## NBS TECHNICAL Note 627

Computation of Spectral Data for a Josephson Junction Circuit
U.S.

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# Computation of Spectral Data for a Josephson Junction Circuit 

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A computer program has been developed to study power flow between different frequency channels in a Josephson junction circuit. This paper discusses the mathematical assumptions used to get such results. They are the trapezoidal approximation from spline theory and the use of a finite range of frequencies to characterize the frequency spectrum. This paper describes the program and provides the FORTRAN listing, flow charts, and discusses how to use the program. A discussion of possible sources of errors is also included.

Key Words: Differential equation; fast Fourier transform; Josephson junction; nonlinear integral-differential equation; spline theory.

## 1. INTRODUCTION

The central problem of interest is studying the power flow between various frequency channels in a Josephson junction circuit that is appropriate for harmonic mixing. To be realistic, the circuit must include shunt capacitance and shunt resistance as well as source and load impedances all of which may be frequency dependent. The general circuit we have chosen to investigate is shown in figure 1.

If the impedance $Z_{i}$, of figure $]$, are not frequency dependent, it is relatively easy to solve the differential equation for the circuit and this solution is included as an appendix. The result is what we term the transient program, i.e., it gives the behavior of the circuit as a function of time starting from specified initial conditions. After a few oscillations, the transient behavior reaches a good approximation to a steady state oscillation and thus provides an estimate of a consistent set of currents, of currents, voltages, and phases which we use in the main program, the steady state program.

The steady state program imposes strict periodicity on the solution as required for Fourier harmonic analysis. Steady state solutions are generated in two steps, first with the $Z_{i}$ frequency independent and second, using the full integro-differential equation for general $Z_{i}$. In the notation of figure 1 , we want the steady state power for each frequency present in the detector, $Z_{3}$, and in the Josephson junction when we are given the applied signal $V_{s}$. The $V_{s}$ is specialized to be

$$
\begin{equation*}
V_{s}=V_{o}\left[b_{o}+b_{L} \cos \left(N_{L} s+\theta_{L}\right)+b_{U} \cos \left(N_{U} s+\theta_{U}\right)\right] \tag{1}
\end{equation*}
$$

Here $\mathrm{V}_{\mathrm{o}}$ is the voltage scale factor, $\mathrm{b}_{\mathrm{o}}$ is the dimensionless d . c . voltage, $\mathrm{b}_{\mathrm{L}}$ is the lower applied r.f. amplitude, $\theta_{L}$ is its phase, and $b_{U}, \theta_{U}$ are the same quantities for the upper applied r.f. . The integers, "frequencies" $N_{L}$ and $N_{U}$ generate multiples of the phase $s, o s s=\omega_{o} t \leq 2 \pi$ where $\frac{2 \pi}{\omega_{0}}$ is the base period of the system. Since our strategy is to perform a discrete Fourier analysis, obviously we must adjust the discrete spectrum to be consistent with the applied frequencies $N_{L}$ and $N_{U}$, the rotation frequency $N_{1}$ of the self oscillation $\omega=\frac{2 e V}{\hbar}=N_{1} \omega_{0}$ and all possible linear combinations of these. The $b_{o}$ is adjusted to allow only periodic solutions with each frequency related to this base frequency, $w_{o}$, by an integer. Note that $b_{o}$ is a function of $b_{U}, b_{L}, \theta_{L}, \theta_{U}$, and $N_{1}$, as well as the impedances $Z_{i}$.

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This report is subdivided into six additional sections. (2) The scale adjustment is defined and the basic equation system is determined. (3) The approximations necessary to solve the basic equation system are given. (4) The FORTRAN Program and flow chart used to solve the equation system are described. (5) We present a sample collection of output from the computer showing the following for each case: (a) the time plot of $i_{4}$, (b) the magnitude of the frequency spectra for $i_{2}$, (c) the magnitude of the frequency spectra for $V_{\text {out }}$, ( $d$ ) the relative phase of ( $i_{2}-V_{\text {out }}$ ) at eachfrequency, (e) the amplitude spectrum of $\frac{d \varphi}{d s}$, ( $f$ ) the steady state power into the detector $\mathrm{Z}_{3},(\mathrm{~g})$ the steady state power across the junction device, and finally (h) a summary chart showing such things as the scales of amperes, volts, resistance, frequency, the $b_{o}, b_{L}, b_{U}, \theta_{U}, \theta_{L}$ and the power and other items at the following selected frequencies, namely d. c., $N_{L}$, $N_{U}$, the direct beat frequency $N_{B}$, and the rotation frequency $N_{1}$. (6) The summary gives general conclusions and recommendations. (7) Finally, we have the appendix describing the transient program.

## 2. THE BASIC EQUATION SYSTEM

### 2.1 Definitions

The following definitions are noted:
$Z_{i}$ are the frequency dependent impedances for the circuit.
$R$ is the shunt resistance of the Josephson Junction.
$Z_{2}$ represents any residual impedance of the Junction device.
$C$ is the capacitance of the Josephson Junction.
$\varphi$ is the Josephson phase function.
$J$ is the Josephson critical current.
$\mathrm{V}_{\mathrm{o}}$ is the voltage scale.
$\mathrm{W} \phi,\left(\omega_{o}\right)^{*}$ is the radian frequency scale.
$R_{o}$ is the resistance scale.
$A_{0}$ is the current scale.
$P_{0}$ is the power scale.
K is the impedance operator.
$K_{o}$ is the d. c. part of the impedance operator.
$D$ is the integral equation operator in the circuit system.
$U$ is the net $d$. current needed to drive the system at the steady state rotation frequency $N_{1}$.
a is the characteristic radian frequency squared for critical current in dimensionless form.

ANU $(v)$ is the characteristic frequency for loss due to RC type effects.
ANUP ( $\nu_{p}$ ) is the characteristic frequency for loss due to the rest of the circuit.

ANUT ( $\nu_{T}$ ) is the total loss characteristic frequency.
$\mathrm{V}, \mathrm{VA}\left(\mathrm{v}, \mathrm{v}_{\mathrm{a}}\right)$ are the various voltages in dimensionless form.

* (Note that the first symbol is the FORTRAN one and the second is the algebraic form. This is only shown for the less obvious cases.)

PHA ( $\varphi_{a}$ ) is the phase with the rotation and directly induced r. f. signal removed.

PHB ( $\varphi_{b}$ ) is the directly induced r.f. signal phase plus the rotation and the initial phase at $s=0$.

PHC ( $\varphi_{c}$ ) is the initial phase at $s=0$.
THL, THU $\left(\theta_{L}, \theta_{U}\right)$ is the phase of the lower, upper r.f. frequency.

## 2. 2. Equation System

We define

$$
\begin{gather*}
S=w_{0} t, V_{o} \equiv \frac{w_{0}}{2 e}, A_{0}=V_{0} / R_{o}, Z_{i}=z_{i} R_{o}, J=A_{0} a,  \tag{2}\\
C=c_{1} / R_{o} w_{0}, R=R_{0} r_{1}, i_{\ell}=A_{0} I_{\ell}(\ell=1 \text { to } 6), \text { and } w_{0}=2 \pi / T .
\end{gather*}
$$

Note that $a, R_{o}$, and $\omega_{o}$ are fixed by requiring $c_{1}$ to be unity in the dimensionless system, by requiring ar ${ }_{1}=\frac{10}{a_{1}}$, and by using a ${ }_{1}$ near 1 . This allows a good initial guess for starting slope and d. current in the program. Near the hysteresis region ${ }^{l} a_{l}$ would have to be carefully adjusted to allow study of this region.

Further definitions are:

$$
V_{J}=\frac{\hbar}{2 e} \frac{d \varphi}{d t}=V_{0} \frac{d \varphi}{d s}=V_{0} V, i_{4}=A_{0} a \sin \varphi
$$

$$
\begin{equation*}
i_{5}=C \frac{d V_{2}}{d t}=V_{0} C \frac{d V_{2}}{d s} w_{0}=A_{0} \frac{d V}{d s}, \text { and } i_{6}=\frac{V_{L}}{R}=\frac{A_{o}}{r_{1}} \frac{d \varphi}{d s} \tag{3}
\end{equation*}
$$

Therefore

$$
\begin{gather*}
i_{1}=i_{4}+i_{5}+i_{6}=A_{o}\left[a \sin \varphi+\frac{V}{r_{1}}+\frac{d V}{d s}\right] .  \tag{4}\\
V_{s} \equiv V_{o} V_{1}, i=i_{1}+i_{2}, i_{2}=A_{o} I_{2}, i_{1}=A_{o} I_{1},  \tag{5}\\
i=A_{o} I_{o}, V_{o u t}=V_{o} V_{2} .
\end{gather*}
$$

Therefore we have the dimensionless operator equations:

$$
\begin{align*}
& I_{2}=z_{3}^{-1}\left[z_{2} I_{1}+V\right] \\
& I_{1}=a \sin \varphi+V / r_{1}+\frac{d V}{d s}  \tag{6}\\
& I_{0}=\left(I+z_{3}^{-1} z_{2}\right) I_{1}+z_{3}^{-1} V \\
& V_{1}=z_{1} I_{0}+V_{2} \text {, and } \\
& V_{2}=z_{3} I_{2}=z_{2} I_{1}+V
\end{align*}
$$

These are reduced to a single operator form.

$$
\begin{align*}
& V_{1}=z_{1}\left(1+z_{3}^{-1} z_{2}\right) I_{1}+z_{1} z_{3}^{-1} V+z_{2} I_{1}+V \text { or } \\
& V_{1}=\left[z_{1}+z_{2}+z_{1} z_{3}^{-1} z_{2}\right]\left[a \sin \varphi+\frac{d V}{d s}\right]  \tag{7}\\
& +\left\{\left[z_{1}+z_{2}+z_{1} z_{3}^{-1} z_{2}\right] / r_{1}+z_{1} z_{3}^{-1}+1\right\} V
\end{align*}
$$

We finally define
$K \equiv z_{1}+z_{2}+z_{1} z_{3}{ }^{-1} z_{2}, \nu \equiv 1 / r_{1}, D \equiv K^{-1}\left(z_{1} z_{3}{ }^{-1}+1\right)-\frac{1}{K_{o}}\left(z_{1}{ }^{\circ}\left(z_{3}{ }^{\circ}\right)^{-1}+1\right)$, and $\nu_{P} \equiv \frac{1}{K_{o}}\left(z_{1}{ }^{o}\left(z_{3}\right)^{0}{ }^{-1}+1\right)$
as the d. c. part of operator, where $z_{j}^{\circ}$ is the D. C. value of $z_{j}$, and

$$
v_{\mathrm{T}}=v+v_{\mathrm{P}}
$$

to get the following pair of operator equations

$$
\begin{gather*}
K^{-1}\left(V_{1}\right)=a \sin \varphi+\frac{d V}{d s}+V_{T} V+D V, \text { and } \\
V=\frac{d \varphi}{d s} \tag{9}
\end{gather*}
$$

As the reader can see, if the $z_{i}$ are frequency independent then the $D$ operator is zero and we have a pure differential equation. If, as would usually be the case, there is frequency dependence, then $D$ is not zero. However, if we put in estimates for the values of $z_{i}$, namely $z_{1} \sim 10^{4}, z_{2} \leqq 1$, and $z_{3} \sim 10^{8}$, we have $D \sim 10^{-4}$. If $a \gtrsim 1$, then our problem is dominated by the differential equation and the integral equation part can be used in an iteration procedure. The acid test is to run the program.

As we are interested only in strictly periodic solutions, we as sume the following forms for the $\varphi$ and V. First we $\operatorname{split} \varphi=\varphi_{a}+\varphi_{b} ; V=V_{a}+V_{b}$ and note that $\varphi_{b}=N_{l} s+\varphi_{c}+\varphi_{d}$ and $V_{b}=N_{l}+\frac{d \varphi}{d s} d \equiv \frac{d \varphi_{b}}{d s}$. We require $\varphi_{d}$ to be defined by

$$
\begin{equation*}
K^{-1}\left(V_{1} \text { A. C. }\right)=\frac{d^{2} \varphi d^{2}}{d^{2} s}+\nu_{T} \frac{d \varphi}{d s}+D \frac{d \varphi}{d s} \tag{10}
\end{equation*}
$$

and $\varphi_{a}$ to be defined by

$$
\begin{gather*}
u=\frac{d V_{a}}{d s}+a \sin \left(\varphi_{a}+\varphi_{b}\right)+\nu_{T} V_{a}+D V_{a}  \tag{11}\\
u \equiv K_{o}^{-1} b_{o}-\nu_{T} N_{1} .
\end{gather*}
$$

Here

The purpose of this split is to remove a ramp term and to get the first order induced phase function, so that the remaining part, $\varphi_{a} V a$ are periodic and have a chance for rapid convergence to the correct solution. Since $V_{a}, \varphi_{a}, \varphi_{d}$ are all periodic function of $s$, the exact solution would be periodic. We use the Fourier space form to get an idea of the structure of D. Thus with $-\infty \leq \ell \leq \infty$,

$$
\begin{gather*}
\left.\varphi_{a}(s)=\ell^{\Sigma} e^{-i \ell s} \tilde{\varphi}_{a}(\ell), V_{a}(s)=\ell^{\Sigma} e^{-i \ell s} \tilde{V}_{a}(\ell), \text { (note that } \tilde{V}_{a}(0)=0 .\right) \\
\varphi_{d}(s)=e^{-i N_{L}}{ }^{s} \tilde{\varphi}_{d}\left(N_{L}\right)+e^{i N_{L} s} \tilde{\varphi}_{d}^{*}\left(N_{L}\right)+e^{-i N_{U}^{s}} \tilde{\varphi}_{d}\left(N_{U}\right)+e^{i N_{U}^{s}} \tilde{\varphi}_{U}^{*}\left(N_{U}\right),  \tag{12}\\
\text { and } \left.D(s)=\Sigma e^{i \ell s} \tilde{D}(\ell) . \quad \text { (We have made } D V=\int_{0}^{2 \pi} \frac{d s^{\prime}}{2 \pi} D\left(s-s^{\prime}\right) V\left(s^{\prime}\right)_{0}\right)
\end{gather*}
$$

The equation (10) defines $\tilde{\varphi}_{d}$ as

$$
\begin{equation*}
K^{-1}\left(N_{L}\right) \frac{b_{L}}{2} e^{-i \theta} L=\left(-i N_{L}\right)+\nu_{T}+\tilde{D}\left(N_{L}\right) \quad\left(-i N_{L}\right) \tilde{\varphi}_{d}\left(N_{L}\right) \tag{13}
\end{equation*}
$$

The $\tilde{\varphi}_{d}\left(N_{U}\right)$ is gotten by $U$ replacing $L$ in the above expression. $\tilde{\varphi}^{*}$ is the complex conjugate. We note that $\mathrm{N}_{\mathrm{U}}$ and $\mathrm{N}_{\mathrm{L}}$ cannot be zero.

The equation (11) can be written as

$$
\begin{equation*}
u \delta_{0}^{\ell}=(-i l)^{2}+\nu_{T}(-i l)+\tilde{D}(l)(-i l) \varphi_{a}(l)+a \int_{0}^{2 \pi} \frac{d s}{2 \pi} e^{i l s} \sin \left(\varphi_{a}+\varphi_{b}\right) . \tag{14}
\end{equation*}
$$

where $\delta_{0}^{\ell}$ is the Kronecker delta. We further note that $\tilde{D}(\ell)$ is given as

$$
\tilde{D}(\ell)=\left(z_{1}(\ell) z_{3}^{-1}(\ell)+1\right) / K(\ell)-\left(z_{1}(0) z_{3}^{-1}(0)+1\right) / K(0)
$$

where

$$
\begin{gathered}
K(l)=z_{1}(\ell)+z_{2}(\ell)+z_{1}(\ell) z_{2}(\ell) z_{3}(\ell)^{-1} \text { and } \\
K(0)=K_{0} .
\end{gathered}
$$

It appears to be best to do the calculation in the time domain and then do the spectral transform. Thus we have
and

$$
\begin{gathered}
\varphi_{d}(s)=2 \operatorname{Real}\left[\tilde{\varphi}_{d}\left(N_{L}\right) e^{-i N_{L} s}+\tilde{\varphi}_{d}\left(N_{U}\right) e^{-i N_{U} s}\right] \\
\varphi_{b}(s)=\varphi_{d}(s)+N_{l} s+\varphi_{c} \\
u=\frac{d}{d s} V_{a}+a \sin \left(\varphi_{a}+\varphi_{b}(s)\right)+\nu_{T} V_{a}+\int_{0}^{2 \pi} \frac{d s^{\prime}}{2 \pi} D\left(s-s^{\prime}\right) V_{a}\left(s^{\prime}\right), \\
V_{a}=\frac{d}{d s} \varphi_{a}
\end{gathered}
$$

as our basic equation system. To progress further, we must make numerical approximations.

We have to choose a discrete set of s-points, $N_{2}$, because we must take a finite amount of computation time and must keep the entire data set in the machine to allow fast fourier transformation. The set of points has to be equidistant and equal to $2^{N_{0}}$, where $N_{0}$ is some integer. Practical considerations restrict the range of $N_{o}$ from 8 to 11 .

## 3. 1 Constraints

A discrete set of points implies frequency folding, hence we require the following conditions to minimize this effect.
(a) The highest frequency of interest must be $\leq \frac{\mathrm{N}_{2}}{8}$. All frequencies above this value are to be treated as unreliable, due to the method of analysis. Computational noise must also be small.
(b) Our demand accuracy for analysis is near 5\%. We make checks on the results by increasing $\mathrm{N}_{2}$ as the occasion demands. If the spectrum does not change to the demand accuracy, our analysis is treated as adequate and accurate to that order. Incidently, if the value of $U$ is correct to $5 \%$ or less, then the lower frequency spectra is also correct to that accuracy.
(c) This $5 \%$ computational accuracy level suggests using $N_{2}=256$ or 512 with an occasional extension to 2048 points for special studies.
(d) We allow $z_{i}(\ell)$ to have arbitrary values for $0 \leq \ell \leq \frac{N_{2}}{2}$ and require the remaining $\ell$ values to be fixed by $z_{i}{ }^{*}(\ell)=z_{i}\left(N_{2}-\ell\right)$.
(e) In $\tilde{D}(\ell)$ and $\tilde{\varphi}_{d}(\ell)$, we use in place of $(-i \ell)$ the periodic expression $\left(e^{-i h \ell}-1\right) / h$, where $h \equiv 2 \pi / N_{2}$ is the interval between points in the time space. This allows folding and preservation of reality conditions. The sensitivity to this change is checked by the doubling of points.
(f) Because $D\left(s-s^{\prime}\right)$ is a small correction to the differential equation system, we replace, with no qualms, the integral by a sum using the trapezoidal rule. Thus

$$
\begin{equation*}
\int_{0}^{2 \pi} \frac{d s^{\prime}}{2 \pi} D\left(s-s^{\prime}\right) V_{a}\left(s^{\prime}\right)=\sum_{j^{\prime}=0}^{N_{2}-1} \frac{1}{N_{2}} D\left(j-j^{\prime}\right) V_{a}\left(j^{\prime}\right) \tag{16}
\end{equation*}
$$

(g) To complete the approximations, we must decide how to approximate $\frac{d V}{d s}$ type terms. Finite differencing does not allow high enough precision in the fit. Since the spectral information for power uses $\frac{d V}{d s}$ in its current, we need to have smoothness in the results to at least order $h^{2}$. Additionally, we want to avoid absolute instabilities and permit the application of quasilinear theory ${ }^{3}$ to get the $b_{o}$ value needed to allow strict periodicity. These requirements make it necessary to use the spline method ${ }^{2}$ using at least the trapezoidal rule. We assume the trapezoidal rule everywhere in the calculation. The results of this analysis should be smooth functions (accuracy order, $h^{2}$ ), hence a frequency separation of true results from noise due to the discrete analysis of the system and an inherent overall accuracy of $h^{2}$. Periodicity should tend to force uniform error throughout the time interval. Sharp spikes in the time dependence of $\varphi_{a}$ will negate the uniformity of the error and make higher frequencies spectral information less accurate.
( $h$ ) Now the problem is split into a part ( $\alpha$ ) which neglects $D$ and a part ( $\beta$ ) that takes $D$ into account.

Here we have the simplified system of equations using the trapezoidal rule.

$$
\begin{align*}
& \varphi_{a}(j+1)=\varphi_{a}(j)+\frac{h}{2}\left[V_{a}(j+1)+V_{a}(j)\right]  \tag{17}\\
& V_{a}(j+1)=V_{a}(j)+\frac{h}{2}\left[q_{1}(j+1)+q_{l}(j)\right]  \tag{18}\\
& q_{l}(j)=-a \sin \left[\varphi_{a}(j)+\varphi_{b}(j)\right]-v_{T} V_{a}(j)+u . \tag{19}
\end{align*}
$$

The boundary conditions are

$$
\mathrm{v}_{\mathrm{a}}(0)=\mathrm{V}_{\mathrm{a}}\left(\mathrm{~N}_{2}\right), \text { and } \varphi_{a}(0)=\varphi_{a}\left(\mathrm{~N}_{2}\right)
$$

The $u$ is adjusted to allow this result. The eq (19) can be rewritten with definitions
as

$$
\begin{gathered}
h_{1}=\left(2-h \nu_{T}\right) /\left(2+h \nu_{T}\right), h_{2}=-a h /\left(2+h \nu_{T}\right), h_{3}=2 h /\left(2+h \nu_{T}\right), \\
V_{a}(j+l)=h_{1} v_{a}(j)+h_{3} u+h_{2}[\sin \varphi(j+1)+\sin \varphi(j)]
\end{gathered}
$$

We generate two solutions starting with $\varphi_{a}(0)$ fixed and unchanged for a given calculation run. To avoid numerical noise problems, experience shows that it is absolutely necessary to fix the $\varphi_{a}(0)$, when $u$ is adjustable.

## 3. 3 Calculation Sequence

(1) We iterate for a given $j, j+1$ pair until $\left|\varphi_{a}(j+1)-\varphi_{a}(j)\right|$ of the $n+1$ iteration changes by
EPI $=10^{-8}$. less than EPI $=10^{-8}$.
(2) We compute two quasilinear solutions using the following starting conditions

$$
\delta \varphi_{a}^{(1)}(0)=0, \delta \mathrm{~V}_{\mathrm{a}}^{(1)}(0)=1, \text { and } \delta \varphi_{a}^{(2)}(0)=\delta \mathrm{V}_{\mathrm{a}}^{(2)}(0)=0
$$

In case one $\delta u=0$ and in case two $\delta u=1$. The two equation sets are, with $q_{2}(j) \equiv \cos \varphi(j)$,

$$
\begin{gather*}
\delta \varphi_{a}^{(i)}(j+1)=\delta \varphi_{a}^{(i)}(j)+\frac{h}{2}\left[\delta V_{a}^{(i)}(j+1)+\delta V_{a}^{(i)}(j)\right], \\
\text { and } \quad \delta V_{a}^{(i)}(j+1)=h_{l} \delta V_{a}^{(i)}(j)+h_{3} \delta{ }_{2}^{i}+h_{2}\left[q_{2}(j+1) \delta \varphi^{(i)}(j+1)+q_{2}^{\left.(j) \delta \varphi^{(i)}(j)\right] .}\right. \tag{20}
\end{gather*}
$$

These can be rewritten as

$$
\delta \varphi_{a}^{(i)}(j+1)=\delta \varphi_{a}^{(i)}(j)+\frac{h}{2}\left[\delta v_{a}^{(i)}(j+1)+\delta V_{a}^{(i)}(j)\right]
$$

$$
\delta \mathrm{V}_{\mathrm{a}}{ }^{(\mathrm{i})}(\mathrm{j}+1)=\mathrm{T}_{1}^{(\mathrm{i})} / \mathrm{T}_{2}^{(\mathrm{i})},
$$

with

$$
\begin{equation*}
T_{1}{ }^{(i)} \equiv\left[h_{1}+h_{2} \frac{h}{2} q_{2}(j+1)\right] \delta V_{a}{ }^{(i)}(j)+h_{3} \delta_{2}^{i}+h_{2} \delta \varphi{ }^{(i)}(j)\left[q_{2}(j+1)+q_{2}(j)\right], \tag{21}
\end{equation*}
$$

and

$$
T_{2}^{(i)} \equiv 1-h_{2} \frac{h}{2} q_{2}(j+1) .
$$

(3) We use the values computed at $\varphi_{a}\left(N_{2}\right), \delta \varphi^{(i)}\left(N_{2}\right)$, etc. to satisfy the following requirements and hence fix $\beta$ and $\delta u$ :
and

$$
\begin{gather*}
\varphi_{a}\left(N_{2}\right)-\varphi_{a}(0)+\beta \delta \varphi_{a}^{(1)}\left(N_{2}\right)+\delta u \delta \varphi_{a}^{(2)}\left(N_{2}\right)=0, \\
v_{a}\left(N_{2}\right)-v_{a}(0)+\beta\left(\delta v_{a}^{(1)}\left(N_{2}\right)-1\right)+\delta u \delta v_{a}^{(2)}\left(N_{2}\right)=0 . \tag{22}
\end{gather*}
$$

The new $\varphi_{a}(j), v_{a}(j)$, and $u$ are now defined by

$$
\begin{gather*}
\varphi_{a}(j)=\varphi_{a}(j)(\text { old })+\beta \delta \varphi_{a}^{(1)}(j)+\delta u \delta \varphi_{a}^{(2)}(j), \\
v_{a}(j)=v_{a}^{(j)(o l d)+\beta \delta v_{a}^{(1)}(j)+\delta u \delta v_{a}^{(2)}(j),}  \tag{23}\\
u=u(\text { old })+\delta u .
\end{gather*}
$$

and

In the computer program the corrections to the initial guesses of $u$ and $V_{a}(0)$ are bounded because initially the first guess may not be in the circle of convergence except when the computer program in the appendix is used. Thus, it is necessary to have built into the calculational procedure a means to search for the region where the quasilinear technique will work. The computer program has allowed jumps in $\delta u$ and in $\delta V_{a}(0)$ that shrink automatically as each computation is made with a new $u$ and $V_{a}(0)$. If after ten iterations the convergence has not taken place, then the $V_{a}(0)$ is set equal to $\mathrm{V}_{\mathrm{a}}\left(\mathrm{N}_{2}\right)$ using the idea that it is more likely to be compatible with a periodic solution because the transient should be less. Then we repeat the above quasilinear process and the window shrinking process for up to eight times. If the convergence is particularly difficult, the shift in $V_{a}(0)$ can take place up to ten times. Finally, a relaxed test for acceptance of the results as converged is made so that a spectral analysis can be made.

### 3.4 Problems

There are two main sources for convergence problems: (1) The number of computed points is too small when there are sharp spikes in the $V_{a}$ and $\frac{d V_{a}}{d s}$. This causes extensive computational noise. (2) The initial $V_{a}(0)$ and $u$ are badly guessed. Both problems are particularly bad when there are strong $B_{V}, B_{L}$ terms and/or when $a, \nu_{T}$ are very large numbers.

Here we use the full quasilinear theory. ${ }^{3}$ First, we define

$$
\begin{align*}
& q_{l}(j)=-a \sin \varphi(j)-\sum_{j^{\prime}=0}^{N_{2}-l} D\left(j-j^{0}\right) \frac{1}{N_{2}} v_{a}\left(j^{0}\right)+u-\nu_{T} V_{a}(j) \\
& q_{2}(j)=\cos \varphi(j) \\
& q_{3}(j)=h_{3}\left\{\frac{1}{h}\left(v_{a}(j+1)-v_{a}(j)\right)-0.5\left(q_{l}(j+1)+q_{l}(j)\right)\right\}  \tag{24}\\
& q_{4}(j)=h\left\{\frac{1}{h}\left(\varphi_{a}(j+1)-\varphi_{a}(j)\right)-0.5\left(v_{a}(j+l)+v_{a}(j)\right)\right\}
\end{align*}
$$

We note that $q_{3}, q_{4}$ are the errors due to not fitting the equation system to a quadratic spline. We use these terms to define the error signal in the quasilinear theory. In addition to the terms needed to correct the error signal, we must add effects due to needed changes in $V_{a}(0)$, and $u$ to allow periodic solutions. Since the quasilinear theory integrates over the entire domain, the $u$ is sensitive to the global solution, and periodicity tends to force global accuracy; we should retain a global accuracy of order $h^{2}$, and a smooth spectrum to that same order. This should be accurate enough to separate the true solution from the effects of the finite analysis. Again, spiking will make the higher frequencies less accurate.

The quasilinear differential equations are replaced by the trapezoidal rule. The basic equations become
and

$$
\begin{align*}
\delta \varphi_{a}^{(i)}(j+1) & =\delta \varphi_{a}^{(i)}(j)+\frac{h}{2}\left[\delta v_{a}^{(i)}(j)+\delta v_{a}^{(i)}(j+1)\right]-q_{4}(j+1) \delta \\
\delta V_{a}^{(i)}(j+1)= & h_{1} \delta V_{a}^{(i)}(j)+h_{2}\left[q_{2}(j) \delta \varphi_{a}^{(i)}(j)+q_{2}(j+1) \delta \varphi_{a}^{(i)}(j+1)\right]  \tag{25}\\
& +h_{3} \delta_{2}^{i}-q_{3}(j+1) \delta_{1}^{i}, \text { for } i=1,2,3 .
\end{align*}
$$

The initial conditions are $\delta \varphi_{a}^{(i)}(0)=0$, for all $i$; the $\delta V_{a}^{(i)}=0$ for $i=1,2$, and the $\delta V_{a}^{(3)}(0)=1$. The periodic conditions at $j=N_{2}$ requires

$$
\begin{gather*}
\delta \varphi{ }^{(1)}\left(\mathrm{N}_{2}\right)+\delta u \delta \varphi{ }^{(2)}\left(\mathrm{N}_{2}\right)+\beta \delta \varphi^{(3)}\left(\mathrm{N}_{2}\right)=0, \\
\delta \mathrm{~V}^{(\mathrm{I})}\left(\mathrm{N}_{2}\right)+\delta u \delta \mathrm{~V}^{(2)}\left(\mathrm{N}_{2}\right)+\beta\left(\delta \mathrm{V}^{(3)}\left(\mathrm{N}_{2}\right)-1\right)=0 . \tag{26}
\end{gather*}
$$

We use these equations to fix $\delta u$ and $\beta$. We compute for the different $\alpha$, a new $\varphi_{a} \equiv$ old $\varphi_{a}+\alpha\left(\delta^{(1)} \varphi_{a}+\delta u \delta^{(2)} \varphi_{a}+\beta \delta^{(3)} \varphi_{a}\right.$ ), and a similarily constructed new $V_{a}$ which are then inserted to compute a new $q_{3}$ and $q_{4}$ set for each $\alpha$. The sum of these; namely, $P(\alpha)=\sum_{j=0}^{N_{2}-1}\left(q_{3}{ }^{2}+q_{4}{ }^{2}\right.$ ), gives us an error signal. Our choice for the $\operatorname{cor} r$ rect $\varphi_{a}, V_{a}$ is fixed by the requirement the correct $\alpha$ has; $P(\alpha)<P(0)$. The smallest $P(\varphi)$ improves the calculation the most and guarantees convergence.

The $\alpha$ chosen, is then used to construct the new $\varphi_{a}$ and $V_{a}$ and the entire part ( $\beta$ ) quasilinear analysis is repeated. This process is continued until 10 iterations have been performed or until $P(o)<2 h^{2}$. The results of this process constitutes the best accuracy that can be obtained for a global solution subject to the inherent errors of the trapezoidal method.

Next we compute by fast Fourier transform the desired spectra. This completes the discussion of the assumptions made to get numerical results. The next section has the FORTRAN program listing and a flow chart of the basic program flow. In particular, we show the input data requests.

## 4. THE FORTRAN PROGRAM AND FLOW CHART

Here we have listed the Fortran programs that give the main details of the computation. The microfilm and fast Fourier transform programs are not listed. The former depends on the computer installation and the latter should be available and optimized at each computer installation. After the Fortran list, figure 2, a flow chart, figure 3, shows the bulk data and computational flows between programs.

## 5. THE SAMPLE COLLECTION OF OUTPUT

This section discussed the printed and the microfilm output for a single computed case. The purpose is to give general information to allow use of the Fortran program by an interested party. No attempt is made to report all the possible data that exists as a consequence of the computations. Selected information generated as a consequence of this program will be described in separate reports. To use the program correctly in detail, one must read the Fortran listing.

## 5. 1 The Printed Output

Figure 4 shows the printed output from a single computed case. The first row of data gives the number of data points used, the number (11) corresponds to 2048 points in time and 1024 different frequencies. The $J(1,-4)$ is the supercurrent in amperes. The Al (1.25-0) gives the scale adjustment defined on page 4. $C(8.2277-15)$ is the capacitance in farads. $R(1 . \pm 1)$ is the number of ohms. EBl (1. - 8) is the demand accuracy for each step interation and scales the overall demand accuracies in the various iterations of part ( $\alpha$ ) and part ( $\beta$ ). The EP6 (1. -0) can be set to increase demand accuracy of part ( $\beta$ ). EP6 equal 1. uses the values stated on page 10 for the converged exit. EP6 less than one increases the accuracy demand and EP6 greater than one decreases accuracy demand.

The first Z 1 , etc., row gives, with 19 equal 1 , the complex values of $\mathrm{Z} 1, \mathrm{Z} 2$, and Z 3 for all impedances as if they were frequency independent. Subsequent rows of Zl type data show specific frequency exceptions to the frequency independence assumption. These are overlaid on the previous values. Please note that $I 9$ is one unit larger than the corresponding frequency number, thus 1 is the d. c. case. In our sample output we have changed the first harmonic of the Z 3 impedance from ( $1 .+8,0$. ) to ( $1 .+5,0$ ). The others have remained unchanged.

We continue the row by row analysis of the printed output. After the impedance spectrum is defined, we next specify the lower frequency, $N L(4)$, the upper frequency, $N U(5)$, at which the two r.f.

|  | PROGRAM JOSPHC |
| :---: | :---: |
| $\begin{aligned} & \mathrm{C} \\ & \mathrm{C} \end{aligned}$ | THIS PROGRAM COMPUTES FREQUENCY SPECTRA FOR A CIRCUIT MODELING THE POINT CONTACT JOSEPHSON JUNCTION. THE ANALYSIS USES A TIME |
| c | PERIOD WITH N2 POINTS. NO IS THE ORDER 2**NO=N2 NO.LE. 11 |
|  | DATA(PI = 3.14159265359), (IM=(0., 1.) ),(PHS=2.06778E-15) DIMENSION Z1(1025),Z2(1025),Z3(1025) |
|  | COMMON /5/ BI,S1,52,S3 |
|  | 1, A1 |
|  | COMMON /16/ V16,P16, H 8 |
|  | TYPE COMPLEX $\mathrm{Zl}, \mathrm{Z2,Z3,K,DT,ZF,ZH,ZG,KL,AN2,ATMP,KU,ATN,AN6,AN7,ZI}$ |
|  | 1,CBU,CBL , ZD |
|  | TYPE REAL KO |
|  | COMMON /DAT/ PI,IM,PHS,U |
|  | DIMENSION K(1025), DT(1025), KL(3), KU(3) |
|  | 1 , I2T(1025),VAT(1025) |
|  | EQUIVALENCE (K(1),DPH(1)),(DT(1),DVA(1)) |
|  | COMMON /1/ PHC,VA(2050),DPH(4,2050), |
|  | lDVA (4, 2050), H, H1, H2, H3, H4,EP1,EP2,EP3,LBX(5),LBTR(5),EP4(6) |
|  | EQUIVALENCE (DPH(6151),Z1(1)),(DVA(6151),Z2(1)),(Q4(1),23(1)) |
|  | COMMON /2/D(2050), PHA (2050), PHB(2050), Q1(2050),Q2(2050), Q3 (2050), |
|  |  |
|  | $2, N B, K L, K U, A N 5, H 5, ~ A N U, A N U P, A N U T, A, R 1, R 0, C, W 0, T, ~ R, P I 2, A N 2, A N 4$, |
|  | 3AN3,MM,N5,N6,N7,ISKPD,SC, OUT $6,111, N R, ~ I C O N V, A J, I C A S E, ~$ |
|  | 4POW1, POW2, POWJ, POW IN, K0, EP5, H 7 |
|  | 5,CBU,CBL |
|  | EQUIVALENCE $\mathrm{X}(1)$, DPH(1)),(Y(1), DPH(2051)),(VAT(1),DVA(1)), |
|  | 1(I1T(1), DVA(2051)), (VOUT(1), DVA(4101)),(X1(1), DPH(1025)) |
|  | TYPE COMPLEX VAT,IIT,VOUT,I2T, IM |
|  | DIMENSION $\mathrm{X}(1025), \mathrm{Y}(1025), 11 \mathrm{~T}(1025)$, VOUT 11025$), \mathrm{X} 1(1025), \mathrm{Y} 1(1025)$ |
|  | EQUIVALENCE(Y1(1), DPH(4101)), (12T(1),Q3(1)) |
| 1 | READ 2,N0,AJ,Al, ${ }^{\text {, R,R,EP1,EP6 }}$ |
| 2 | FORMAT(I5,6E12.4) |
|  | PRINT 200,N0,AJ,A1, C,R,EP1,EP6 |
| 200 | FORMAT(* NO,J,A1,C,R,EP1,EP6 */I5,6E12.4) |
|  |  |
|  | N6 $=$ N2+2 \$N7=N2-1 \$AN2=-IM*AN4 \$MM=NO-1 \$SC=.5/AN3 \$NP=NH+1 |
|  | XT=PI2*C \$ R7=PHS/(XT*AJ) \$ $A=R 7 /\left(R^{*} A 1 * * 1\right) * * 2$ |
|  | $R 0=S Q R T F(A * R 7) ~ \$ T=R O * X T ~ \$ W 0=P I 2 / T \quad$ SNR=NH/2 |
|  | $R 1=R / R O$ S $A N U=1 . / R 1 \quad \$ V 0=P H S / T$ SAO $=V 0 / R 0$ \$PO=VO*AO \$AN8=NH |
|  | EP2=EP1*A SEP5=.1*A |
|  | Sl=R1*SQRTF(A) \$S2=A1/10. \$ S3=EXPF(.2-S1) |
| C | SET I8=Igel First Time to Generate all the same zi, $22, Z 3$. |
| C | TO ALLOW SELECTED FREQUENCY DIFFERENCE USE I $8=0$ I 9 THE PARTICULAR |
| C | HARMONIC. TO EXIT USE I8=-1 AND I9 THE LAST PARTICULAR FREQUENCY |
| C | HERE I9=NP AND $\mathrm{Zl} 1, \mathrm{Z2,Z3}$ ARE REAL AS MUST BE Z111) ETC. |
| C | I9=1 IS D.C. CASE ADD 1 TO HARMONIC ORDER TO GET CORRECT |
| C | LOCATION IN DATA STORE. |
| 5 | READ 3,I8,I9,Z1(19),Z2(19), Z3(19)\$PRINT199,19,Z1(19),Z2(19),Z3(19) |
| 3 | FORMAT(I3,I5,3C(E12.3,E12.3)) |
|  | IF (I8)4,5,6 |
| 6 | DO $7 \mathrm{I}=2, \mathrm{NP}$ \$Z1(I)=Z1(1) \$Z2(I)=Z2(1) \$Z3(I)=Z3(1) |
| 7 | CONTINUE \$GO TO 5 |
| 4 | DO $8 \mathrm{I}=1, \mathrm{NP}$ \$Z1(I)=Z1(I)/R0 \$Z2(I)=Z2(I)/R0 \$Z3(I)=Z3(I)/R0 |
|  | ATMP=1.+Z1(I)/Z3(I) SK(I)=Z1(I)+Z2(I)*ATMP \$DT(I)=ATMP/K(I) |
| 8 | CONTINUE \$ANUP=DT(1) \$ANUT=ANUP+ANU \$DO $9 \mathrm{I}=2$,NP |
| 9 | DT(1)=DT(I)-DT(1) \$KO=K(1) \$ DT(1)=0. |
|  | READ 10,NL,NU, ISKPD, ICASE |
| C | NL IS LOWER APPLIED FREQUENCY NU IS UPPER |
| 10 | FORMAT (4I5) |
| C | ISKPD $=+1$ DO THE D STRUCTURES. ISKPD=0 SKIP D STRUCTURES HERE D=0. |
|  | PRINT 201,NL,NU,ISKPD, ICASE |
| 201 | FORMAT (* NL, NU, ISKPD,ICASE *,415) |
| C | ICASE IS INCREMENTED TO LABEL EACH COMPUTED CASE LOAD LAST VALUE |
|  | IF (NL-NU) I1, 12,12 |
| 12 | PRINT 202,NL,NU SCALL EXIT |
| 202 | FORMAT ( ${ }^{\text {* }}$ NL TOO LARGE * 2I5) |

Figure 2. The steady state Fortran program listing.

| 11 | IF (NU-NR) $13,14,14$ |
| :---: | :---: |
| 14 | PRINT 203,NU,NR |
| 203 | FORMAT ( NU TOO LARGE * 2I5) |
| 13 | $N B=N U-(N U / N L) * N L$ <br> $\operatorname{KEY}(1)=0 \quad \operatorname{SKEY}(2)=\operatorname{NB} \operatorname{SKEY}(3)=\operatorname{NL} \operatorname{SKEY}(4)=\operatorname{NU} \quad \operatorname{SKEY}(5)=\mathrm{NH}$ |
|  | IF (ISKPD) 15,19 |
| 15 | DT(NP) = - 5*DT(NP) \$ZH=CEXP(AN2) \$ZG=CONJG(ZH)\$DO $16 \mathrm{~J}=1, N 2$ SD(J) $=0$ 。 |
|  | ZG=ZG*ZH \$ZF=1. \$DO 17 I=2,NP \$ZF=2F*ZG \$ZD=OT(I)*ZF |
| 17 | $D(J)=D(J)+R E A L(Z D)$ |
| 16 |  |
|  | READ 18,IDRED \$GO TO (19,20),IDRED |
| C | DECIDE PRINT OF D IDRED $=2$ YES $=1$ NO |
| 20 | PRINT 204,(D(J), J=1,N2) |
| 204 | FORMAT ( ${ }^{\text {* }}$ S SET * /(10E12.3)) |
| 19 |  |
|  | $\mathrm{Hl}=(2 .-\mathrm{HT}) * \mathrm{HU} \quad \$ \mathrm{H} 5=\mathrm{H} 4 * \mathrm{H} 2$ \$ $\mathrm{H} 7=\mathrm{HU}$ |
|  | AN7 = AN2*NU \$ATMP = (CEXP(AN7)-1.)/H \$ATN=(ATMP+ANUT+DT(NU+1)) |
|  | $K U=0.5 /(A T M P * K(N U+1) * A T N) \quad \$ K U(2)=A T M P * K U ~ \$ K U(3)=A T M P * K U(2)$ |
|  | AN6 = AN2*NL \$ATMP=(CEXP(AN6)-1.)/H \$ATN=(ATMP+ANUT+DT(NL+1)) |
|  |  |
| C | N1 IS RAMP FREQUENCY RELATIVE TO BASE USE PHC=-PI IS NI NOT ZERO |
| C | PHO =PHC AT S=O. BL, THL IS LOWER APPLIED RELATIVE VOLT AND PHASE |
| C | BU, THU ARE THE UPPER |
|  | CALL CHARTA |
| C | NOTE BL=KO*BL USED IN TRANSIENT PROGRAM. SAME FOR BU. HERE KO IS |
| C | Z1/RO• HERE RO=R*A1*A/10. |
| 2022 | READ $21, N 1, P H C, B L, T H L, B U, T H U$ |
|  | $B L=K 0 * B L$ \$ $B U=K 0 * B U$ |
| C | ADD EACH TIME ICOUNT $=1$ IGUS $=-1$ FOR SMALL POWER INCREMENTS. |
| C | ADD ICOUNT $=1$ IGUS $=1$ AND USE RESULTS OF TRANSIENT PROGRAM. |
| C | ADD ICOUNT $=1$ IGUS $=0$ TO START NON RF GUESS SEQUENCE. |
| C | ADD ICOUNT $=0,-1$ TO GET NEW POWER VALUES |
| 2020 | READ 2041, ICOUNT, IGUS,BI,V16,P16 |
| 2041 | FORMAT (2I5,3E20.9) |
|  | IF (EOF, 60) 22,23 |
| 22 | CALL EXIT |
| 23 | IF (ICOUNT)2022,2022,2023 |
| 2023 | AN5 $=\mathrm{H} *$ N1 $\$ \operatorname{KEY}(6)=\mathrm{N} 1$ |
|  | S4=N1*S2 |
|  | PRINT 2111,N1,PHC, BL, THL, BU, THU |
| 2111 | FORMAT (* N1, PHC, BL, THL, BU, THU * I5,5E12.4) |
| 21 | FORMAT (I5,E20.9 /2E20.9/2E20.9) |
|  | CBU $=$ BU*CEXP(-IM*THU) \$ CBL=BL*CEXP (-IM*THL) |
|  | ZF=CEXP (AN6) \$ZG=CEXP (AN7) \$ZH=CONJG(ZF)*KL*CBL |
|  | ZI=CONJG(ZG)*CBU*KU \$DO $24 \mathrm{I}=1, \mathrm{~N} 5$ \$AI=I-1 |
|  | ZH=ZH*ZF \$ZI=ZI*ZG \$ E=2.*(ZH+ZI) |
| 24 | PHB (I) $=\mathrm{E}+\mathrm{PHC}+\mathrm{AI*AN5}$ S \$ICONV=30 SICASE=ICASE+1 |
|  | IF (IGUS) $302,300,303$ |
| 302 | BI=U \$V16=VA(1) \$P16=PHA(1) \$ GO TO 303 |
| 300 | $\mathrm{BI}=A * S 3 *(S Q R T F(1 \cdot+S 4 * S 4)-54)$ |
| 301 | V16=P16=0. |
| 303 | CONTINUE |
|  | PRINT 2113 , ICOUNT, IGUS, BI, V16,P16 |
| 2113 | FORMAT(* ICOUNT, IGUS,BI,V16,P16 */2I5,3E20.9) |
|  | CALL PARTA \$IF (ICONV) $25,26,26$ |
| 25 | PRINT 205, N1, ICASE |
| 205 |  |
|  | GO TO 2020 |
| 26 | ICONV=10 \$CALL PARTS \$IF(ICONV)27,31,31 |
| 27 | PRINT 205,N1,ICASE |
| 31 | CALL CHARTB |
|  | GO TO 2020 |
| 18 | FORMAT (I5) |
| 199 | FORMAT(* $\mathrm{Z}_{1}, 22, \mathrm{Z3}(19)=$ * I5,3C(E12.3,E12.3)) |
|  | END |

SUBROUTINE SCALE(I9,I99)
PURPOSE IS TO CONVERT THE DC DOWN BY 10-7 AND FORM LOGIO(Y1)
THE AMPLITUDE OR POWER SPECTRA 5 DECADES ONLY.
TYPE REAL MIN,MAX
COMMON /4/ ICLE $(8,2,6)$, IDC(8)
COMMON /1/E(6151),Y1(1025) PF(11299)
GO TO $(1,2,3,4,5,6)$ I 99
$\operatorname{IDC}(3)=\operatorname{ICLE}(I 9+1,1,3)=I \operatorname{CLE}(I 9+1,2,3)=0 \quad \$$ RETURN CONTINUE
CONTINUE
CONTINUE
CONTINUE
CONTINUE
POWER CASE AND VOLT CASE
$M I N=0$ - 5 MAX $=0$.
IF (I9) 8,9
$\operatorname{IDC}(199)=-7 \quad \$ \quad Y I(3)=1 \cdot E-7 * Y 1(3)$
DO $11 \mathrm{I}=1,1024$
IF(YI(I)) 12,11,13
IF (MIN-Y1(I)) 11,11,14
MIN=Y1(I) \$ GO TO 11
IF(MAX-YI(I))15,11,11
MAX=Y1(I)
CONTINUE
IF (MIN) 31,32,32
IMIN=ALOG10(-MIN) \$ MIN=-1॰*(10.)**IMIN SICLE(I9+1,1,I99)=IMIN-5
IF (MAX) $33,33,34$
IMAX = ALOG10 (MAX) \$MAX $=(10) * * I M A X ~ \$ I C L E.(I 9+1,2, I 99)=I M A X-5$
DO $21 \mathrm{I}=1,1024 \mathrm{~S} \mathrm{IF}(\mathrm{Y} 1(\mathrm{I})) 22,21,23$
Y1(I) $=-($ ALOG10(Y1(I)/MIN)+5.) $\$ \mathrm{IF}(\mathrm{Y} 1(\mathrm{I}) / 21,21,25$
Y1(I) $=0$. $\$$ GO TO 21
Y1(I) $=A \operatorname{LOG1O}(Y 1(1) / M A X)+5 . \quad \$ I F(Y 1(I)) 24,21,21$
Y1(I) $=0$. S GO TO 21
CONTINUE $\$$ RETURN $\$ E N D$

```
    SUBROUTINE PART
    DATA( ALPHA=0.,1.,.2,.1,.01,.005)
    DIMENSION ALPHA(6)
    DIMENSION Z1(1025),Z2(1025),23(1025)
    COMMON /5/ BI,S1,52,S3
1,A1
    COMMON /16/ V16,P16,H8
    TYPE COMPLEX Z1,Z2,Z3,K,DT,ZF,ZH,ZG,KL,AN2,ATMP,KU,ATN,AN6,AN7,ZI
1,CBU,CBL
    TYPE REAL KO
    COMMON /DAT/ PI,IM,PHS,U
    DIMENSION K(1025),DT(1025),KL(3),KU(3)
1 ,I2T(1025),VAT(1025)
    EQUIVALENCE (K(1),DPH(1)),(DT(1),DVA(1))
    COMMON /1/ PHC,VA(2050),DPH(4,2050),
1DVA(4,2050),H,H1,H2,H3,H4,EP1,EP2,EP3,LBX(5),LBTR(5),EP4(6)
    EQUIVALENCE (DPH(6151),Z1(1)),(DVA(6151),Z2(1)),(Q4(1),Z3(1))
    COMMON /2/D(2050),PHA(2050),PHB(2050),Q1(2050),Q2(2050),Q3(2050),
1Q4(2050),KEY(6),NO,N1,N2,NH,NP,PO,VO,AO,V, THL,THU,BO,BL,BU,NL,NU
2,NB,KL,KU, AN5,H5, ANU,ANUP,ANUT,A,R1,RO,C,WO,T, R,PI2,AN2,AN4,
3AN3,MM,N5,N6,N7,ISKPD,SC, OUT(6,11),NR, ICONV,AJ,ICASE,
4POW1,POW2,POWJ,POWIN,KO,EP5,H7
5,CBU,CBL
    EQUIVALENCE( X(1),DPH(1)),(Y(1),DPH(2051)),(VAT(1),DVA(1)),
1(11T(1),DVA(2051)),(VOUT(1),DVA(4101)),(X1(1),DPH(1025))
    TYPE COMPLEX VAT,IIT,VOUT,I2T,IM
    DIMENSION X(1025),Y(1025),11T(1025),VOUT(1025),X1(1025),Y1(1025)
    EQUIVALENCE(Y1(1),DPH(4101)),(I2T(1),Q3(1))
    DIMENSION T1(4),T2(4),T3(4)
```

ENTRY PARTA

```
    PART USES PURE DIFFERENTIAL EQUATION AND TRAPAZOIDAL SPLINE
    APPROACH TO GET PHA,VA SHAPE IN TIME
    DO 1 I=1,N6 $VA(I)=0.
l
    PHA(I)=0.
    EP77=EP5
    EP78=.01*A $ZP=EP78**3
    EP74=EP78
    IF(NI)2,3
    U= A*SINF(PHB(I)) $GOTO 483
    U=BI
    VA(N5)=V16
483 ICONVI=12
    DO 460 I 19=1,10
    BET =GAM=XI=0.
    PHA(1)=P16 $ VA(1)=VA(N5) $ VA(N5)=0.
    ICONV1=ICONVI-2 $ IF(ICONVI)61,61,620
61 ICONV=2 $ GO TO 4
620 ICONV=ICONVI
4 ICONV=ICONV-1 $IF(ICONV) 460,461,461
460 CONTINUE $ IF(ABSF(PHA(N5)-PHA(1))-100.*EP2)411,413.413
411 IF(ABSF(VA(N5)-VA(1))-100.*EP2)412,413,413
412 ICONV=1 $ GO TO 405
413 ICONV=-1$ GO TO 405
461 U=U+GAM
    VA(1)=VA(1)+XI
    TA=H3*U $TB=PHA(1)+PHB(1) $Q3(1)=SINF(TB)
401 Q2(1)=COSF(TB) SDPH(1,1)=\operatorname{DPH}(2,1)=\operatorname{DVA}(2,1)=0.
    DVA(1,1)=1. $ T3(2)=H3 $T3(1)=0.
    DO 6 I=I,N2 $I3=I+I $TD=HI*VA(I)+TAtH2*Q3(I) $ IV=5
8 TB=PHA(I3)+PHB(I3) $Q2(13)=COSF(TB) $ TC=1•-H5*Q2(I3)
    IV=IV-1
    IF( IV)10,11,11
    PRINT 13 SCALL EXIT
    FORMAT(* CONVERGENCE PROBLEMS IN PART A* )
    Q3(I3)=SINF(TB)
    VA(I3)=TD+H2*Q3(I3)
    TB=PHA(I)-PHA(I 3)+H4*(VA(I3)+VA(I))
    DEL=TB/TC
    PHA(I3)=PHA(13)+DEL $ DEL=ABSF(DEL)
    IF(DEL-EP1)7,8,8
7 TB=PHA(I3)+PHB(I3)
    Q3(I3)=SINF(TB)$TC= Q2(I)+Q2(I3)$VA(I3)=TD+H2*Q3(I3)
    Q2(13)=COSF(TB)
    TB=H5*Q2(I3) $DO 5 11=1.2
    Tl(II)=DVA(I1,I)*(H1+TB)+T3(I1)+H2*DPH(11,I)*TC
    T2(I1)=1.-TB
    DVA(I1,I3)=T1(11)/T2(II)
    DPH(I1,I3):= DPH(11,I)+H4*(DVA(I1,13)+DVA(11,1))
    CONTINUE $GO TO 12
    B1=PHA(1)-PHA(N5) $B2=VA(1)-VA(N5)
    Al1=DPH(1,N5) $A12=DPH(2,N5)
    A21=DVA(1,N5)-1. $A22=DVA(2,N5)
    DET=Al1*A22-A12*A2l $Bl=Bl/DET $ B2=B2/DET
    XI =A22*B1-A12*B2 $GAM=A11*B2-A21*B1
    1F(ABSF(GAM)-EP5) 17,17,19
    IF(ABSF(XI)-EP74)14,14,19
17
    GAM=SIGNF(EP77,PHA(1)-PHA(N5))
    IF(ABSF(PHA(1)-PHA(N5))-.1)199,199,198
    XI=SIGNF(ZP,VA(4)*(VA(1)+VA(4)-VA(3)-VA(2)))
    GO TO 200
    XI=SIGNF(EP78,(VA(1)+VA(4)-VA(3)-VA(2))*VA(4))
        EP78=.9*EP78
        EP77=.90*EP77
        PRINT 1000,U,VA(1),PHA(1),GAM,BET,XI
    1,VA(N5),PHA(N5)
1000 FORMAT(* U,VA,PHA * /(6E13.4))
    IF(ABSF(GAM)-EP2)400,4,4
400 IF(ABSF(XI)-EPI)40,4,4
```

Figure 2. (Continued)
RETURN
PRINT 999,(PHA(I7),VA(I7),I7=1,N5)
FORMAT(8E12.4)
ENTRY PARTB
C THIS PART GETS THE INTEGRAL EQUATION PART USING QUASILINEAR
C ANALYSIS AND THE TRAPAZOIDAL METHOD FOR THE ERROR ESTIMATE
DPH(1,1)=DPH(2,1)=DPH(3,1)=DVA(2,1)=DVA(1,1)=0. \$DVA(3,1)=1.
6 0
30
21
$GO TO 44
GO TO (21,28,28,28,28,24,24,60,35)IWAY
Al2=-H4$All=A22=DET1= l. \$ T2(2)=T2(3)=0. \$T3(2)=H3 \$T3(3)=0.
DVA(2,1)=DPH(2,1)=0.
DO 22 I=1,N2 \$I2=I+1 \$A2l=-H2*Q2(I2) \$DET=DET1-A12*A21
T3(1)=-Q3(12)
T2(1)=-Q4(I2) \$DO 23 I6=1,3
Bl=(DPH(I6,I)+H4*DVA(I6,I)+T2(I6))/DET
B2=(H1*DVA(I6,I)+T3(I6)+H2*Q2(I)*DPH(I6,I))/DET
DPH(I6,I2)= B1*A22 -B2*A12
DVA(I6,12)= B2*A11 -B1*A21
CONTINUE
CONTINUE
B1= -DPH(1,N5) \$B2= -DVA(1,N5)
Al1=DPH(3,N5) \$A12=DPH(2,N5)
A21=DVA(3,N5)-1. \$A22=DVA(2,N5)
DET=A1l*A22-Al2*A2l \$Bl=B1/DET \$ B2=B2/DET
XI =A22*B1-A12*B2 SGAM=A11*B2-A21*B1
I 99=2 \$ GO TO 42
TEMP=EP4(1)$I99=1$ICONV=ICONV-1 $J6=1$IWAY=8 \$IF(ICONV) 39,26,26
DO 250 L6=2,6 \$ IF(EP4(L6)-TEMP)25,250,250
TEMP=EP4(L6) \$J6=L6
CONTINUE \$ IF (J6-1) 62,62,28
DO 32 I=1,N2
DPH(2,I)= ALPHA(J6)* DPH(1,I)+PHA(I)
DVA(2,I)= ALPHA(J6)* DVA(1,1)+VA(I) \$TC=U+ALPHA(J6)*GAM
GO TO (46,49)I99
29 EP4(J6)=0. SDO 34 I=1,N2
SMALL=Q4(I)**2+Q3(I)**2 \$ IF(SMALL-EP1 )341,341,344
PRINT 345,SMALL,U,I,J6
FORMAT(* ERROR AT POINT * ,2E12.3.2I5)
EP4(J6)=EP4(J6)+SMALL
CONTINUE
TEST=EP4(J6) \$ IWAY=J6 \$ J6= J6+1
IF( TEST -EP3)51,30,30
IWAY=9 \$GO TO 46
PRINT 63,EP4,EP3 SRETURN
PRINT 41,EP4,EP3
FORMAT( * 10 LOOPS IN QUASI NO CONVERGENCE */(6E13.4))
RETURN
DO 43 I=1,N2
DPH(1,I)=DPH(1,I)+XI *DPH(3,I)+GAM*DPH(2,I)
DVA(1,I)=DVA(1,I)+XI *DVA(3,I)+GAM*DVA(2,I)
IWAY=J6=2 \$ GO TO 30
DO 45 I= 1,N2
DPH(2,I)=PHA(I) SDVA(2,I)=VA(I)
CONTINUE \$ TC=U \$GO TO 49
DO 47 I=1,N2 \$PHA(I)=DPH(2,I) \$VA(I)=DVA(2,I)
CONTINUE \$U=TC \$GO TO 30
DO 52 I=1,N2 \$TB=DPH(2,I)+PHB(I) \$Q2(I)=COSF(TB)
Q1(I)=-A*SINF(TB)+TC -ANUT*DVA(2,I)
IF(ISKPD)54,52
DO 53 13=1,N2
L=I-13
IF(L)55,56,56
L=N2+L
CONTINUE
Q1(I)=Q1(I)-D(L+1)*DVA(2,I3)
CONTINUE
DVA(2,N5)=DVA(2,1)

```
\(\operatorname{DVA}(2, N 6)=\operatorname{DVA}(2,2)\)
\(\operatorname{DPH}(2, N 5)=\operatorname{DPH}(2,1)\)
\(\operatorname{DPH}(2, N 6)=\operatorname{DPH}(2,2)\)
\[
1
\]

FORMAT(* NO CONVERGENCE POSSIbLE * /(6E13.4))
Q2(N5)=Q2(1) SQ1(N5)=Q1(1) \$Q1(N6)=Q1(2)
Do \(50 \quad \mathrm{I}=1\), N 2 \$ \(12=\mathrm{I}+1\) \$ \(13=12+1\)
Q4(12) \(=\) DPH(2,12)-DPH(2,I)-H4*(DVA12,12) + DVA(2,1))

Q3(1) \(=\) Q3(N5)\$Q4(1) \(=Q_{4}\left(\mathrm{~N}_{5}\right)\) \$GO TO 29 SEND
subroutine chart
DIMENSION Z1(1025),Z2(1025),23(1025)
COMMON /5/ BI,S1,52,53
1, A1
COMMON /4/ ICLE \(8,2,6\) ),IDC(8)
COMMON /DDC/LU,LUC,IFL
COMMON IDD/ IN,IOR,IT,IS,IC,ICC,IX,IY
COMMON /TAB/ V1(60),V2(121),V3(19)
TYPE COMPLEX \(Z 1, Z 2, Z 3, K, D T, Z F, Z H, Z G, K L, A N 2, A T M P, K U, A T N, A M G, A N T, Z I\)
\(1, \mathrm{CBU}, \mathrm{CBL}\)
TYPE REAL KO
COMMON /DAT/ PI,IM,PHS,U
DIMENSION K(1025), DT(1025),KL(3),KU(3)
1 , I2T(1025),VAT(1025)
EQUIVALENCE (K(1), DPH(1)),(DT(1),DVA(1))
COMMON /1/ PHC,VA(2050),DPH(4,2050),
1DVA ( 4,2050 ), \(\mathrm{H}, \mathrm{H} 1, \mathrm{H} 2, \mathrm{H} 3, \mathrm{H} 4, E P 1, E P 2, E P 3, \operatorname{LBX}(5), \operatorname{LBTR}(5), E P 4(6)\)
EQUIVALENCE (DPH(6151),Z1(1)),(DVA(6151),Z2(1)),(Q4(1),Z3(1))
COMMON /2/D (2050), PHA(2050),PHB(2050),Q1(2050),Q2(2050),Q3(2050),
1Q4(2050), KEY(6),NO,N1,N2,NH,NP,PO,VO,AO,V, THL,THU,BO,BL,BU,NL,NU \(2, N B, K L, K U, A N 5, H 5, ~ A N U, A N U P, A N U T, A, R 1, R O, C, W O, T, R, P I 2, A N 2, A N 4\), 3AN3,MM,N5,N6,N7,ISKPD,SC, OUT(6,11),NR, ICONV,AJ,ICASE, 4POW1, POW2, POWJ, POWIN,KO,EP5,H7 \(5, \mathrm{CBU}, \mathrm{CBL}\)
EQUIVALENCE ( \(\mathrm{X}(1)\), DPH(1)),(Y(1), DPH(2051)), (VAT(1), DVA(1)),
1(I1T(1), DVA(2051)), (VOUT(1), DVA(4101)),(X1(1),DPH(1025))
TYPE COMPLEX VAT,IIT,VOUT,I2T, IM
DIMENSION X(1025),Y(1025),I1T(1025),VOUT(1025),X1(1025),Y1(1025)
EQUIVALENCE(Y1(1),DPH(4101)),(12T(1),Q3(1))
Entry charta
REWIND 1
WRITE TAPE 1,Z1,22,Z3
REWIND 1
IN=IOR=ITT=IC=0 \$IS=1 \$LU=LUC=0 \$IFL=30
CALL DDINIT(3,2OH EGJOHNSON,EXT 3706 )
SI=S1/(R1.*ANUT)
PRINT 100,A1,S1 \$ PRINT 110,A,ANUT
FORMAT(* A,ANUT *, 2 E12.4)
ENCODE \(\left(40,100, V_{1}\right) A_{1}, S_{1}\)
FORMAT(16H A1,SQR(A)/NUTO/, 2E12.4)
IX \(=89\) \$ \(I Y=500\) \$ CALL DDTAB \$ CALL DDTEXT \((5, \mathrm{~V} 1)\)
ENCODE \((480,1, V 1) N O, N 2, N B, N L, N U, A N U, A N U T, A, V O, A O, P O, W O, R O, C, R, H, A J\), \(1 T\)
FORMATIB9H THE JOSEPHSON SPECTRUM CONTROL PARAMETERS THAT REMAIN ICONSTANT 1/FOR THE ENTIRE RUN. \(1 /\), 40 H NO, DATA POINTS,BEAT,LOWER,U 2PPER FREQ.1/ \(5515,48 \mathrm{HI} / \mathrm{NU}, \mathrm{NU}\) TOTAL,SCALED A,VOLTS,AMPERES, WATTS 31/ , 6E12.4, 62H 1/RADIAN FREQ,OHMS, CAPACITANCE, RESISTANCE, ST 4EP INTERVAL'/, 5E12.4,36H "/SUPER CURRENT AMP.,PPERIOD SEC.1/ 5 2E12.4,24H \(1 /\) CASES TO FOLLOW. \(\cdot\)
ENCODE 32,99, LBX(1) \()\)
FORMAT 32 H JOSEPHSON CURRENT IN TIME CASE
)
IX \(=89 \$\) IY \(=950\) \$CALL DDTAB \$CALL DDTEXT( \(60, V 1\) ) SCALL DDFR SRETURN ENTRY CHARTB
READ TAPE 1,Z1,Z2,Z3
REWIND 1
C Q2(I) Is
A*SINF(PH)

Figure 2. (Continued)
```

$C$ DROP NP POINT IN THESE PLOTS. SEE PRINT OUT FOR IT. POW $1=P O W 2=P O W J=0$.
NS=NH/129+1 \$DO 15 I8=1,NS \$I9=128*(I8-1) \$K7=I8-1
DO 21 I= 1,128 \$ I2=(I-1)*7+3 \$II=I +I 9
21 Yl(I2)=CABS(I2T(II)) SENCODE(40,97,LBTR(I))ICASE,I8
97 FORMAT(26H AMPLITUDE I2TRAN, ICASE= 2I5, 4HPART,
97 FORMAT(26H AMPLITUDE I2TRAN, ICASE= 2I5, 4HPART,
CALL SCALE(K7,1)
CALL FILMGRAF(X1,Y1,1024,LBTR,1,0)
PLOT OF CURRENT,I2, AMPLITUDE THROUGH DETECTOR 23, AS A FUNCTION
C PLOT OF CUR
DO 22 I= I, 128 \$I 2=(I-I)*7+3 \$I I=I +I 9
22 Y1(I2)=CABS(VOUT(I1)) SENCODE(40,96,LBTR(1))ICASE,I8
96 FORMAT(26H AMPLITUDE VOUTTR, ICASE= 2I5, 4HPART)
DO 2 I = 1,N2
Q2(I) =A*SINF(PHA(I)+PHB(I))
XI(I)=0. \$DO 3 I =2,256
XI(I)=XI(I-1)+4. \$IFL=60
ENCODE(8,98,LBX(5))ICASE
FORMAT (I5)
NS=N2/257+1 SDO 4 I8=1,NS \$I9=256*(I8-1) \$DO 5 I=1,256 \$II=I+I9
Yl(I)=Q2(II) \$CALL FILMGRAF(XI,YI,256,LBX,1,0)
CONTINUE
DO 6 I= 1,N2
Q2(I)=Q2(I)+Q1(I)+ANU*VA(I)
DO }7\textrm{I}=1,NH \$X(I)=Q2(I
CONTINUE SDO }8\mathrm{ I=NP,N2
Y(I-NH)=Q2(I) \$CALL REORDER(X,Y,MM) \$CALL CFFTRC(X,Y,MM,SC,I)
CALL REALTRAN(X,Y,MM,I,I) \$DO 9OI=I,NP \$ IIT(I)=CMPLX(X(I),Y(I))
CONTINUE
CALL REORDER(A,B,M)
REVERSIBLE PERMUTATION OF REAL SEQUENCE
FROM FIRST-LAST-NORMAL SEQUENCE
TO ODD-EVEN-REVERSE BINARY SEQUENCE.
SEQUENCE LENGTH IS N = 2**M
CALL REALTRAN(A,B,M,NE,INV)
IF(INV.GT•O) UNSCRAMBLE THE TRANSFORM OF A REAL SEQU'_INCE.
IF(INV.LT.O) SCRAMBLE THE TRANSFORM OF A REAL SEQUENCE.
INPUT AND OUTPUT ARE IN NORMAL SEQUENCE,
SEQUENCE LENGTH IS N = 2**M
NE MUST AGREE WITH SIGN OF EXPONENT IN TRANSFORM DEFINITION.
INNER LOOP SINES AND COSINES COMPUTED
RECURSIVELY BY SINGLETONS 2ND DIFFERENCE ALGORITHM,
INITIALIZED FROM A DATA TABLE.
DISCRETE COMPLEX FAST FOURIER TRANSFORM.
CALL CFFTRC(A,B,M,SC,NX)
INPUT A(J) + I*B(J) IN REVERSE BINARY SEQUENCE.
OUTPUT A(K) + I*B(K) IN NORMAL SEQUENCE.
SEQUENCE LENGTH IS N = 2**M
SC IS REAL SCALING MULTIPLIER.
NX IS THE SIGN OF THE EXPONENT IN THE TRANSFORM DEFINITION.
INNER LOOP SINES AND COSINES COMPUTED
RECURSIVELY BY SINGLETONS 2ND DIFFERENTE ALGORITHM,
INITIALIZED FROM A DATA TABLE.
IIT(I)= IIT(1)+N1*ANU \$IIT(NU+1)=IIT(NU+1* (KU(2)*ANU+KU(3))*CBU
IIT(NL+1)=I1T(NL+1)+(KL(2)*ANU+KL(3))*CBL
DO 10 I= I,NH \$X(I)=VA(I)
CONTINUE \$DO ll I=NP,N2
Y(I-NH)=VA(I) \$CALL REORDER(X,Y,MM) \$CALL (r:FTRC(X,Y,MM,SC,I)
CALL REALTRAN(X,Y,MM,I,I) SDO 12 I=I,NP \$VAT(I)=CMPLX(X(I),Y(I))
CONTINUE
VAT(1)=VAT(1)+N1 \$VAT(NU+1)=VAT(NU+1)+KU(2) *CBU
VAT (NL+1)=VAT(NL+1)+KL(2) *CBL
VJUN=VO*VAT(1) \$ AIJUN=AO*I1T(1) \$ ENCODE(64,9111,VI)VJUN,AIJUN
FORMAT(30H VJUN(VOLTS), IJUN(AMPERES)'/ , 2E17.5)
DO 13 I= 1,NP
VOUT(I)= Z2(I)*IIT(I)+VAT(I) \$I2T(I)=VOUT(I)/Z3(I)
CONTINUE \$DO 14 I=1,1024 \$XI(I)=I-1
YI(I)=0.
DROP NP POINT IN THESE PLOTS. SEE PRINT OUT FOR IT.
POW1 = POW2=POWJ=0.

```
```

    CALL SCALE(K7,2)
    CALL FILMGRAF (X1,Y1,1024,LBTR,1,0)
    C PLOT OF OUTPUT VOLTAGE ACROSS DETECTOR
        DO 23 I= l, 128 $I2=(I-I)*7+3 $II=I +I9
    CONT INUE
        ENCODE(968,9,V2)POWJ,POW1,POW2,POWIN,KO,(KEY(I),(OUT(I,J),J=1,11)
        I,I=1,6)
    FORMAT(28H POWJ,POW1,POW2,POWIN,KO,1/ 5E12.3,2HI/ ,
    138HKEY,B,PHC,POWIN,POW1,POW2,'/Z1,Z2,Z31/
    2 6(I4,5El2.3.2H//,6El2.3.2H*/ )।
    ENCODE(152,91,V3)EP1,EP2,EP3,EP4
    ```

Figure 2. (Continued)
        CALL DDTAB \$ CALL DDTEXT(60,V2) \$ CALL DDFR

911 FORMAT 158 H SCALE EXP ONLY DC SET THEN MIN,MAX SET FOR EACH GRAPH 1\% , 6I4.2H1/,6(16I4.2H1/))
RETURN
END



THE FAST FOURIER TRANSFORM ROUTINES

REORDER CFFTRC
REALTRAN


THE PART SUBROUTINE GENERATES THE INITIAL GUESS FOR PH AND THE QUASILINEAR PROCEDURE FOR THE SELFCONSISTANT PH.

ENTRY PARTA
ENTRY PARTB

Figure 3. The flow chart for Josephson steady state program.
NO, J, A1, C,R,EP1,EP6
8.2277-015 1.0000+001 1.0000-008 1.0000+000

signals are to be applied, whether we want to compute the integral equation part, (ISKP is zero means that the frequency spectra of impedance is to be ignored. ISKP equal 1 means that it is considered and the integral equation is now used.); and finally what the initial case number label for this computation run is to be, ICASD(5). The remaining cases are incremented by one for each new case.

The N1, PHC, etc., row is repeated for each new case. Here \(\mathrm{Nl}(4)\) is the chosen rotation frequency for the Josephson oscillator, PHC ( \(-3.1416-0\) ) is the chosen initial phase assuming PHA is zero initial value. BL is the corrected value for the lower frequency r.f. applied voltage as defined on page 1. This \(B L(2.4+5)\) is changed from the \(\overline{B L}(7.68+2)\) put into the program and used in the transient program in the appendix. The lower r.f. phase is THL( \(0_{0}\) ) in this example. The BU( \(0_{0}\) ) is likewise scaled from \(\overline{\mathrm{BU}}(0\).\() and finally we have the upper initial phase THU(0.).\)

The next data row has two control variables ICOUNT and IGUS as well as input data that is used as initial guesses for \(U, V_{a}(1)\), and \(\varphi_{a}(1)\). The ICOUNT(1) has two options--if it is zero, it causes the program to request a new set of \(\mathrm{NL}, \mathrm{PHC}, \overline{\mathrm{BL}}, \mathrm{THL}, \overline{\mathrm{BU}}, \mathrm{THU}\); and if it is equal to one it causes the program to proceed to compute the case just read in. The IGUS has three options--if it is -1 , it causes the program to use the computed previous values of U, VA(1), PHA(1) as the initial guesses for the new set of data; (this option is useful if you are simply making small changes in previous data cases), if IGUS is zero, it uses some computed guesses based on the approximation that the r.f. power is not importantly affecting the initial data, hence the d. c. case can be a good basis for initial data guesses; and finally, if IGUS is one, it uses the currently read in values BI for U, V16 for VA(1), and Pl6 for PHA(1). These values are gotten from the program in the appendix or from other calculations.

The final rows of information in the printed output depends on how the calculation proceeds. If the calculation is normal, then you get what is shown in the sample case. If the calculation does not work normally, then you will get error messages and some detailed printouts. Please look at the Fortran listings to understand these messages and printouts. The U, VA, PHA sequence is repeated until there is convergence in the PARTA subroutine. The data printed out is the current \(U(-2.6+1)\), \(\mathrm{VA}(1)(4.96), \operatorname{PHA}(1)\left(-3.62, \delta u(1.6177)\right.\), change in \(\varphi_{a}(1)\), the change in VA(1), the VA(N5) (4.6127), and finally PHA(N5) (-3.6492). Please note that convergence occurs when the VA(1) equals VA(N5) and PHA(1) equals PHA(N5). Also note that the phase PHA(1) is unchanged for the entire calculation sequence. Here \((1, N 5)\) is equivalent to \(\left(0, N_{2}\right)\) in the analysis.

\section*{5. 2 The Microfilm Output}

We look at selected plots of the microfilm output of a single computer run. This data corresponds to the same run as the printed output. Figure 5 has information that remains constant throughout the different runs. The only information that may not be self evident is the SQR(A)/ NUT term (. 5) which is the characteristic frequency ratio of the circuit. It is \(\sqrt{a} / \nu_{T}\) and it measures the capacitance level; the smaller it is the less important is the capacitance term in the circuit.

Figure 6 has 8 plots of 256 points each of a single period of computation. Please note that the FILMGRAF routine will independently scale the individual plots so that the variations are full scale in the \(y\) axis. If the range of \(y\) variables are both negative and positive then the scale is plus and minus one. If the \(y\) variables are only one sign then the range is zero to minus one or one to zero
 con int lwitet un

it 2040 :
wu.nu totab.scales a, voits. anerets. nets




1.0500-004 1.6542.011
casts - 8 OLi20

\section*{-1.3certalati}
1. 8180.181 8 181-01

Figure 5. Sample microfilm output - the control parameters.
depending on the sign. We emphasize that it is possible for the different plots to have very different scales for their respective data. In this sample case the figure 6 set has exactly the same scale for all 8 plots.

Figures 7 and 13 are the first two plots of a series of 8 for the amplitude at each frequency of the current through the detector. The first figure gives the first 128 harmonics, the latter figure gives the next 128 harmonics. The remaining 768 harmonics are not shown. The d. c. line amplitude is reduced by \(10^{-7}\) in the first plot. Then, in this set of 128 harmonics as well as each succeeding set of 128 harmonics, the largest positive value is found, the \(\log _{10}\) is taken and the integer part is extracted. Then all positive nonzero y values are scaled relative to this integer. Only the top five decades of \(\log _{10}\) amplitude values are computed and put into the \(y\) axis. The remaining \(y\) values that are less than this allowed 5 dacades are set equal to zero and is thus plotted that way. For those plots which have negative nonzero \(y\) values the largest negative value is extracted, the integer part is removed, these y values are scaled down by that number, and the \(\log _{10}\) is taken and those numbers that are 5 decades less than this negative maximum are set to zero. The numbers are then made negative. Both the negative and positive values that have been processed as above are now subject to further scaling in the subroutine FILMGRAF. That scaling is the same as already discussed for the time series plots.

The above scaling takes place for all the figures 7 through 18 as well as those plots that have not been shown. Figure 19 has at its bottom a set of scaling information which in conjunction with the information given at the bottom of each figure permits the relative determination of the amplitudes for each spectral line that has been plotted. The row of -7's with a zero mean that only the phase difference series of plots does not scale the d. c. frequency line; the rest are scaled by this \(10^{-7}\). This was done to allow detailed plot information of the higher frequencies without the d. c. values causing scaling problems. After the -7 row are six rows of numbers. The first eight numbers of each row defines the scaling used on the negative numbers in the spectrum. Note that only the last row which corresponds to the power spectra across the junction has negative values. The last eight values correspond to the scaling of positive numbers. The rows are related to the plots as follows.
\begin{tabular}{ll}
\(\frac{\text { Row }}{1}\) & \multicolumn{1}{c}{\begin{tabular}{l} 
Plot \\
2
\end{tabular}} \\
3 & \begin{tabular}{l} 
Amplitude of the current across detector (AMPLITUDE I2TRAN)
\end{tabular} \\
4 & Phase difference between detector current and voltage \\
5 & (PHASE I2-VOUT TRAN) \\
6 & Voltage across R (AMPLITUDE VPHTR) \\
6 & Power across detector (POWER I2VOUT SSTR) \\
& Power across Josephson junction (POWER Il VOUR SSTR)
\end{tabular}

We illustrate the reading of the plots to get an absolute number for the power across the junction at the 4,8 harmonics and at the 132,133 harmonics respectively. The first two are gotten from figure 12 and the latter from figure 18. The conversion to watts is given in figure 5 as (4.8830 -11). The fourth harmonic has value full scale (5.918). The eighth harmonic has value (4.7 \(\times(4.4410) / 6.2\) \(=3.34)\).

\[
\begin{aligned}
& \text { JOSEPHSON CURRENT IN TIME CASE } 6 \\
& x \text { - Min Me. } \\
& \text { Puncrion } \\
& \text { Scaie 0.000.00 } \\
& 0.000 .000 \\
& 1.020 .003 \\
& \text {-. Min } \\
& -2.560 .002 \\
& \text { Me: } \\
& \text { 2. } 560.002 \\
& \text { 2. } 560.002 \\
& 6
\end{aligned}
\]

Run :972 Mey 65 (124) © \(17: 50: 52\)
い.:2 \(2=1.0\)
256
ints, Frese !

Figure 6. The next 8 plots give one complete period, each plot 256 points in time.


Figure 6. (Continued)


Figure 6. (Continued)

\[
\begin{aligned}
& \text { JOSEPHSON CURRENT IN TIME CASE } 6 \\
& \begin{array}{lllll} 
& X \cdot M i n & \text { Mas } & \text { Min } & \text { Min } \\
\text { Function } & 0.000 .000 & 1.020 .005 & -2.560 .002 & \text { Mas } \\
\text { Scoio } & 0.000 .000 & 1.020 .005 & -2.560 .002 & 2.560 .002
\end{array}
\end{aligned}
\]

Figure 6. (Continued)


Figure 6. (Continued)


Figure 6. (Continued)


\title{
JOSEPHSON CURRENT IN TIME CASE \\ 6 \\ \(X \cdot \operatorname{Min}\) \\ 0.000 .000 \\ Maı \\ 0.000 .000 \\ 1.020 .005 \\ \(r\) - Min \\ -2.560.002 \\ He: \\ 2.560 .002 \\ \(1.020 .003 \quad-2.560 .002\) \\ 2.560 .002
}
Run 1972 May 03 11241 of 17:50:53 L1.628 1.0 256 Poinis. Preme y

Figure 6. (Continued)


Figure 6. (Continued)


Figure 7. The first frequency block for the \(\mathrm{LOG}_{10}\) amplitude of the current through the detector. (Notice only \(\ell=4\), 8 , etc., a re non zero.) \(l\) is the "frequency" relative to \(w_{0}\).


Figure 8. The first frequency block for the \(\mathrm{LOG}_{10}\) amplitude of the voltage across the detector.


Figure 9. The first frequency block for the relative phase between the current and voltage of the detector.


Figure 10. The first frequency block for the LOG 10 amplitude of the voltage across the junction without \(z_{2}{ }^{\circ} 10\)


Figure 11. The first frequency block for the LOG \(_{10}\) power into the
detector.

POWER IIVOUT SSTR, ICASE=

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0.0000000
\(\gamma\). Min
- 4.1100000 - 4.110 .000

6 IPART

Me:
5. 910.000
5. 910.000

Figure 12. The first frequency block for the \(\mathrm{LOG}_{10}\) power into the Josephson junction.


Figure 13. The second frequency block ( \(128 \leq \ell<256\) ) for the LOG 10 current through the detector.


Figure 14. The second frequency block for the LOG \(_{10}\) voltage across the detector.


Figure 15. The second frequency block for the relative phase.



Run 1972 Koy 05 11241 of 17:51:05 b1.62: 1.0 1024 Poinsp. Ficae 10

Figure 16. The second frequency block for the LOG \({ }_{10}\) voltage across junction without \(z_{2}\).


Figure 17. The second 128 lines for \(\mathrm{LOG}_{10}\) power into detector.
```

POWER IIVOUT SSTR, ICASE=
6 2PART

|  | X - Min | Mo: | V-Min | Mos |
| :---: | :---: | :---: | :---: | :---: |
| Funcrion | 0.000 .000 | 1.023.005 | -4.350.000 | 4. 562.000 |
| Scalo | 0.000 .000 | 1.0250005 | -1. 262.000 | 4. 562.000 |

Run 1972 Rey 05 11241 of 17:51:04 L1.62: 1.0 1024 Points. Piene 20

```

Figure 18. The second 128 lines for LOG \({ }_{0}\) power out of junction. (Noise pulses are seen as enelgy sources here.)
(Here we have compared the length of the 8 th harmonic to the 16 th harmonic which is given as full scale of (4.410). We take the absolute value in this effort. The minus sign is simply a code to tell us that the power is out instead of in. The 132 th harmonic has value \((5.5 \times(4.350) / 6.2=3.85)\). The 133 rd harmonic has value ( \(3.5 \times(4.963) / 7.1=2.45\) ). The above values are extracted from the plots and the given scales in the plots. To complete the determination of the power we need to look at the conversion table of figure 19 to find that the positive 4 th harmonic is reduced by -2 , the negative 8 th is reduced by -6 , the positive 133 rd is reduced by -24 , and the negative 132 is reduced by -16 . Thus the power is given as
\begin{tabular}{cc} 
Harmonic & \(\frac{\text { Power }}{3.918}\) \\
8 & \(10^{3.9}\) \\
132 & \(-10^{-2.66}\) \\
133 & \(-10^{-12.15}\) \\
& \(10^{-21.55}\)
\end{tabular}

To scale these numbers to watts just multiply them by the \(4.8830 \times 10^{-11}\). Notice that the noise signal in the 133 harmonic is approximately 9 decades less than the corresponding non-noise signal of the 132 harmonic.

It should be obvious that if one wants extensive detailed tables of the various spectra, it is better to put in appropriate print statements into the programs. Our purpose has been mainly qualitative visual information of the power spectra. Figure 19 has a table of power and other information at selected harmonics. These numbers are in the dimensionless form used in the calculation. To convert to watts, volts, amperes, etc., it is necessary to use the constants given in figure 5.

We finish this section by describing the remaining details shown in figure 19. The first row of data gives the total steady state power across the junction without the \(Z_{2}\) part (POWJ), the total steady state power across the full junction (POW1), the power across the detector (POW2), the power put in by the applied voltages (POWIN), and finally the effective impedance KO. The next block of six different row pairs give the harmonic number to be applied to the row pair information (KEY), the applied voltages ( \(B\) ), the net initial phase at \(s=0\) ( PHC ), the power in (POWIN), the power across the junction (POW1), and the power across the detector (POW2). The second row of a given row pair has the dimensionless values for the complex number impedances of that frequency. The frequencies chosen to be printed out are d. c., the direct beat frequency, the lower applied r.f. frequency, the upper applied r.f. frequency, the highest frequency used as a consequence of the number of data points computed, and finally the rotation frequency.

The remaining undiscussed information has the net \(d\). c. voltage across, and the current through the junction without the \(Z_{2}\) term, the data summarizing the error controls used in the program (EPl, EP2, EP3) and finally the last computed P( \(\alpha\) ) set EP4(I6). Please look at the program to see what these errors limits mean in detail.

\section*{6. THE SUMMARY}

In this report we are documenting the computer program, the mathematical techniques and approximations used, and the procedures for use and interpretation of the computer program. The


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Figure 19. Scaling information for a particular steady state run.
results using this program are to be shown in other reports.
Here the discussion is directed toward the procedures used to check the correctness of the computer program. The strict \(d\). c. cases have been selectively compared with the Stewart voltagecurrent model. The r.f. applied voltage case has been selectively compared with the Fack, et al, 4 model and their reported results. Both cases agree within the accuracy of the techniques and the graphs of the data reported in those papers. The final check has been a comparison with the voltage driven model. Again the results are in agreement. The integral equation part has not been checked against any other model or results since thereare no such published results. In terms of the computer program, the program has been checked to be sure that the \(D\) is correctly generated. The only remaining part for possible programming errors is in the integral equation section in subroutine PART. The behavior and visual check shows no obvious errors. A detailed quantitative test is not known.

The most important numerical conclusion was to establish that it is very necessary to fix the \(\varphi_{a}(0)\) and to allow the \(V_{a}(o)\) to be adjusted when one wants to pin the rotation frequency. If one wants to use only time domain information, then one should use the much simpler programs such as used by Fack, Stewart, and is used in the appendix and fix \(b_{o}\) instead of the rotation frequency. If one is interested in high order harmonic mixing, and detailed harmonic spectra, then this program is best. If noise effects on the system of equations are of interest, then it is necessary to use the time sequence procedures and to allow selected Fourier transforms on the various time series. That analysis will be discussed in another report.

Because the analysis has been done for discrete uniformly spaced time points, we have discrete induced frequencies that are rationally related to each other. The consequences of this situation are additional steps induced in the d. c. current-voltage plots. The excursion of those steps is reduced by increasing the number of time points used. To get an idea on the range of any d. c. step for a given rotation frequency, it is necessary to make at least four calculations. First we chose a particular PHA (1) \(\equiv \varphi_{0}\) and compute the \(B_{1}\) for that case. If the program converges, then we have one point for that case. To get a convergence case, we use the transient program with the chosen \(B_{1}\) and deduce an adequate PHA(1). The steady state program just improves the accuracy of the results. Second, we compute a \(\mathrm{B}_{ \pm}\)with PHA(1) now \(\varphi_{0} \pm \frac{\pi}{2}\). At least one of these will converge. If both converge, then \(\left(B_{+}+B_{-}\right) / 2\) is the center and \(\left|\left(B_{+}-B_{-}\right)\right|\)is the range of the step. For example, if the \(B_{+}\left(B_{-}\right)\)does not converge because the phase location denies steady state solution, then it is necessary to compute \(\varphi_{0}+\frac{\pi}{4},\left(\varphi_{0}-\frac{\pi}{4}\right)\) case to get the \(B_{+\frac{1}{2}}\), (B \(\left.\frac{1}{2}\right)\) value. We assume that the range of a step is defined by the following simple formulae-- \(B=X+R \sin \theta\), where \(X\) is the center of the step, \(R\) is the range and \(\theta\) is the phase relation relative to the center of step. Normally \(\theta\) is not equal to \(\varphi_{o}\) and normally the formulae is not strictly accurate. It is only true if the dominate processes creating a step are due to a single frequencing mixing. For estimating the range's order of magnitude, we can accept this approximate formulae to get the range and center of the step for the case considered, namely \(B_{-}\left(B_{+}\right)\), \(B\) respectively are given by
\[
\begin{aligned}
& B_{1}=X+R \sin \theta \\
& B_{\mp \frac{1}{2}}=X+R \sin (\theta \mp \pi / 4) \\
& B_{\mp}=X+R \sin (\theta \mp \pi / 2)
\end{aligned}
\]
\[
\begin{aligned}
& \mathrm{X}=\left(\mathrm{B}_{1}+\mathrm{B}_{\mp}-\sqrt{2} \mathrm{~B}_{\mp \frac{1}{2}}\right) /(2-\sqrt{2}) \\
& \mathrm{R}=\left[\left(\mathrm{B}_{1}-X\right)^{2}+\left(\mathrm{B}_{\mp}-X\right)^{2}\right]^{\frac{1}{2}} \\
& \mathrm{~B}_{1}-X \\
& \mathrm{R}
\end{aligned} \mathrm{X}^{-\mathrm{Sin} \theta \text { or } \theta \text { is gotten to give center. For more }} \text { accurate analysis, it is necessary to do empirical work. }
\]

\section*{7. APPENDIX}

In this section we discuss briefly a simpler computer program called a transient program that can be used to get improved estimates for the effects of r.f. and d. c. applied voltages to the Josephson equations without the operator D. The equations used are closely related to those given in references land 4. We use the trapezoidal spline method with variable step size so that the accuracy per step is kept uniform throughout. Figure 20 gives the Fortran listing for this program. This program simply generates a time series of selected points. This data can then be used to estimate the average rotation frequency, \(\mathrm{N}_{1}\), and to fix a PHA(1), VA(1), and a U .

\section*{7. 1 The Transient Equation}

Here we show the equations used in the program and describe how they relate to equation 11. The transient equations are:
\[
\begin{gather*}
\frac{d \varphi}{d s}=V \\
\frac{d V}{d s}=-\nu_{T} V-a \sin \varphi+\bar{B}+\bar{B}_{L} \cos \left(N_{L} s+\theta_{L}\right)+\bar{B}_{U} \cos \left(N_{U} s+\theta_{U}\right)  \tag{Al}\\
\varphi=\varphi_{a}+\varphi_{b}+\varphi_{c}
\end{gather*}
\]

Here
where
\[
\frac{d \varphi_{b}}{d s}=V_{b}, \frac{d V_{b}}{d s}=-\nu_{T} V_{b}+\bar{B}_{L} \cos \left(N_{L} s+\theta_{L}\right)+\bar{B}_{U} \cos \left(N_{U} s+\theta_{U}\right)
\]
and \(\varphi_{c}\) is the initial value of \(\varphi_{a}\) at \(s=0\); thus here \(\varphi_{a}=0\) at \(s=0\). We note that \(\bar{B}=b_{o} / K_{0}, \bar{B}_{L}=b_{L} / K_{0}\), and \(\bar{B}_{U}=b_{U} / K_{0}\). The best method to fix \(a, \nu_{T}\), and \(K_{0}\) is to use the main program first for a given \(C, R, Z_{i}, a_{l}, J\), etc., with a d. c. case. We then use this program called the transient calculation for the various \(\bar{B}_{L}, \bar{B}_{U}, \theta_{L}, \theta_{U}\) cases. Please note that the scale used for \(b_{L}, b_{U}\) necessarily gives different phases and amplitudes in the main program if the \(K\left(N_{L}\right), K\left(N_{U}\right)\) are different from \(K_{0}\). When the programs are used, this difference needs to be considered. Before we look at the input and output of this program, we briefly consider the criteria for the step size. It simply restricts the absolute change in \(\varphi\) to be less than STEP/ (3Vl-V2) and the absolute change in \(V\) to be less than STEP*a/ [ \(\left.\nu_{T}(3 Q 1-Q 2)\right]\). The STEP is the control accuracy for this allowed change per step; the Vl and V2 are the last two \(V^{\prime}\) s used; the \(Q 1\) and \(Q 2\) are the last two \(q\) 's used. The \(q\) is \(\frac{d V}{d s}\).

This program is organized under the assumption that a given computer run will fix the r.f. and selected control parameters and will sweep the applied d.c. current \(B\) over many values. Thus the first read statement has a parameter which fixed the range of \(s\). The initial \(V\), and PH has no direct correspondence to the expected output VA and PHA, hence they are used to just start the calculation. The IPRINT defines how many steps are skipped before the usual print mode takes place. Figure 21 shows the flow of input data. Figure 22 gives an example of printed output data. The first three rows of printed output simply gives the summary of the input data use in this particular computer run. The list of names \(S\), VA, etc., are the labels for the six columns of numbers that follow the Warning statement. The \(S\) is the time variable. The VA is the \(\varphi_{a}\) with the average steady state slope of the ramp term not removed. PH is the total phase \(\varphi\). PHB is the r.f. phase including the initial PHC. VR is the integral of \(V\); it is used to estimate the slope of the ramp. PHA is the \(\varphi_{a}\) phase with the ramp term still present. To illustrate how to get the necessary initial data for the main program from this transient program, we run through the sample output to get the proper numbers. Please note that the printed output far exceeds the necessary information for these variables. It may be possible to suppress all printing except the few numbers of interest. The current calculation assumes that the transients are unimportant after \(S\) exceeds \(2 \pi\), hence we let the computation only run about one more \(2 \pi\) interval. This last interval is used to estimate the ramp and the needed initial PHA(1), VA(1), and \(U\) in the main program. Note that \(\mathrm{P} 16, V 16\), and \(B I\) are respectively used as the input variables for this initial data. In our sample output data we have the following needed data:
\begin{tabular}{clccc}
\(S\) & VA & PHB & VR & PHA \\
6.28245 & 8.665 & -.756 & 20.74 & 21.486 \\
12.57033 & 8.933 & -.645 & 46.03 & 46.66
\end{tabular}
\(B=102\), and \(\nu_{T}=32\). We use subscript 1 and 2 to label each data now. We chose these points because we want PHB to be nearly the same in both this program and in the main program. If you change to other THL and THU then you can consider other data points. Keep in mind that it is not necessary to have accurate values. The purpose is to get a guess that is at least accurate to 10 to \(20 \%\) so that the circle of convergence for the quasilinear procedure will operate effectively in the main program and so that some information is available to estimate potential effects due to spiking. We can then choose the necessary accuracy in the main calculation. The ramp slope is estimated as
\[
\tilde{N}_{1}=\left(V R_{2}-V R_{1}\right) /\left(s_{2}-s_{1}\right)+4.0
\]

The initial data is thus estimated as
\[
\begin{aligned}
\mathrm{BI} & =\mathrm{B}-\nu_{\mathrm{T}} * \widetilde{\mathrm{~N}}_{1}=-26 \\
\mathrm{~V} 16 & =\left(\mathrm{VA}_{1}+\mathrm{VA}_{2}\right) / 2 .-\tilde{\mathrm{N}}_{1}=4.8 \\
\mathrm{P} 16 & =\left(\mathrm{PHA}_{1}-\tilde{\mathrm{N}}_{1} * 2 \pi\right)+\left(\mathrm{PHA}_{2}-\tilde{\mathrm{N}}_{1} * 4 \pi\right) / 2 . \\
& =\left(\mathrm{PHA}_{1}+\mathrm{PHA}_{2} / 2 .-\tilde{\mathrm{N}}_{1} * 3 \pi=-3.626\right.
\end{aligned}
\]

PROGRAM JOSPA
C TO GET INTIAL WAVE SHAPE FOR STARTING GUESSES IN JOSPHC C PHDOT \(=V\) \$ VDOT \(=B-V * A N U T-A * S I N F(P H)+B L * C O S F(N L * S+T H L)+\) C \(\quad B U * \operatorname{COSF}(N U * S+T H U)\)

COMMON/1/B,ANUT,BL,THL,BU,THU,NL,NU,A,STEP,STEP1
READ 1 , RANGE, \(V, P H, B L, B U, T H L, T H U, S T E P, A N U T, A, I P R I N T, N L, N U\)
READ 1, B
70
1
FORMAT (5E12.3/5E12.3/3I10)
IF (EOF, 60) 71,72
71 CALL EXIT
```

SUBROUTINE COM(PH,V,Q,S,PHX,VX,QX,SX,HX,PHY)
COMMON /1/ B,ANUT,BL,THL,BU,THU,NL,NU,A,STEP,STEP1
DATA(ERO=1•E-5)
H=HX
H3=H*ANUT/2. \$H4=1\bullet-H3 \$H5=1\bullet+H3 \$S=SX +HX
CALL QR(O,QY,QZ,PHY) \$ H6=(H*H**25)/H5 \$H9=H4/H5 \$H10=H/(2.*H5)
H7=1./(1.- -H6*QZ) \$H8=VX*H/H5+PHX+H6*QX \$QZ1=QF(S) \$QY=QY+QZ1
IV=5
4 IV=IV-1 \$ IF(IV)5,6,6
6 DEL = (H8-PHY +H6*OY)*H7
PHY=PHY +DEL \$CALL QR(1,QY,O.,PHY) SQY=QY+QZ1
IF(ABSF(DEL)-ERO) 8,4,4
PRINT 7,S \$CALL EXIT
FORMAT (* NO CONVERGENCE *,E12.3)
PH=PHY \$Q=QY \$V=(H9*VX+H1O*(QY+QX)) SRETURN \$END

```

Figure 20. The transient Fortran program listing.

SUBROUTINE PROJ(PHY,H1,H,V2,V1,Q2,Q1,PH1) COMMON /I/ B,ANUT,BL,THL,BU,THU,NL,NU,A,STEP,STEP1 \(\mathrm{H} 1 \mathrm{~A}=\operatorname{STEP} / \operatorname{ABSF}\left(3 \cdot * \mathrm{~V}_{1}-\mathrm{V}_{2}\right)\)
\(\mathrm{H} 1 \mathrm{~B}=\mathrm{STEP} /\) /ABSF (3.*Q1-Q2)
\(\mathrm{HI}=\mathrm{MINIF}(\mathrm{HI} A, \mathrm{HIB})\)
PHY \(=\mathrm{PH} 1+\mathrm{Hl}\) * \(\left(\mathrm{V}_{1}+\left(\mathrm{V}_{1}-\mathrm{V} 2\right) * \mathrm{Hl} /(2 . * H)\right) \quad\) SRETURN SEND

SUBROUTINE START(PHX,SX,QX)
COMMON /I/ B,ANUT,BL,THL,BU,THU,NL,NU,A,STEP,STEPI CALL QR(1,QX,O.,PHX) \$QX=QX+QF(SX) \$RETURN \$END

FUNCTION QF(Z)
COMMON /I/ B,ANUT,BL,THL,BU,THU,NL,NU,A,STEP,STEPI \(Q F=B+B L * \operatorname{COSF}(N L * Z+T H L)+B U * \operatorname{COSF}(N U * Z+T H U)\) SRETURN SEND

SUBROUTINE QR(I,QA,QB,X)
COMMON /1/ B,ANUT,BL,THL,BU,THU,NL,NU,A,STEP,STEPI IF(I)I, I, 3
\(Q B=-A * \operatorname{Cos} F(X)\)
\(Q A=-A * \operatorname{SINF}(X)\) SRETURNSEND
```

PROGRAM JOSPA
READ 1, RANGE, ETC.
70 READ 1, B
IF(EOF) EXIT
IF(S-RANGE)4,4,70
4 COMPLETE
PRINT

```

\section*{SUBROUTINES CALLED}

COM
PROJ
START
QF
QR

Figure 21. The flow chart for Josephson transient program.


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Note that only integer \(\widetilde{\mathrm{N}}_{1}\) can be used in the main program and that the choice of \(B\) in the transient program does not necessarily correspond to such integers.

\section*{8. ACKNOWLEDGMENTS}

We wish to thank Stephen Jarvis and Clark Hamilton for useful discussions.

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16. ABSTRACT (A 200-word or less factual summary of most significant information. If document includes a significant bibliography or literature survey, mention it here.)

A computer program has been developed to study power flow between differen frequency channels in a Josephson junction circuit. This paper discusses the mathematical assumptions used to get such results. They are the trapezoidal approximation from spline theory and the use of a finite range of frequencies cha racterize the frequency spectrum. This paper describes the program and provides the FORTRAN listing, flow charts, and discusses how to use the program. A discussion of possible sources of errors is also included.
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[^0]:    ${ }^{1}$ Heal marters and Laboratories at Gaitlersburg, Maryland, unless otherwise noted; mailinh phes Washiugton, D.C. 20234.
    ${ }^{2}$ Part of the Center for Radiation Research.
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    ${ }^{1}$ Part of the Center for Building Technology.

