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Comparison of Theoretical and Experimental Photoeffect Data 0.1 keV to 1.5 MeV

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Graphs of calculated and experimental atomic photoeffect cross sections as a function of photon energy 0.1 keV to 1.5 MeV are presented for all elements Z = 1 to 94. The calculated results presented are (a) the nonrelativistic Hartree-Fock self-consistent-field (SCF) results of Veigele, Henry, et al., over the range 0.1 keV to between 1.0 and 8.0 keV for all elements Z = 1 to 94 and (b) the relativistic Hartree-Slater SCF results of Scofield over the range 1.0 keV to 1.5 MeV for all elements Z = 1 to 101. The "experimental" data-points are derived by subtracting theoretical scattering cross sections from total attenuation coefficient measurements in the literature. Differences between theoretical and experimental photoeffect data are typically a factor of two from 0.1 to 1.0 keV, 5-10% from 1.0 to 5.0 keV and 1-5% from 5.0 keV up to energies, ranging from 20 keV for carbon up to 500 keV for lead, above which the photoeffect cross section becomes fractionally too small to be accurately determined from the total attenuation coefficient.

Key words: Attenuation coefficients, cross sections, gamma rays, photoelectric effect, photons, x-rays.

1. Introduction

The atomic photoeffect, in which an atom absorbs a photon and ejects an electron, is the predominate process by which photon radiation is attenuated by matter in the vacuum ultraviolet (VUV) (10 eV to 100 eV) and soft x-ray (100 eV to 10 keV) portions of the electromagnetic spectrum. At higher energies the attenuation is dominated by two other processes, first incoherent scattering commencing in importance at around 10 keV for the lowest-Z elements and at a few hundred keV for the highest-Z elements,

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then pair production which proceeds from a threshold of 1.022 MeV $(2 \text{ m} \text{ c}^2)$ to dominate over all other processes in the high energy limit.

As an example, figure 1 shows the total measured attenuation cross section (small circles) for photons from 10 eV to 100 GeV in copper, decomposed into the theoretical contributions from the above three processes: photoeffect (T), incoherent (bound-electron Compton) scattering (σ_{INCOH}) and pair production (K), indicating the region where each predominates. For lower-Z elements the incoherent scattering cross section dominates over a broader range (for hydrogen, incoherent scattering dominates the total cross section from a few keV to 100 MeV) and for higher-Z elements the range of σ_{INCOH} dominance narrows to the point where for plutonium (Z = 94) its contribution to the total reaches a maximum of 70% at 1.5 - 2.0 MeV and its relative importance rapidly decreases toward lower and higher energies. Also shown in figure 1 are two additional processes: coherent scattering (σ_{COH}) and the photonuclear giant resonance ($\sigma_{PH,N}$) of some significance but which contribute less than 10% to the total cross section at any given energy.

Pratt, Ron and Tseng $[1]^{\perp}$ have recently extensively reviewed the existing calculations and measurements of the atomic photoeffect over the range 10 keV to 100 MeV (including material extending down to 1 keV), complementing an earlier review by Fano and Cooper [2] of theoretical and measured photoeffect data over the range 10 eV to 10 keV.

In the latter (soft x-ray) region, where the atomic photoeffect can be taken directly from measured total attenuation coefficients with only small or negligible corrections for the scattering contribution, additional recent reviews and surveys of experimental photoeffect cross section data include those by Henke and Tester [3], Hudson and Kieffer [4], [5] and Zimkina and Fomichev [6]. Higher-energy photoeffect measurements, both by direct methods (e.g., photoelectron and fluorescence x-ray coincidences) and by indirect methods (subtraction of scattering and pair-production contributions from total attenuation measurements) have also been recently reviewed by Parthasaradhi [7].

In the above and other recent photoeffect reviews, the graphical and/ or numerical comparisons of theoretical and photoeffect data have in each case been for a limited number of specific elements of interest, over limited ranges of photon energy. Thus we consider it a useful adjunct to the above reviews, for an over-all perspective, to systematically compare available experimental total and photoeffect cross sections in the extensive data-collections of Veigele [8] and Hubbell [9] with present state-ofthe-art theoretical photoeffect cross sections for all elements Z = 1 to 94 over the photon energy range 0.1 keV to 1.5 MeV.

For this purpose we have selected the low-energy photoeffect calculations by Veigele, Henry and Bates [8], [10] for Z = 1 to 94 and photon energies 0.1 keV to from 1.0 to 8.0 keV, and the higher-energy calculations

Figures in brackets indicate the literature references at the end of this paper.

by Scofield [11] for Z = 1 to 101, 1.0 keV to 1.5 MeV. These and some additional photoeffect calculations are listed and described in the following section.

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2. Photoeffect Calculations

Prior to 1960, as discussed in more detail in the review article by Pratt et al. [1], theoretical numerical tabulations of the atomic photoeffect cross section, including those in the attenuation coefficient tables of Davisson and Evans [12] and of Grodstein [13], were based primarily on work (reviewed by Hall [14]) dating from the period 1930 to 1935:

(1) The Stobbe [15] exact nonrelativistic K- and L-shell results in the point-Coulomb potential,

(2) The Sauter [16], [17] relativistic K-shell Born result in the point-Coulomb potential,

(3) The Hall [18] and Hall and Rarita [19] high-energy limit results for the K- and L-shells in the point-Coulomb potential,

(4) Exact relativistic point-Coulomb K-shell numerical results for six cases (Z = 26, 50 and 84 at 0.35 and 1.13 MeV) by Hulme, McDougall, Buckingham and Fowler [20], and

(5) A semi-empirical prescription by Slater [21] to take into account inner-electron screening effects.

Renewed interest in theoretical understanding of the photoeffect process began in the late 1950's and continues to the present, resulting from (a) increasingly exacting photon-interaction needs by the medical, scientific and engineering communities, (b) availability of high-speed computers and (c) new higher-accuracy photoeffect and total attenuation coefficient measurements. New analytical work in this period includes corrections by Erber [22], Pratt [23,24] and by Gorshkov and Mikhailov [25] to the theoretical high-energy limit, extension of the relativistic Born approximation by Gavrila [26,27] and Nagel [28], dipole-approximation calculations in the ultra-soft x-ray region by Cooper [29], estimates of the quadrupole contribution by Guttmann and Wagenfeld [30], and additional work more exhaustively reviewed by Pratt et al. [1] and Fano and Cooper [2].

Based on the above work and further refinements, recent pilot and systematic numerical calculations of the photoeffect cross section include the following, listed approximately in chronological order:

(1) Hultberg et al. [31], [32]: relativistic, point-Coulomb, K shell, 21 elements Z = 1 to 100, 1 keV to 10 MeV,

(2) Pratt et al. [33]: relativistic, point-Coulomb, K shell, 6 elements Z = 13 to 92, 200 keV to 2 MeV, (3) Alling and Johnson [34]: relativistic, point-Coulomb, K-L-shells, Pb and U, 81 keV to 1.33 MeV,

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(4) Matese and Johnson [35]: relativistic, exponential screening, K-, L-shells, Ag, Hg, Pb and U, 81 keV to 1.33 MeV,

(5) Hall and Sullivan [36]: relativistic, Thomas-Fermi screening, K-shell, Fe, Sn, Po and U, 300 keV to 1.23 MeV,

(6) Rakavy and Ron [37], [38]: relativistic, Fermi-Amaldi screening for Al, Fe, Sn and W, Thomas-Fermi screening for U, 1 keV to 2 MeV,

(7) Schmickley and Pratt [39]: relativistic, Hartree-Fock-Slater screening, 14 elements Z = 13 to 92, 10 keV to 5 MeV,

(8) Bell and Kingston [40], [41]: non-relativistic, dipole-length and -velocity formulations, Hartree-Fock screening, He, 25 to 310 keV,

(9) Combet-Farnoux and Héno [42], [43]: non-relativistic, Hartree-Fock-Slater screening, Ta, Pt, Au and Bi, 50 to 500 eV,

(10) Manson and Cooper [44]: non-relativistic, Hartree-Fock-Slater screening, 9 elements Z = 18 to 100, 10 eV to 2 keV,

(11) Brysk and Zerby [45]: relativistic, Dirac-Slater screening, Al and U, 1 to 150 keV,

(12) McGuire [46]: non-relativistic, dipole-length formulation, Hartree-Fock-Slater screening, Z = 2 to 54, 6 eV to 62 keV,

(13) Amusia et al. [47], [48]: non-relativistic, Hartree-Fock screening, Ar, Kr and Xe, 15 eV to 1 keV,

(14) Storm and Israel [49]: Brysk-Zerby program [45], Z = 1 to 100, 1 to 200 keV; extended table to 100 MeV by interpolation and extrapolation of Schmickley and Pratt [39] and Rakavy and Ron [38] theoretical results, above,

(15) Cromer and Liberman [50]: relativistic, Dirac-Slater, Kohn-Sham screening, Z = 3 to 98, 5.4 to 22.2 keV,

(16) Veigele et al. [8], [10]: non-relativistic, Hartree-Fock-Slater screening, Z = 1 to 94, 100 eV to 8 keV,

(17) Kennedy and Manson [51]: non-relativistic, Hartree-Fock and Hartree-Slater, dipole-length and -velocity formulations, Ne, Ar, Kr and Xe, 10 eV to 1.5 keV,

(18) Scofield [11]: relativistic, Hartree-Slater screening, Z = 1 to 2° 101, 1 keV to 1.5 MeV, and

(19) Hildebrandt et al. [52], [53]: non-relativistic, revised Slater screening constants, dipole and quadrupole components given, Z = 6 to 54, 5 to 25 keV.

In the most recent NBS x-ray cross section tabulation (NSRDS-NBS 29, covering the energy range 10 keV to 100 GeV) [54], photoeffect cross sections in the range 10 to 100 keV were obtained from total attenuation coefficient measurements [9] by subtracting theoretical scattering contributions [49] except for uranium which was taken directly from the above Rakavy-Ron [37], [38] theoretical results. Above 100 keV, for all 23 elements in the NSRDS-NBS 29 tabulation [54], the Rakavy-Ron results were combined with the high-energy numerical calculations and analytical formulas of Gavrila [26], Pratt et al. [23], [33] and Hultberg [32] to extend the tabulations up to 100 MeV.

In the McMaster et al. [55], [56] LLL (Lawrence Livermore Lab.) - NBS x-ray cross section tabulation, source of the 1 keV-1 MeV photoeffect values in the present ENDF (Evaluated Nuclear Data File, ERDA) tape library [57], the theoretical values of Schmickley and Pratt [39] were weighted into the iterative cubic log-log least-squares fit of photoeffect cross sections which were otherwise obtained by subtraction of theoretical scattering contributions from total attenuation coefficient measurements.

3. Explicit Photoeffect Measurements

In the photon energy range \sim 100 keV to 2.6 MeV, where the above "subtraction method" (total attenuation coefficient measurement minus theoretical scattering contribution) ceases to be useful, there exist a few explicit measurements of total and subshell photoeffect cross sections, as summarized by Pratt et al. [1] and by Parthasaradhi [7], with uncertainties of the order of 4% to 15% or greater. Additional new measurements beyond those listed in these summaries include:

(1) Sahota [58]: L-shell, fluorescence x-ray detection, U, 145, 192, 662 and 1250 keV,

(2) Allawadhi and Sood [59], [60]: K-shell, fluorescence x-ray detection, 15 elements Z = 33 to 74, 37 and 74 keV,

(3) Gowda and Sanjeevaiah [61] - [65]: total and K-shell, photoelectron detection, 8 elements Z = Al to Pb, 145, 279, 412 and 662 keV, and

(4) Goswami and Chaudhuri [66]: total photoeffect, cylindrical-shell adaptation of the spherical-shell transmission method of Ghose [67], Pb, 662 and 1250 keV.

Because of the substantial inherent difficulties and uncertainties in the above and earlier explicit measurements, this source of photoeffect data has as yet been given little weight, as compared with theory, in constructing photon cross section tabulations, and is not included in the graphical comparisons in figures 2 to 95.

> 4. Graphical Comparison of Experimental Points with Veigele et al. and Scofield Theoretical Results

As can be seen in figure 1, the photoeffect cross section, τ , decreases approximately as E^{-3} from the K absorption edge up to 1 MeV, with a similar \sqrt{E} dependence at lower energies in the segments between absorption edges

down to 1 keV. Above 1 MeV the photoeffect behavior is seen changing to the high-energy-limit E⁻¹ dependence [17], [22], [23]. Below 1 keV the slope of the cross section again decreases, and the cross section, particularly for medium- and high-Z elements, becomes characterized by outer-subshell maxima and minima with little interpolable regularity as a function of E and Z.

In the region 1 keV to 1 MeV where τ decreases rapidly as a function of E, relative differences between photoeffect data are difficult to see on an ordinary log-log plot of τ vs. E as in figure 1. Thus, in figures 2 to 95 we have plotted $\tau \cdot \tilde{E}^2$, instead of τ , which reduces the number of logarithmic cycles required to display the data and has the effect of magnifying the relative differences. The quantity $\tau \cdot E^3$ is less suitable for graphing purposes because of the absorption edges which, for all but the lightest elements, superimpose an effective positive power of E on the energy-range 1 keV to 1 MeV as a whole for a net over-all dependence $\tau \propto v E^{-2}$ including absorption edges.

In figures 35-95 (Z \geq 34) the number of logarithmic cycles used was determined for each figure by the $\tau \cdot E^2$ maximum and minimum values from the Scofield data in the range 1.0 keV to 1.5 MeV. As a result, some experimental points and portions of the Veigele et al. calculated results near 0.1 keV are off-scale and not shown in these figures.

For elements H(Z = 1) to K(Z = 19), in figures 2 to 20, the Veigele et al. calculated results extend only over the range 0.1 to 1.0 keV and hence have only the 1.0-keV point in common with the 1.0 keV to 1.5 MeV range of the Scofield results. For the remaining elements Ca(Z = 20) to Pu(Z = 94) the Veigele et al results overlap the Scofield results by amounts ranging from 1.0 - 1.6 keV for Kr(Z = 36), Rb(Z = 37) and Sr(Z = 38) up to 1.0 - 8.0 keV for Pu (Z = 94).

Since the Veigele et al. and Scofield values in this overlap region are not identifiable as to which is which, numerical values in kilobarns/atom (without the factor $[E(keV)^2]$) at 1.0 and 2.0 keV are listed in table I. Also in table I are percent differences obtained at 100% x (V-S)/S, where V and S are the Veigele et al. and Scofield values, respectively, for a given element and energy.

Comparing the Veigele et al. results with those of Scofield at 1 keV, the agreement is seen to be better than 10% for all elements except for -16%for Cs(Z = 55) and +31% for Nd(Z = 60), the latter two differences attributable to displacements in absorption-edge energies. For elements Z = 1 to 59 (excluding Z = 55) at 1 keV, where the photoeffect cross section is primarily due to the K-, L- and M-shell atomic electrons, the results of Veigele et al. are systematically lower than those of Scofield by an average of 2%. For the higher-Z elements Z = 61 to 94, where the photoeffect cross section in the overlap region is primarily due to N-shell electrons, the trend reverses, with the Veigele et al. results systematically higher than Scofield by an average of 6% at 1 keV for Z = 61 to 89, and higher by an average of 9% at 2 keV for Z = 90 to 94. The experimental points (x's) in figures 2 to 95 are taken from the automated 8000-data-point file of Veigele et al. obtained from 153 sources for the period 1920 - 1970 cited and cross-referenced as to element in reference [8], updated to Jan. 1974. Excluded from this file are data deduced for elements from measurements made on compounds, and a number of discrepant data-points. An additional relevant measurement-set reported by Del Grande and Oliver [68] for uranium at 1 to 10 keV is shown as the circles in figure 93.

A characteristic feature of figures 2 to 95 is the diverging scatter of experimental points toward higher energies. This divergence, commencing at a few hundred keV for the heaviest elements, occurs when the photoeffect, obtained by subtracting theoretical scattering cross sections from total attenuation coefficients, becomes small and similar in magnitude to the uncertainties in the latter two quantities. Thus figures 2 to 95 indicate, as a function of Z, approximate upper energy limits for obtaining meaningful values of the photoeffect cross section from total attenuation coefficient measurements.

5. Conclusions

In the graphs in figures 2 to 95, the theoretical results of Veigele et al. in the region 0.1 to 1.0 keV are seen to represent the experimental data to within typically 10% to a factor of two, the best agreement being found for elements available as gases. Differences of this order are attributable at least in part to chemical binding effects on the outer shell electrons which determine the photoionization cross section in this energy region.

In the region 1.0 to 5.0 keV the Veigele et al. and Scofield values represent the experimental data to within typically 5% to 10%. For the high-Z elements where the theoretical sets systematically differ from each other by 5% to 10% in this energy region, the available measurements for Au, Pb, Bi and U suggest that the theoretical sets would be best joined by taking the Veigele et al. set below and the Scofield set above the M₅ absorption edge.

Above 5 keV, up to energies ranging from 20 keV for carbon up to 500 keV for lead, the Scofield data, with a few exceptions, can be seen to represent the available experimental data to within the 1% to 5% "envelope of uncertainty" obtainable from measurements of the total attenuation coefficient. Among the exceptions is tellurium (figure 53) for which the Scofield values are systematically 10% to 30% lower than the extensive measurements of Biermann (1936) [69], Schulz (1936) [70] and Wrede (1939) [71] over the range 1.3 to 97 keV, but for which the Scofield values agree within 1% with the L-edge cross section measurements by Nordfors and Noreland (1961) [72] in the region 4.3 to 5.0 keV.

The above (Te) and other examples in figures 2 to 95 point out elements and energy-regions where further attenuation coefficient measurements, taken with modern detectors and sample-purities, would afford a more critical test of present theory and would reduce the uncertainties in present photon cross section tabulations.

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Table I. Comparison of ¹Veigele et al. [8], [10] and Scofield [11] theoretical photoeffect cross section values (in kilobarns/atom) at 1.0 and 2.0 keV. The Veigele et al. values in parenthesis for 2 keV photons in Z = 70, 71 and 72 are interpolated from adjacent energies.

	1.0 keV			2 0 kov			
7	Voicele et el	Secticid	9 1166	Voicele et el	Coofdold	9 11EE	
4	vergere et ar.	T (Lb)	% ulli.	vergere et ar.	Scorleid	% dlli.	
	((KD)	((KD)		((KD)	(KD)		
1	0.0109	0 011414	-4 5%	(Veigele et al		1 0 k v	
2	404	40177	- -	(vergere et ar.	max. energy	I.U KEV	
2	2 62	2 6000	2.2	11			
5	2.03	2.0909	-2.5				
4	0.00	9.0314	-1.9	**			
5	21.4	22.046	-2.9				
6	42.9	44.0/3	-2./				
7	76.0	76.990	-1.3				
8	119.	121.92	-2.4	11			
9	177.	178.17	7	11			
10	257.	248.20	+3.5	11			
11	24.5	24.900	-1.6	11			
12	36.5	37.146	-1.7	11			
13	51.7	52.999	-2.5	11			
14	71.0	73.101	-2.9	11			
15	96.5	98.248	-1.8	11			
16	127.	129.19	-1.7	11			
17	163	166 55	-2 1	11			
18	210	211 01	_ 5	11			
10	250	211.01	5	11			
20	237.	203.22	-1.0	50 1	52 024	1 0 %	
20	515.	323.07	-2.1	J2.1	55.054	-1.0%	
21	382.	- 390.72	-2.2	63.8	64.767	-1.5	
22	452.	466.34	-3.1	76.7	78.151	-1.9	
23	530.	549.09	-3.5	91.3	93.275	-2.1	
24	629.	638.99	-1.6	108.	109.77	-1.6	
25	715.	737.89	-3.1	126.	129.30	-2.6	
26	821.	842.08	-2.5	147.	150.45	-2.3	
27	931.	958.20	-2.8	170.	173.71	-2.1	
28	924	959.94	-3.7	195.	199.25	-2.1	
29	1110.	1114.8	4	226.	226.79	4	
30	166.	168 07	-1 2	256	257 33	5	
	2001	200107		2301	207100		
31	190.	195.91	-3.0	289.	290.62	6	
32	222.	227.55	-2.4	325.	326.35	4	
33	258.	263.23	-2.0	365.	364.8	+ .3	
34	299.	303.12	-1.4	409.	405.59	+ .8	
35	342.	347.37	-1.6	455.	451.41	+ .8	
36	398.	396.30	+ .4	(Veigele et al.	max. energy	1.6 keV)	
37	444	449.60	-1.3	11			
38	499	507.44	-1 7	11			
39	562	569 52	_1 3	107	108.75	-1.6	
40	627	636 72	_1.5	120	122 06	_1 7	
-10	027.	000.12	-T.J	140.	122.00		

	1.0 keV			2.0 keV			
Z	Veigele et al.	Scofield	% diff.	Veigele et al.	Scofield	% diff.	
	τ (kb)	τ (kb)		τ (kb)	τ (kb)		
41	707.	708.56	2%	135.	136.28	9%	
42	778.	786.27	-1.1	150.	151.86	-1.2	
43	847.	869.65	-2.6	165.	168.71	-2.2	
44	948.	958.39	-1.1	185.	186.81	-1.0	
45	1040.	1052.9	-1.2	205.	206.31	6	
46	1160.	115.4	+ .5	227.	227.10	0	
47	1250.	1259.4	8	249.	249.61	2	
48	1370.	1370.6	0	273.	273.72	3	
49	1480.	1487.4	5	299.	299.50	2	
50	1600.	1606.4	4	326.	326.83	3	
51	1720.	1733.5	8	355.	355.83	2	
52	1740.	1861.6	-6.5	387.	386.71	0	
53	1880.	1915.0	-1.8	420.	419.20	+ .2	
54	2040.	2050.4	5	459.	453.49	+1.2	
55	1740.	2065.0	-15.7	496.	489.50	+1.3	
56	1860.	1946.1	-4.4	533.	527.02	+1.1	
57	1960.	2094.1	-6.4	567.	566.65	0	
58	2200.	2257.3	-2.5	620.	604.84	+2.5	
59	2370.	2474.0	-4.2	663.	645.73	+2.7	
60	2080.	1585.1	+31.2	704.	687.44	+2.4	
61	501.	492.46	+1.7	750.	731.37	+2.6	
62	537.	523.56	+2.6	801.	776.95	+3.1	
63	570.	556.64	+2.4	850.	825.04	+3.0	
64	600.	595.74	+ .7	892.	875.04	+1.9	
65	655.	629.70	+4.0	963.	923.16	+4.3	
66	702.	670.17	+4.8	952.	932.95	+2.0	
67	750.	/13.61	+5.1	1010.	980.77	+3.0	
68	808.	760.23	+6.3	1070.	1037.0	+3.2	
69	871.	810.25	+7.5	1130.	1031.3	+9.6	
70	931.	863.88	+7.7	(932.)	1088.1	-14.4	
71	977.	922.72	+5.9	(978.)	999.98	-2.2	
72	1050.	985.30	+6.6	(1040.)	1063.4	-2.2	
73	1120.	1051.3	+6.5	(Veigele et al.	max. energy	1./ keV	
74	1190.	1120.9	+6.2	11			
/5	1260.	1193.6	+5.6	/ 17 4 3		1.0.1	
/6	1330.	1269.6	+4.8	(Veigele et al.	max. energy	1.9 KeV	
//	1430.	1350.5	+5.9				
78	1530.	1432.3	+6.8	200	260 10	15.0	
/9	1630.	160/ 0	+/.4	390.	308.12	+5.9	
80	1750.	1604.9	+9.0	41 5 .	390.53	+0.3	

	1.0 keV			2.0 keV		
Z	Veigele et al.	Scofield	% diff.	Veigele et al.	Scofield	% diff.
	T (kb)	τ (kb)		T (KD)	T (KD)	
81	1820.	1695.5	+7.3%	439.	413.96	+6.0%
82	1920.	1788.4	+7.4	466.	438.35	+6.3
83	2010.	1883.7	+6.7	492.	463.61	+6.1
84	2120.	1981.9	+7.0	524.	489.76	+7.0
85	2130.	2041.9	+4.3	555.	516.93	+7.4
86	2250.	2143.0	+5.0	588.	545.01	+7.9
87	2360.	2248.1	+5.0	621.	573.97	+8.2
88	2470.	2322.7	+6.3	654.	603.79	+8.3
89	2520.	2433.8	+3.5	681.	634.72	+7.3
90	2300.	2543.5	-9.6	717.	666.43	+7.6
91	2440.	2660.6	-8.3	761.	698.71	+8.9
92	2530.	2613.8	-3.2	798.	732.29	+9.0
93	2650.	2730.7	-3.0	839.	767.22	+9.4
94	2830.	2850.5	7	889.	803.28	+10.7

Figure 1.

Total measured cross sections (σ_{TOT} , circles), and theoretical photoeffect (τ), coherent scattering (σ_{COH}), incoherent scattering (σ_{INCOH}), and pair production (κ) cross sections for copper for photon energies 10 eV to 100 GeV. The total photonuclear cross section ($\sigma_{PH.N.}$) measured by Wyckoff et al. [73] is also shown.

Figures 2 - 95.

Comparisons of theoretical photoeffect data (curves) of Veigele et al., [8], [10] extending below 1 keV and Scofield [11] 1 keV to 1.5 MeV with experimental data (points) for elements H (Z = 1) to Pu (Z = 94). The data₂ (in units of barns/atom) in each case have been multiplied by $[E(keV)]^2$. The number of vertical-scale graph cycles was in each case determined by the Scofield data. For elements Z > 34 this results in an apparent low-energy cut-off in the Veigele et al. data which in reference [8] and [19] extends down to 0.1 keV for all elements Z=1 to 94.





















































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