides of the Cauchy-Riemann equations

$$l_1 = \frac{\partial u_1}{\partial x_1} - \frac{\partial u_2}{\partial x_2}, \quad l_2 = \frac{\partial u_2}{\partial x_1} + \frac{\partial u_1}{\partial x_2}, \quad (3)$$

we observe that this statement is an immediate consequence of the relations

$$\Delta u_1 = \frac{\partial l_1}{\partial x_1} + \frac{\partial l_2}{\partial x_2}, \quad \Delta u_2 = \frac{\partial l_2}{\partial x_1} - \frac{\partial l_1}{\partial x_2}, \quad (4)$$

that is to say, that the Laplacians of u_1, u_2 are linear combinations of the derivatives of the left sides of the Cauchy-Riemann equations.

In 1939 Olga Taussky-Todd³ studied the following general problem. Let $u_1, u_2, \dots, u_n$ be functions of the independent variables $x_1, x_2, \dots, x_n$. Is it possible to find a system of n linear partial differential equations with constant coefficients

$$l_j = \sum_{k=1}^n a'_{jk} \frac{\partial u_k}{\partial x_j} = 0, \quad (5)$$

in such a way that

$$\Delta u_j = \sum_{k=1}^n b'_{jk} \frac{\partial l_k}{\partial x_j}, \quad (6)$$

the b'_{jk} again being constant coefficients? If a set of functions $u_1, u_2, \dots, u_n$ satisfies (5), it follows then from (6) that they are harmonic. So we may say that (5) are *Cauchy-Riemann equations in n -dimensional space* and generate a *theory of functions in this space reasonably related to potential theory*. O. Taussky proved that this problem can only be solved in spaces of dimension $n=2^n$. In this paper the better result is established that n must be 1, 2, 4, or 8 and moreover all Cauchy-Riemann systems (5) will be classified. In our discussion we admit that the x_j, u_k and the coefficients in (5), (6) are complex. We will use methods of representation theory introduced by Wigner and Eckmann⁴ for the solution of problems of an analogous type, but we shall simplify matters a little by dealing with algebras instead of groups.

1. Introducing the n -row matrices

$$A_i = (a'_{ik}), \quad B_k = (b'_{jk}), \quad i, k = 1, 2, \dots, n, \quad (7)$$

and the vectors

$$l = (l_1, l_2, \dots, l_n), \quad u = (u_1, u_2, \dots, u_n), \quad (8)$$

relations (5), (6) can be written

$$l = \sum_{i=1}^n A_i \frac{\partial u}{\partial x_i}, \quad \Delta u = \sum_{k=1}^n B_k \frac{\partial l}{\partial x_k} \quad (9)$$

Inserting the second equation into the first we get

$$\Delta u = \frac{\partial^2 u}{\partial x_1^2} + \frac{\partial^2 u}{\partial x_2^2} + \dots + \frac{\partial^2 u}{\partial x_n^2} = \sum_{k=1}^n B_k A_k \frac{\partial^2 u}{\partial x_k \partial x_k} \quad (10)$$

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² University of California, Los Angeles, Cal., and Eidgenössische Technische Hochschule, Zürich, Switzerland.

³ O. Taussky, An algebraic property of Laplace's differential equation, *Quart. J. Math.* 10, 66 (1939).

⁴ B. Eckmann, Gruppentheoretischer Beweis des Satzes von Hurwitz-Radon über die Komposition quadratischer Formen, *Comment. Math. Helv.* 15, 368 (1942).

This identity holds for every vector u . Comparing coefficients it turns out that

$$B_i A_i = I, \quad B_h A_i + B_i A_h = 0, \quad i \neq h, \quad (11)$$

I being the n -row unit matrix. Thus the matrix B_i is the inverse of A_i , and our question is reduced to the problem of constructing n matrices A_i of n rows satisfying the relation

$$A_i^{-1} A_i + A_i^{-1} A_h = 0, \quad i \neq h, \quad (12)$$

2. In order to solve this problem we observe that (12) is invariant under a general equivalence transformation

$$A_i \rightarrow S A_i T, \quad i=1, 2, \dots, n, \quad (13)$$

where S, T are two matrices with nonvanishing determinants. We do not distinguish between two Cauchy-Riemann systems (5) related to each other by such a transformation but call them essentially equal. Using this equivalence, we may transform by (13) one of our matrices—say A_n —into the unit matrix. For $h=n$ we have then

$$A_i + A_i^{-1} = 0 \quad \text{or} \quad A_i^{-1} = -A_i \quad \text{or} \quad A_i^2 = -I$$

and for $h=1, 2, \dots, n-1$ this gives

$$A_h A_i + A_i A_h = 0.$$

Thus we may restrict ourselves to the problem of constructing $(n-1)$ matrices $A_1, A_2, \dots, A_{n-1}$ having the property

$$A_i^2 = -I, \quad A_h A_i + A_i A_h = 0, \quad i \neq h. \quad (14)$$

If we have only these special systems under consideration the general equivalence transformation (13) will be restricted to a similarity transformation

$$A_i \rightarrow S A_i S^{-1}, \quad (15)$$

because the unit matrix must be left invariant. Matrices of the type (14) have been studied at first by Hurwitz⁵ in the special case where the A_i are real and orthogonal.

3. From the basic relations (14) it follows that the 2^{n-1} matrices

$$\left. \begin{aligned} I, A_1, A_2, \dots, A_{n-1}, \\ A_1 A_2, A_1 A_3, \dots, A_{n-2} A_{n-1}, \\ A_1 A_2 A_3, \dots, A_{n-3} A_{n-2} A_{n-1}, \\ \dots \\ A_1 A_2, \dots, A_{n-1} \end{aligned} \right\} \quad (16)$$

this is to say, all products with increasing subscripts of the factors, form a matrix algebra of order 2^{n-1} . Indeed, the product of two matrices of the set (16) is (up to the sign) again a matrix of the set. Let us now consider the abstract associative algebra H of order 2^{n-1} over the complex field given by the basic elements

$$\left. \begin{aligned} 1, e_1, e_2, \dots, e_{n-1}, \\ e_1 e_2, e_1 e_3, \dots, e_{n-2} e_{n-1}, \\ e_1 e_2 e_3, \dots, e_{n-3} e_{n-2} e_{n-1}, \\ \dots \\ e_1 e_2, \dots, e_{n-1} \end{aligned} \right\} \quad (17)$$

and the multiplication rules

$$e_i^2 = -1, \quad e_h e_i + e_i e_h = 0, \quad i \neq h. \quad (18)$$

Thus our problem is finally to construct a representation of the algebra H by n -row matrices. In order to do this we use the following well-known theorems of representation theory:

Theorem I. There is—up to similarity transformation (15)—only a finite number m of irreducible representations, where m is the order of the center of the given algebra H . Any representation is the sum of irreducible representations.

Theorem II. Let f be the degree of a representation (number of rows of the representing matrices). Then the degrees $f_1, f_2, \dots, f_m$ of the irreducible representations satisfy the relation

$$f_1^2 + f_2^2 + \dots + f_m^2 = \text{order of } H = 2^{n-1}.$$

4. Let us discuss first the case that the number n of Cauchy-Riemann eq (5) is even. Then the last element $(e_1 e_2 \dots e_{n-1})$ of the sequence (17) commutes with $e_1, e_2, \dots, e_{n-1}$ and is therefore a center element of the algebra H . It is not difficult to show that the elements 1 and $(e_1 e_2 \dots e_{n-1})$ span the center of H , that is to say that the general center element is of the form $\alpha + \beta(e_1 e_2 \dots e_{n-1})$, where α, β are complex numbers. The order of this center being $m=2$, we learn from theorem I, that our algebra H has exactly two irreducible representations, D_1 and D_2 . They are related in the following way. If D_1 is given by

$$D_1: \quad 1 \rightarrow I, \quad e_i \rightarrow E_i, \quad (19)$$

(E_i being the representing matrices) then D_2 is given by

$$D_2: \quad 1 \rightarrow I, \quad e_i \rightarrow -E_i. \quad (20)$$

In order to prove this let us observe that if the E_i satisfy the basic relations (14)

$$E_i^2 = -I, \quad E_h E_i + E_i E_h = 0,$$

⁵ A. Hurwitz, Über die Komposition der quadratischen Formen, collected papers 3, 641 (1903).

the same is true for the matrices $(-E_i)$. Hence, if (19) is a representation, then (20) is another. Furthermore, if D_1 is irreducible, the same is true for D_2 and, finally, D_1, D_2 are not similar. Indeed, the center element $(e_1 e_2 \dots e_{n-1})$ is represented in D_1 by the matrix $E_1 E_2 \dots E_{n-1}$, which must be a multiple cI of the unit matrix, since it commutes with the whole irreducible set of representing matrices. In D_2 the representing matrix is

$$-E_1 E_2 \dots E_{n-1} = -cI.$$

But (cI) and $(-cI)$ are not similar because, certainly, $c \neq 0$. This finishes our proof.

This discussion shows, in particular, that D_1 and D_2 have the same degree f . From theorem II we have

$$2f^2 = 2^{n-1} \quad \text{and so} \quad f = 2^{\frac{n-2}{2}}. \quad (21)$$

As stated in section 3, our basic problem is to find a representation D of the algebra H by n -row matrices. From theorem I it follows that D must be the sum of the representations D_1, D_2 , each of them perhaps repeated several times. So the degree n of D is a multiple of f :

$$n = k \cdot 2^{\frac{n-2}{2}}. \quad (22)$$

But this is only possible for $n=2, 4, 8$, and the multiplicities k are 2, 2, 1, respectively. In the cases $n=2, 4$ we have for D the three inequivalent possibilities

$$D_1 + D_1, \quad D_1 + D_2, \quad D_2 + D_2 \quad (23)$$

and in the case $n=8$ the two possibilities D_1, D_2 .

5. The case of odd n is rather trivial because the center of H in this case is formed only by the multiples of the unity element. So we have only one irreducible representation. Its degree f is given by

$$f^2 = 2^{n-1}, \quad f = 2^{\frac{n-1}{2}},$$

and the wanted representation D must be a multiple of this unique irreducible representation:

$$n = k \cdot 2^{\frac{n-1}{2}}.$$

This leads to $n=1$.

Collecting the results we get the theorem:

A system of n Cauchy-Riemann equations of the type (5), (6) is only possible for $n=1, 2, 4, 8$. In the cases $n=2, 4$, there are three inequivalent systems, in the case $n=8$, only two.

6. In this section we establish the Cauchy-Riemann systems explicitly and discuss especially the real ones.

(a) In the case $n=2$ the degree f of the irreducible representation D_1 is $f=1$ according to (21). So D_1, D_2 may be given by

$$D_1: e_1 \rightarrow i, \quad D_2: e_1 \rightarrow -i, \quad (24)$$

i being the imaginary unit. Our first possibility (23) $D=D_1+D_1$ is then

$$e_1 \rightarrow \begin{pmatrix} i & 0 \\ 0 & i \end{pmatrix} = A_1, \quad A_2 = I \quad (25)$$

and yields the Cauchy-Riemann system (5)

$$i \frac{\partial u_1}{\partial x_1} + \frac{\partial u_1}{\partial x_2} = 0, \quad i \frac{\partial u_2}{\partial x_1} + \frac{\partial u_2}{\partial x_2} = 0 \quad (26)$$

consisting of two separate equations for u_1 and u_2 . Hence, we may restrict ourselves to the single equation

$$i \frac{\partial u}{\partial x_1} + \frac{\partial u}{\partial x_2} = 0, \quad (27)$$

which expresses the fact (if x_1, x_2 are real variables) that u is a complex analytic function of $x_1 - ix_2$. By differentiation of (27) it follows of course $\Delta u = 0$. The two other possibilities (23) may be established by changing i into $(-i)$ in one or both equations (26).

In order to find the real Cauchy-Riemann systems—that is to say real representations of H —we must form the sum of one of our irreducible representations and its complex conjugate representation. Taking into account that D_1 and D_2 are complex conjugates, we have finally only the unique real representation $D=D_1+D_2$ and only one real Cauchy-Riemann system, which is, of course, the system (1).

(b) For $n=4$ the representation D_1 may be given by the so-called Pauli-matrices

$$D_1: e_1 \rightarrow E_1 = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix}, \quad e_2 \rightarrow E_2 = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}, \\ e_3 \rightarrow E_3 = \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} \quad (28)$$

with

$$e_1 e_2 e_3 \rightarrow \begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix}. \quad (29)$$

D_2 is obtained by changing the sign of those matrices. The Cauchy-Riemann system corresponding to $D=D_1+D_1$ splits again into the equations

$$\left. \begin{aligned} i \frac{\partial u_2}{\partial x_1} - \frac{\partial u_2}{\partial x_2} + i \frac{\partial u_1}{\partial x_3} + \frac{\partial u_1}{\partial x_4} &= 0 \\ i \frac{\partial u_1}{\partial x_1} + \frac{\partial u_1}{\partial x_2} - i \frac{\partial u_2}{\partial x_3} + \frac{\partial u_2}{\partial x_4} &= 0 \end{aligned} \right\} \quad (30)$$

for u_1, u_2 alone and the same equations for u_3, u_4 . In this case, however, D_1 and D_2 are not conjugate complex, but D_1 is similar to its own conjugate complex and the same is true for D_2 . This follows from the fact that in D_1 the center-element $e_1 e_2 e_3$ is represented by a real matrix according to (29). So we have two nonequivalent real Cauchy-Riemann systems corresponding to the representations $D_1 + D_1$ and $D_2 + D_2$. They are R. Fueter's equations for

right and left regular functions of a quaternion variable, and they are closely related to Dirac's equations in quantum mechanics. The first system may be derived from our equations (30) in the following way. Let x_1, x_2, x_3, x_4 be real variables and u_1, u_2 complex functions:

$$u_1 = v_1 + i v_2, \quad u_2 = v_3 + i v_4.$$

Splitting the eq (30) into their real and imaginary parts we get the four equations wanted for v_1, v_2, v_3, v_4 . The second system follows in the same way if we replace (30) by the equations corresponding to $D_2 + D_2$.

(c) The case $n=8$ is entirely different from the previous cases, because we found in section 4 that the representation D solving our problem is either D_1 or

D_2 and hence irreducible. Thus the Cauchy-Riemann systems of this case will not split into systems of fewer equations. We omit the computation of the matrices of D_1 , which is closely related to the so-called Cayley numbers building a nonassociative algebra and mention only the result that those matrices may be constructed as real matrices. They yield two Cauchy-Riemann systems.

As a final result we have the following statement: *The only generalized real Cauchy-Riemann systems are 1. the classic system of two equations; 2. the two systems of Dirac-Fueter equations, each system having four equations; and 3. two systems of eight equations.*

LOS ANGELES, December 4, 1951.

