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## Corrections and Calculations on

 an X-ray Diffraction Line Profile: A Computer Program
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# Corrections and Calculations on an X-ray Diffraction Line Profile: A Computer Program 

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Corrections and Calculations on an X-ray
Diffraction Line Profile: A Computer Program
Raymond E. Schramm
This computer program was written to perform corrections and make calculations on an x-ray diffraction profile before Fourier analysis. The corrections are for background and for variations of intensity with the Bragg angle. Also calculated are the separation of the $K \alpha_{1}-K \alpha_{2}$ doublet and the centroid and position of peak maximum with their standard deviations. There is also an option to smooth the profile.

Key Words: Computer programs, Fourier analysis, nickel steels, x-ray analysis, x-ray diffraction, statistics.

1. Introduction

Annealed and heavily strained x-ray diffraction line profiles (fig. 1) can be compared to provide information on parameters such as particle size, dislocation density, and stacking fault energy. One scheme to do this is Fourier analysis. Gazzara, Stiglich, Meyer, and Hansen ${ }^{[1]}$ wrote a computer program "UNFOLD" to perform the mathematical operations; several other programs are also available. ${ }^{[2,3]}$ In working with a series of iron-nickel alloys, it became necessary to make corrections and calculations on the $x-r a y$ data in addition to those available in the program of Gazzara, et al. The program "PREP" was written to unify the preparation for programs such as "UNFOLD" by making all of the following calculations.

Within subroutines corrections are made for:

1. Background intensity.
2. Variation of intensity with the Bragg angle due to geometrical factors.

Subroutines are also used to determine:

1. Rachinger separation ${ }^{[4]}$ of the $K \alpha_{1}-K \alpha_{2}$ doublet.
2. Centroids of the $K \alpha$ and $K \alpha_{1}$ profiles with estimated standard deviations.
3. Peak maxima of the $K \alpha$ and $K \alpha_{1}$ profiles with estimated standard deviations.
${ }^{1}$ Figures in brackets indicate the literature references at the end of this paper.


Table 1. List of Symbols Used.

Symbol
Page of
Meaning
First Appearance

| a | 4 | Background intensity on low angle side of diffraction lines. |
| :---: | :---: | :---: |
| b | 4 | Background intensity at an arbitrary point. |
| c | 7 | Composite angular correction factor. |
| d | 8 | Angular separation of $K \alpha_{1}-K \alpha_{2}$ doublet. |
| $f$ | 5 | Atomic scattering factor. |
| h | 7 | Cylindrical specimen height. |
| m | 7 | Specimen density. |
| n | 4 | Number of experimental points or pairs of $2 \theta_{i}$ and $I\left(2 \theta_{i}\right)$. |
| r | 5 | Cylindrical specimen radius. |
| x | 7 | Mass fraction of Nickel |
| z | 4 | Background intensity on high angle side of diffraction line. |
| A | 5 | Sample absorption factor. |
| B, C, D, | 13 | Parameters of parabolic fit to peak. |
| G | 2 | Profile center of gravity. |
| I, $I_{1}, I_{2}$ | 4 | Reflected intensity (counts/unit time) due to the radiations $K \alpha, K \alpha_{1}$, and $K \alpha_{2}$ respectively. |
| $\mathrm{K} \alpha, \mathrm{K} \alpha_{1}, \mathrm{~K} \alpha_{2}$ | 1 | Characteristic radiations of the X-ray tube target. |
| LP | 5 | Lorentz-polarization factor. |
| $\alpha$ | 5 | Bragg angle of crystal monochromator. |
| $52 \theta$ | 14 | Increment in $2 \theta$ between individual measurements. |
| $\Delta 2 \theta$ | 10 | Angular span between first and last measurements. |
| $\theta$ | 1 | Specimen Bragg angle of reflection; $2 \theta$ is twice this angle. $2 \theta_{\text {max }}$ is the position of the peak maximum. |
| $\lambda, \lambda_{1}, \lambda_{2}$ | 8 | W avelengths of $K \alpha, K \alpha_{1}$, and $K \alpha_{2}$ radiations respectively. |
| $\mu$ | 5 | Linear X-ray absorption coefficients. |
| $\rho$ | 7 | Specimen mass density |
| $\rho_{s}$ | 7 | Mass density of powder sample. |
| $\sigma$ | 11 | Standard deviation of center of gravity and peak maximum. |

It has been found advantageous to smooth the profiles of cold-worked samples before Fourier analysis so an option to do this has been provided. It is also possible to have the corrected profile data punched in a format ready to enter "UNFOLD".

Angular positions are entered invalues of $2 \theta$ (twice the Bragg angle). The diffraction intensities must be in the fixed-time mode (intensity count at each angular increment determined for a fixed length of time) and measured at equal increments of $2 \theta$ across the entire profile. Several constants are written into the program assuming Debye-Scherrer geometry, the use of a monochromator with Co $K \alpha$ radiation, and iron-nickel or silver samples. The changes necessary when using other geometries, radiations, or samples are listed in the Appendix.

This program has been written in FORTRAN language and executed on a CDC $3800^{*}$ computer.
2. Background Calculation, "BACK"

The background intensity is the measured radiation that is scattered into the detector by means other than primary Bragg reflection of the $\mathrm{K}_{\alpha}$ component by the specimen. It generally arises from a number of sources: fluorescent radiation from the specimen, diffraction of the continuous spectrum, incoherent scattering from the specimen, and diffraction or scatter from something other than the specimen itself (e.g., collimator, air, etc.).

This calculation is made by arbitrarily averaging the intensities of the first five points and the last five points and then drawing a straight line through these two averaged points.

## Line Profile


*The use of trade names in this paper in no way implies endorseme nt or approval by NBS and is included only to define the procedure.

Set:

$$
\begin{aligned}
& a=\frac{1}{5} \sum_{i=1}^{5} I\left(2 \theta_{i}\right) \\
& z=\frac{1}{5} \sum_{i=0}^{4} I\left(2 \theta_{n-i}\right)
\end{aligned}
$$

Then

$$
b_{i}=\left(\frac{z-a}{2 \theta_{n}-2 \theta_{1}}\right) 2 \theta_{i}+\frac{a\left(2 \theta_{n}\right)-z\left(2 \theta_{1}\right)}{2 \theta_{n}-2 \theta_{1}}
$$

where $b_{i}$ is the background intensity at $2 \theta_{i}$.
Background subtraction is an available option in "UNFOLD".

## 3. Angular Correction, "ANG"

There are several factors affecting the diffracted intensity which are functions of the Bragg angles. This angular correction is divided into three parts.

### 3.1. Lorentz-polarization Factor (LP)

When a monochromator is used ${ }^{[5]}$

$$
L P=\frac{1+\cos ^{2} 2 \alpha \cos ^{2} 2 \theta}{\sin ^{2} \theta \cos \theta}
$$

where $\alpha$ is the Bragg angle of the crystal in the monochromator and $\theta$ is the Bragg angle of the specimen.

## 3. 2. Atomic Scattering Factor (f)

The atomic scattering factor $(f)$ is a function of $\frac{\sin \theta}{\lambda}, \lambda$ being the wavelength of the radiation being used. Tables of this factor for various values of $\frac{\sin \theta}{\lambda}$ and for various atomic models are available ${ }^{[6]}$. The values used are given in table 2 in a normalized form (from self-consistent wave functions) for iron, nickel, and silver.

To determine $f$ at intermediate values of $\frac{\sin \theta}{\lambda}$ a power series polynomial was fitted to the values in table 2. For a better fit $\frac{\sin \theta}{\lambda}$ was divided into three ranges. The values for this fit are given in table 3.

Table 2. - Normalized atomic scattering factors, $f$

| $\frac{\sin \theta}{\lambda}$ | Fe | Ni | Ag |
| :--- | :--- | :--- | :--- |
| 0.00 | 1.000 | 1.000 | 1.000 |
| 0.05 | 0.9731 | 0.9768 | 0.9817 |
| 0.10 | 0.9108 | 0.9214 | 0.9336 |
| 0.15 | 0.8404 | 0.8568 | 0.8698 |
| 0.20 | 0.7727 | 0.7925 | 0.8017 |
| 0.25 | 0.7077 | 0.7300 | 0.7362 |
| 0.30 | 0.6450 | 0.6689 | 0.6762 |
| 0.35 | 0.5323 |  | 0.6234 |
| 0.40 | 0.4412 | 0.5557 | 0.5781 |
| 0.50 | 0.3735 | 0.4611 | 0.5066 |
| 0.60 | 0.3258 | 0.3875 | 0.4547 |
| 0.70 |  | 0.3332 | 0.4151 |

Table 3. - Power series coefficients for the atomic scattering factors

$$
f=a+b\left(\frac{\sin \theta}{\lambda}\right)+c\left(\frac{\sin \theta}{\lambda}\right)^{2}+d\left(\frac{\sin \theta}{\lambda}\right)^{3}
$$

| Coefficient | Fe | Ni | Ag | Range of $\frac{\sin \theta}{\lambda}$ |
| :---: | :---: | :---: | :---: | :---: |
| a | +1.000 | +1.000 | $+1.000$ | $0 \leq \frac{\sin \theta}{\lambda} \leq 0.12$ |
| b | -0. 184 | -0.142 | -0.068 |  |
| c | -7.080 | -6.644 | -5.960 |  |
| a | +1.0637002 | +1.0645004 | +1.0685171 | $0.12<\frac{\sin \theta}{\lambda} \leq 0.35$ |
| b | -1.6056701 | -1.4746725 | -1.2033810 |  |
| c | +0.86001593 | +0.68002691 | -1.2685714 |  |
| d | -0.53335684 | -0.53337307 | +3.0666667 |  |
| a | +1. 1987014 | +1.1781009 | +1.2061000 | $\frac{\sin \theta}{\lambda}>0.35$ |
| b | -2.3833411 | -2.1006716 | -2.4973333 |  |
| c | +2.0200144 | +4.4750092 | +2.8050000 |  |
| d | -0.56667538 | -0.28333892 | -1.2166667 |  |

In the case of the iron-nickel alloys it was arbitrarily decided to use the atomic scattering factor for Fe for all bcc alloys and the atomic scattering factor for Ni for all fcc alloys. This gives a reasonable approximation to the electron density distribution in the alloy.

### 3.3. Sample Absorption (A)

For a flat specimen in a focusing geometry the sample absorption factor is independent of $2 \theta^{[7]}$. However, for a Debye-Scherrer geometry ${ }^{[8,9]}$ a cylindrical sample is used and Bradley ${ }^{[10]}$ gives the absorption factor:

$$
A=\frac{1}{\pi \mu r}\left\{1-\frac{\ell n(\sec \theta+\tan \theta)}{\sec \theta \tan \theta}\right\}+\frac{\sin 2 \theta}{2 \pi \mu^{2} r^{2}}+\frac{1}{\pi \mu^{3} r^{3}}\left[-\frac{1}{4}+\frac{3}{4} \cos ^{2} \theta\left\{\frac{1}{2}+\frac{1}{2} \frac{\ell n(\sec \theta+\tan \theta)}{\sec \theta \tan \theta}\right\}\right]
$$ for $\mu r>10$, where $\mu$ is the linear absorption coefficient and $r$ is the radius of the cylindrical specimen.

Using trigonometric identities this factor can be restated:
$A=\frac{1}{\pi \mu r}\left\{1-\frac{\cos ^{2} \theta \ell n\left(\frac{1+\sin \theta}{\cos \theta}\right)}{\sin \theta}\right\}+\frac{\sin 2 \theta}{2 \pi(\mu r)^{2}}+\frac{1}{\pi(\mu r)^{3}}\left[-\frac{1}{4}+\frac{3}{8} \cos ^{2} \theta\left\{1+\frac{\cos ^{2} \theta \ell n\left(\frac{1+\sin \theta}{\cos \theta}\right)}{\sin \theta}\right\}\right]$.

Bradley notes that for a diluted specimen (as in this case, we have used a powder mixed with cement):

$$
\mu r=\frac{\mu}{\rho} \frac{m}{\pi r h},
$$

where $\frac{\mu}{\rho}$ is the mass absorption coefficient, $m$ is the mass of the powder in a specimen of height $h$ and radius $r$. For a powder with three components (as $\mathrm{Fe}, \mathrm{Ni}$, and inactive cement)

$$
\mu r=\left[\left(\frac{\mu}{\rho}\right)_{N i} x+\left(\frac{\mu}{\rho}\right)_{F e}(1-x)\right] \frac{m}{\pi r h}
$$

where x is the mass fraction of the Ni. Now:

$$
\frac{m}{\pi r h}=r \rho_{s}, \quad \rho_{s}=\text { the mass density of the final sample. }
$$

mass absorption coefficients here for Co $K \propto$ radiation are ${ }^{[11]}$ :

$$
\left(\frac{\mu}{\rho}\right)_{\mathrm{Ni}}=75.1, \quad\left(\frac{\mu}{\rho}\right)_{\mathrm{Fe}}=59.5, \quad\left(\frac{\mu}{\rho}\right)_{\mathrm{Ag}}=332
$$

Also: $\quad \mathrm{r}=0.1 \mathrm{~cm}, \rho_{\mathrm{s}}(\mathrm{Fe}-\mathrm{Ni}$ alloys $)=3.65 \mathrm{~g} / \mathrm{cm}^{3}, \quad \rho_{\mathrm{s}}(\mathrm{Ag})=3.92 \mathrm{~g} / \mathrm{cm}^{3}$.
The total angular correction factor used to multiply the intensity is composed of the preceeding three elements.

$$
\mathrm{c}=\frac{1}{f^{2}(\mathrm{LP}) \mathrm{A}}
$$

This is normalized to unity at $2 \theta_{1}$.
"UNFOLD" has the option to make this correction taking into account the Lorenzpolarization and atomic scattering factors. This is all that is necessary when using a flat sample but the Debye-Scherrer geometry requires the inclusion of the absorption factor.

## 4. Rachinger Separation, "RACH"

The Rachinger method
[4] is a graphical means of separating the $\mathrm{K} \alpha_{1}-\mathrm{K} \alpha_{2}$ line doublet.

$$
\begin{aligned}
& I_{1}\left(2 \theta_{i}\right)=I\left(2 \theta_{i}\right)-I_{2}\left(2 \theta_{i}\right) \text { where } \\
& I=\text { intensity due to } \alpha_{1}+\alpha_{2} \text { (recorded intensity) } \\
& I_{1}=" \quad \text { " } \quad \text { " } \alpha_{1} \text {. } \\
& I_{2}=" \quad " \quad " \alpha_{2} .
\end{aligned}
$$

The usual assumption [ is made that the $K \alpha_{2}$ intensity is half that of the $K \alpha_{1}$ so that:

$$
\mathrm{I}_{2}\left(2 \theta_{\mathrm{i}}\right)=\frac{1}{2} \mathrm{I}_{1}(2 \theta-\mathrm{d}),
$$

where $d$ is the doublet separation. Now

$$
I_{1}\left(2 \theta_{i}\right)=I\left(2 \theta_{i}\right)-\frac{1}{2} I_{1}\left(2 \theta_{i}-d\right) .
$$

It is important that data be collected sufficiently far out on the low angle side of the profile so that the calculation is started at a diffraction angle sufficiently small that $I_{2} \approx 0$, or the measured intensity is due mainly to the $K \alpha_{1}$ component. Values of $I_{1}$ at higher angles can then be determined by successive calculations.

The doublet separation, $d$, is determined from an approximate peak position of the $\mathrm{K} \alpha$ curve and the wavelengths $\lambda, \lambda_{1}$, and $\lambda_{2}$ of the $K \alpha, K \alpha_{1}$, and $K \alpha_{2}$ radiations respectively.

From the Bragg equation:

$$
\dot{\lambda}=2 \mathrm{~d}^{\prime} \sin \theta \text { and } \lambda_{1}=2 d^{\prime} \sin \theta_{1}
$$

( $\mathrm{d}^{\prime}$ = interplanar spacing),

$$
\sin \theta_{1}=\frac{\lambda_{1}}{\lambda} \sin \theta_{0}
$$

Similarly for the $\mathrm{K} \alpha_{2}$ component:

$$
\begin{aligned}
& \sin \theta_{2}=\frac{\lambda_{2}}{\lambda} \sin \theta,
\end{aligned}
$$

and the doublet separation in $2 \theta$ is

$$
d=2 \theta_{2}-2 \theta_{1}
$$

or

$$
d=2\left\{\sin ^{-1}\left[\frac{\lambda_{2}}{\lambda} \sin \theta\right]-\sin ^{-1}\left[\frac{\lambda_{1}}{\lambda} \sin \theta\right]\right\}
$$

For cobalt radiation

$$
\begin{aligned}
& \lambda=1.79021 \AA, \\
& \lambda_{1}=1.78892 \AA . \\
& \lambda_{2}=1.79278 \AA .
\end{aligned}
$$

In general, the angle $\left(2 \theta_{i}-d\right)$ will not coincide with any of the angular values $2 \theta_{j}$ so it is necessary to interpolate. Assume:

$$
2 \theta_{j} \leq 2 \theta_{i}-d \leq 2 \theta_{k} .
$$

A straight line is fitted between $I_{1}\left(2 \theta_{j}\right)$ and $I_{1}\left(2 \theta_{k}\right)$ and:

$$
I_{1}\left(2 \theta_{i}-d\right)=\frac{I_{1}\left(2 \theta_{k}\right)-I_{1}\left(2 \theta_{j}\right)}{2 \theta_{k}-2 \theta_{j}}\left(2 \theta_{i}-d\right)+\frac{\left[I_{1}\left(2 \theta_{j}\right)\right]\left(2 \theta_{k}\right)-\left[I_{l}\left(2 \theta_{k}\right)\right]\left(2 \theta_{j}\right)}{2 \theta_{k}-2 \theta_{j}}
$$

Although it is not apparent, the Rachinger separation has a tendency to oscillate or become assymmetrical on the high angle side of the profile. This is especially true for a sharp, well separated $\mathrm{K} \alpha$ line. For this reason the options for card punching and mathematical smoothing are applied only to the corrected $K \alpha$ profile in this program. This section is still included, however, for those occasions when it might be useful.
5. Center of Gravity Calculation, "CENTROID"

The centroid or center of gravity is one measure of the position of the diffraction line. It is of value in measuring lattice parameters ${ }^{[13]}$, twin fault probabilities ${ }^{[14]}$, etc. According to Pike and Wilson ${ }^{[13]}$ the center of gravity is determined by

$$
\begin{equation*}
G=\frac{\sum_{1}\left(I_{i}-b_{i}\right) c_{i}\left(2 \theta \theta_{i}\right)}{\sum_{i}\left(I_{i}-b_{i}\right) c_{i}} \tag{1}
\end{equation*}
$$

where

$$
\begin{aligned}
& G=\text { center of gravity, } \\
& 2 \theta_{i}=\text { angular position }, \\
& I_{i}=\text { measured intensity }, \\
& b_{i}=\text { background (calculated in sec. } 2 \text { ), } \\
& c_{i}=\text { angular correction factor (calculated in sec. 3). } \\
& \text { The summations are over all points. }
\end{aligned}
$$

Equation (1) is modified from the equation given by Pike and Wilson by the presence of the $c_{i}$

$$
\begin{aligned}
& \text { From the calculation of the background (sec. 2) } \\
& \qquad \text { setting } \Delta 2 \theta_{i}=2 \theta_{n}-2 \theta_{1}, \\
& \qquad b_{i}=\left(\frac{z-a}{\Delta 2 \theta}\right) 2 \theta_{i}+\frac{a\left(2 \theta_{n}\right)-z\left(2 \theta_{1}\right)}{\Delta 2 \theta^{2}} ;
\end{aligned}
$$

or

$$
\begin{equation*}
b_{i}=z\left(\frac{2 \theta_{i}-2 \theta_{1}}{\Delta 2 \theta}\right)-a\left(\frac{2 \theta_{i}-2 \theta_{n}}{\Delta 2 \theta}\right) \tag{2}
\end{equation*}
$$

Substituting eq (2) into eq (1)

$$
G=\frac{\Sigma \sum_{i}\left[I_{i}-z\left(\frac{2 \theta_{i}-2 \theta_{1}}{\Delta 2 \theta}\right)+a\left(\frac{2 \theta_{i}-2 \theta_{n}}{\Delta 2 \theta}\right)\right] c_{i}^{\left(2 \theta_{i}\right)}}{\Sigma_{i}\left[I_{i}-z\left(\frac{2 i_{i}-2 \theta}{\Delta 2 \theta}\right)+a\left(\frac{i_{i}-2 \theta}{\Delta 2 \theta}\right)\right] c_{i}},
$$

$$
\begin{equation*}
\left.G=\frac{\sum_{i} I_{i} c_{i}\left(2 \theta_{i}\right)-z \sum_{i}\left(\frac{2 \theta_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i}\left(2 \theta_{i}\right)+a \sum_{i}\left(\frac{2 \theta_{i}-2 \theta_{n}}{\Delta 2 \theta}\right) c_{i} \cdot\left(2 \theta_{i}\right)}{\sum_{i} I_{i} c_{i}-z \sum_{i}\left(\frac{2 \theta_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i}+a \sum_{i}\left(\frac{2 \theta}{i}-2 \theta n\right.} \frac{\Delta 2 \theta}{\Delta 2 \theta}\right) c_{i} \quad . \tag{3}
\end{equation*}
$$

An estimate of the standard deviation is given by the standard propagation of error.

$$
\begin{equation*}
\sigma^{2}(G)=\sum_{i}\left(\frac{\partial G}{\partial I_{i}}\right)^{2} \sigma^{2}\left(I_{i}\right)+\left(\frac{\partial G}{\partial a}\right)^{2} \sigma^{2}(a)+\left(\frac{\partial G}{\partial z}\right)^{2} \sigma^{2}(z) \tag{4}
\end{equation*}
$$

Now

$$
\begin{aligned}
& \frac{\partial G}{\partial I_{i}}=\frac{\left(2 \theta_{i}\right) c_{i}}{\sum_{i} I_{i} c_{i}-z \sum_{i}\left(\frac{2{ }_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i}+a \sum_{i}\left(\frac{2 \theta_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i}} \\
& \frac{c_{i}\left[\sum_{i} I_{i} c_{i}\left(2 \theta_{i}\right)-z \sum_{i}\left(\frac{2{ }_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i}\left(2 \theta_{i}\right)+a \sum_{i}\left(\frac{2 \theta_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i} \cdot\left(2 \theta_{i}\right)\right]}{\left[\sum \sum_{i} I_{i} c_{i}-z \sum_{i}\left(\frac{2 i_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i}+a \sum_{i} c_{i}\left(\frac{2 \theta_{i}-2 \theta n}{\Delta 2 \theta}\right)\right]^{2}} .
\end{aligned}
$$

On substituting from eq (3)

$$
\begin{equation*}
\left.\frac{\partial G}{\partial I_{i}}=\frac{\left(2 \theta i_{i}-G\right) c_{i}}{\sum_{i} I_{i} c_{i}-z \sum_{i}\left(\frac{2 \theta i}{\Delta 2 \theta}-2 \theta\right.}\right) c_{i}+a \sum_{i}\left(\frac{2 \theta i}{\Delta 2 \theta} n{ }_{i}\right) c_{i} \quad . \tag{5}
\end{equation*}
$$

Also

$$
\frac{\partial G}{\partial a}=\frac{\sum_{i}\left(\frac{2 \theta_{i}-2 \theta_{n}}{\Delta 2 \theta}\right) c_{i} \cdot\left(2 \theta_{i}\right)}{\sum_{i} I_{i} c_{i}-z \sum_{i}\left(\frac{2 \theta_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i}+a \sum_{i}\left(\frac{2 \theta_{i}-2 \theta_{n}}{\Delta 2 \theta}\right) c_{i}}
$$

$$
-\frac{\sum_{i}\left(\frac{2 \theta_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i}\left[\sum_{i} I_{i} c_{i} \cdot\left(2 \theta_{i}\right)-z \sum_{i}\left(\frac{2 \theta_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i} \cdot\left(2 \theta_{i}\right)+a \sum_{i}\left(\frac{2 \theta_{i}-2 \theta_{n}}{\Delta 2 \theta}\right) c_{i} \cdot\left(2 \theta_{i}\right)\right]}{\left[\sum_{i} I_{i} c_{i}-z \sum_{i}\left(\frac{2 \theta_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i}+a \sum_{i}\left(\frac{2 \theta i_{i}-2 \theta{ }_{n}}{\Delta 2 \theta}\right) c_{i}\right]^{2}} .
$$

On substituting from eq (3)

$$
\frac{\partial G}{\partial a}=\frac{\sum_{i}\left(\frac{2 \theta_{i}-2 \theta n}{\Delta 2 \theta}\right) c_{i} \cdot\left(2 \theta_{i}\right)-G \sum_{i}\left(\frac{2 \theta_{i}-2 \theta n}{\Delta 2 \theta}\right) c_{i},}{\sum_{i} I_{i} c_{i}-z \sum_{i}\left(\frac{2 \theta_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i}+a \sum\left(\frac{\sum_{i}-2 \theta n}{\Delta 2 \theta}\right) c_{i}}
$$

Or

$$
\begin{equation*}
\frac{\partial G}{\partial a}=\frac{\sum_{i}\left(\frac{2 \theta_{i}-2 \theta_{n}}{\Delta 2 \theta}\right) c_{i} \cdot\left(2 \theta_{i}-G\right)}{\sum_{i} I_{i} c_{i}-z \sum_{i}\left(\frac{{ }_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i}+a \sum_{i}\left(\frac{2 i_{i}-2 \theta n}{\Delta 2 \theta}\right) c_{i}} . \tag{6}
\end{equation*}
$$

Also:

$$
\begin{aligned}
& \frac{\partial G}{\partial z}=\frac{-\sum_{i}\left(\frac{2 i_{i}-2 \theta 1}{\Delta 2 \theta}\right) c_{i} \cdot\left(2 \theta_{i}\right)}{\sum_{i} I_{i} c_{i}-z \sum_{i}\left(\frac{2 \theta i^{-2 \theta} 1}{\Delta 2 \theta}\right) c_{i}+a \sum_{i}\left(\frac{2 \theta i^{-2 \theta} n}{\Delta 2 \theta}\right) c_{i}} \\
& +\frac{\sum_{i}\left(\frac{2 \theta_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i} \cdot\left[\sum_{i} I_{i} c_{i} \cdot\left(2 \theta_{i}\right)-z \sum_{i}\left(\frac{2 \theta_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i} \cdot\left(2 \theta_{i}\right)+a \sum_{i}\left(\frac{2 \theta_{i}-2 \theta_{n}}{\Delta 2 \theta}\right) c_{i} \cdot\left(2 \theta_{i}\right)\right]}{\left[\sum I_{i} c_{i}-z \sum_{i}\left(\frac{2 \theta_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i}+a \sum_{i}\left(\frac{2 \theta{ }_{i}-2 \theta_{n}}{\Delta 2 \theta}\right) c_{i}\right]^{2}},
\end{aligned}
$$

or

$$
\begin{equation*}
\frac{\partial G}{\partial z}=\frac{\sum_{i}\left(\frac{2 \theta_{i}-2 \theta}{\Delta 2 \theta}\right) c_{i} \cdot\left(G-2 \theta_{i}\right)}{\sum_{i} I_{i} c_{i}-z \sum_{i}\left(\frac{2 \theta i^{-2 \theta} 1}{\Delta 2 \theta}\right) c_{i}+a \sum_{i}\left(\frac{2 \theta i^{-2 \theta} n}{\Delta 2 \theta}\right) c_{i}} . \tag{7}
\end{equation*}
$$

And, since the standard deviation of a random variable is the square root of that variable,

$$
\begin{equation*}
\sigma^{2}\left(I_{i}\right)=I_{i} \tag{8}
\end{equation*}
$$

also $a=\frac{1}{5} \sum_{i=1}^{5} I_{i} \quad$ (see sec. 2),

$$
\sigma^{2}(\mathrm{a})=\sum_{\mathrm{i}}\left(\frac{\partial \mathrm{a}}{\partial \mathrm{I}_{\mathrm{i}}}\right)^{2} \sigma^{2}\left(\mathrm{I}_{\mathrm{i}}\right) \quad \text { where } \quad \frac{\partial \mathrm{a}}{\partial \mathrm{I}_{\mathrm{i}}}=\frac{1}{5}
$$

so $\quad \sigma^{2}(a)=\frac{1}{25} \sum_{i=1}^{5} I_{i}=\frac{1}{5} a$.
In a similar manner for $z$,

$$
\begin{equation*}
\sigma^{2}(z)=\frac{1}{5} z \tag{10}
\end{equation*}
$$

The center of gravity is calculated from eq (1) and the estimate of its standard deviation is determined from eq (4) using eq (5) through eq (10).

## 6. Peak Calculation, "PEAK"

The position of peak maximum is another measure of the profile's angular loca-
tion. It also is used to measure lattice parameter and twin fault probabilities, as well as stacking fault probabilities ${ }^{[15]}$.

A portion of the profile on either side of the maximum (we generally used a total of 7 to 9 points) is fitted to a parabola ${ }^{[16]}$ and the maximum of that curve is determined.

$$
I_{i}=B+C \cdot\left(2 \theta_{i}\right)+D \cdot\left(2 \theta_{i}\right)^{2}
$$

Differentiating:

$$
\frac{d I_{i}}{d\left(2 \theta_{i}\right)}=C+2 D \cdot\left(2 \theta_{i}\right)=0
$$

so that

$$
\begin{equation*}
\text { peak }=2 \theta_{\max }=-\frac{C}{2 D} \tag{11}
\end{equation*}
$$

where $B, C$, and $D$ are the parameters of the fitted curve.
The standard propagation of error with correlation between the variables $C$ and $D$ gives:

$$
\begin{aligned}
\sigma^{2}\left(2 \theta_{\max }\right)= & \left(\frac{\partial\left(2 \theta_{\text {max }}\right)}{\partial \mathrm{C}}\right)^{2} \sigma^{2}(\mathrm{C})+\left(\frac{\partial\left(2 \theta_{\text {max }}\right)}{\partial \mathrm{D}}\right)^{2} \sigma^{2}(\mathrm{D}) \\
& +2\left(\frac{\partial\left(2 \theta_{\text {max }}\right)}{\partial \mathrm{C}}\right)\left(\frac{\partial\left(2 \theta_{\text {max }}\right)}{\partial \mathrm{D}}\right) \sigma(\mathrm{C}) \sigma(\mathrm{D}) .
\end{aligned}
$$

Differentiating eq (ll) this becomes

$$
\begin{aligned}
\sigma^{2}\left(2 \theta_{\max }\right) & =\left(\frac{1}{4 D^{2}}\right) \sigma^{2}(C)+\left(\frac{C^{2}}{4 D^{4}}\right) \sigma^{2}(D)+2\left(\frac{-1}{2 D}\right)\left(\frac{C}{2 D^{2}}\right) \sigma(C) \sigma(D) \\
& =\frac{\left(2 \theta \max ^{\prime}\right)^{2}}{C^{2}} \sigma^{2}(C)+\frac{\left(2 \theta \max ^{2}\right)^{2}}{D^{2}} \sigma^{2}(D)-2 \frac{\left(2 \theta \max ^{2}\right.}{C D} \sigma(C) \sigma(D) .
\end{aligned}
$$

A modification of a computer program written by W. J. Hall was used here to compute the curve parameters and their standard deviations. This program performs an orthonormal curve fitting using Bjork's modification ${ }^{[17,18]}$ of the Gram-Schmidt technique.

## 7. Curve Smoothing, "SMOOTH"

In the course of experimenting on several iron-nickel alloys, it was found that the Fourier analyses of the diffraction profiles were far more repeatable if the data from the cold-worked sample were first smoothed. This subroutine affords the option of performing this operation.

The intensity curve is mathematically smoothed by a least-squares fit of a quartic equation to seven points, taking the value at the midpoint and then advancing the group of points by one unit and repeating the process. [19]

$$
I_{i}(\text { smoothed })=\left(I_{i-3}-18 I_{i-2}+63 I_{i-1}+164 I_{i}+63 I_{i+1}-18 I_{i+2}+I_{i+3}\right) / 256
$$

Three points are added before the initial intensity and three after the final intensity and they are all set equal to zero.

This smoothing operation can be performed repeatedly (limited to 99 times by format). It is possible to have up to ten of these smoothed curves printed and to have one of them card punched for input to "UNFOLD".
8. Data Input and Output and a Sample

Card Input

Card
1

2

3

5
6.

Smoothing numbers to be printed. (This card is 1015 omitted if $3 a \neq 2$. The number in $3 c$ must be one of these numbers).

Lower and upper angular limits for parabolic
2F10. 2 fit to determine profile maximum.

Intensities
10F7. 0

Output data includes:

## 1. Identification

2. Table containing angular positions, recorded intensities, backgrounds, angular corrections, and the corrected $\mathrm{K} \alpha$ and $\mathrm{K} \alpha_{1}$ profiles.

[^1]3. Line maximum and center of gravity with their standard deviations for $\mathrm{K} \alpha$ and $K \alpha_{1}$ profiles.
4. Peak value used for punching "UNFOLD" input and smoothing number used (if applicable).
5. Table containing angular positions and smoothed $\mathrm{K} \alpha$ profile (if applicable). A sample of input and output follows.

## Punched Card Input



Printed Output

FE-28NI COLD-WORKED (111)

| 2 THETA | REC $C N T$ | BACK | ANG COR | ALPHA | ALPHA I |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: |
|  |  |  |  |  |  |
| 50.60 | 270 | 264 | 1.0000 | 6 | 0 |
| 50.62 | 260 | 264 | 1.0004 | -4 | 0 |
| 50.64 | 260 | 264 | 1.0009 | -4 | 0 |
| 50.66 | 270 | 264 | 1.0013 | 6 | 0 |
| 50.68 | 260 | 264 | 1.0017 | -4 | 0 |
| 50.70 | 290 | 264 | 1.0022 | 26 | 0 |
| 50.72 | 240 | 264 | 1.0026 | -24 | 0 |
| 50.74 | 290 | 264 | 1.0030 | 26 | 0 |
| 50.76 | 310 | 264 | 1.0035 | 46 | 46 |
| 50.78 | 310 | 264 | 1.0039 | 46 | 46 |
| 50.80 | 290 | 265 | 1.0044 | 26 | 26 |
| 50.82 | 290 | 265 | 1.0048 | 26 | 26 |
| 50.84 | 310 | 265 | 1.0052 | 46 | 46 |
| 50.86 | 300 | 265 | 1.0057 | 35 | 35 |
| 50.88 | 270 | 265 | 1.0061 | 5 | -17 |
| 50.90 | 260 | 265 | 1.0065 | -5 | -28 |
| 50.92 | 320 | 265 | 1.0070 | 56 | 43 |
| 50.94 | 280 | 265 | 1.0074 | 15 | 2 |
| 50.96 | 330 | 265 | 1.0078 | 66 | 43 |
| 50.98 | 330 | 265 | 1.0083 | 66 | 48 |
| 51.00 | 300 | 265 | 1.0087 | 35 | 43 |
| 51.02 | 330 | 265 | 1.0092 | 65 | 79 |
| 51.04 | 360 | 265 | 1.0096 | 96 | 75 |


| 51.06 | 340 | 265 | 1.0100 | 75 | 74 |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 51.08 | 370 | 265 | 1.0105 | 106 | 85 |
| 51.10 | 330 | 265 | 1.0109 | 65 | 42 |
| 51.12 | 350 | 265 | 1.0113 | 86 | 64 |
| 51.14 | 360 | 265 | 1.0118 | 96 | 56 |
| 51.16 | 340 | 266 | 1.0122 | 75 | 38 |
| 51.18 | 350 | 266 | 1.0127 | 85 | 48 |
| 51.20 | 390 | 266 | 1.0131 | 126 | 84 |
| 51.22 | 400 | 266 | 1.0135 | 136 | 115 |
| 51.24 | 400 | 266 | 1.0140 | 136 | 104 |
| 51.26 | 370 | 266 | 1.0144 | 106 | 77 |
| 51.28 | 420 | 266 | 1.0148 | 156 | 137 |
| 51.30 | 450 | 266 | 1.0153 | 187 | 163 |
| 51.32 | 490 | 266 | 1.0157 | 228 | 186 |
| 51.34 | 470 | 266 | 1.0162 | 207 | 150 |
| 51.36 | 510 | 266 | 1.0166 | 248 | 196 |
| 51.38 | 510 | 266 | 1.0170 | 248 | 209 |
| 51.40 | 530 | 266 | 1.0175 | 268 | 200 |
| 51.42 | 600 | 266 | 1.0179 | 340 | 259 |
| 51.44 | 620 | 266 | 1.0184 | 360 | 267 |
| 51.46 | 710 | 266 | 1.0188 | 452 | 377 |
| 51.48 | 740 | 266 | 1.0192 | 483 | 385 |
| 51.50 | 740 | 266 | 1.0197 | 483 | 378 |
| 51.52 | 830 | 267 | 1.0201 | 575 | 475 |
| 51.54 | 910 | 267 | 1.0206 | 657 | 528 |
| 51.56 | 950 | 267 | 1.0210 | 698 | 564 |
| 51.58 | 940 | 267 | 1.0214 | 688 | 500 |
| 51.60 | 1010 | 267 | 1.0219 | 760 | 567 |
| 51.62 | 1170 | 267 | 1.0223 | 923 | 734 |
| 51.64 | 1260 | 267 | 1.0228 | 1016 | 779 |
| 51.66 | 1330 | 267 | 1.0232 | 1088 | 824 |
| 51.68 | 1410 | 267 | 1.0236 | 1170 | 888 |
| 51.70 | 1550 | 267 | 1.0241 | 1314 | 1063 |
| 51.72 | 1620 | 267 | 1.0245 | 1386 | 1103 |
| 51.74 | 1720 | 267 | 100250 | 1489 | 1123 |
| 51.76 | 1840 | 267 | 1.0254 | 1613 | 1224 |
| 51.78 | 1940 | 267 | 1.0258 | 1716 | 1304 |
| 51.80 | 2050 | 267 | 1.0263 | 1830 | 1386 |
| 51.82 | 2130 | 267 | 1.0267 | 1912 | 1382 |
| 51.84 | 2220 | 267 | 1.0272 | 2006 | 1454 |
| 51.86 | 2270 | 267 | 1.0276 | 2058 | 1496 |
| 51.88 | 2350 | 267 | 1.0280 | 2141 | 1530 |
| 51.90 | 2380 | 268 | 1.0285 | 2173 | 1521 |
| 51.92 | 2390 | 268 | 1.0289 | 2184 | 1491 |
| 51.94 | 2390 | 268 | 1.0294 | 2185 | 1494 |
| 51.96 | 2360 | 268 | 1.0298 | 2155 | 1428 |
| 51.98 | 2320 | 268 | 1.0303 | 2114 | 1366 |
| 52.00 | 2280 | 268 | 1.0307 | 2074 | 1309 |
| 52.02 | 2190 | 268 | 1.0311 | 1982 | 1221 |
| 52.04 | 2040 | 268 | 1.0316 | 1828 | 1082 |
| 52.06 | 1940 | 268 | 1.0320 | 1726 | 979 |
| 52.08 | 1840 | 268 | 1.0325 | 1623 | 908 |
| 52.10 | 1600 | 268 | 1.0329 | 1376 | 692 |
| 52.12 | 1330 | 268 | 1.0334 | 1097 | 442 |
| 52.14 | 1160 | 268 | 1.0338 | 922 | 311 |
| 52.16 | 1050 | 268 | 1.0342 | 809 | 266 |
| 52.18 | 950 | 268 | 1.0347 | 705 | 215 |
| 52.20 | 860 | 268 | 1.0351 | 612 | 158 |
| 52.22 | 810 | 268 | 1.0356 | 561 | 213 |
| 52.24 | 760 | 268 | 1.0360 | 509 | 286 |
| 52.26 | 720 | 269 | 1.0365 | 468 | 312 |
| 52.28 | 670 | 269 | 1.0369 | 416 | 283 |
| 52.30 | 610 | 269 | 1.0374 | 354 | 246 |
| 52.32 | 560 | 269 | 1.0378 | 302 | 223 |
|  |  |  | 18 |  |  |


| 52.34 | 510 | 269 | 1.0382 | 250 | 144 |
| ---: | ---: | ---: | ---: | ---: | ---: |
| 52.36 | 490 | 269 | 1.0387 | 230 | 87 |
| 52.38 | 470 | 269 | 1.0391 | 209 | 53 |
| 52.40 | 440 | 269 | 1.0396 | 178 | 36 |
| 52.42 | 420 | 269 | 1.0400 | 157 | 34 |
| 52.44 | 400 | 269 | 1.0405 | 136 | 25 |
| 52.46 | 390 | 269 | 1.0409 | 126 | 53 |
| 52.48 | 380 | 269 | 1.0414 | 115 | 71 |
| 52.50 | 370 | 269 | 1.0418 | 105 | 78 |
| 52.52 | 360 | 269 | 1.0422 | 95 | 76 |
| 52.54 | 350 | 269 | 1.0427 | 84 | 67 |
| 52.56 | 340 | 269 | 1.0431 | 74 | 61 |
| 52.58 | 330 | 269 | 1.0436 | 63 | 37 |
| 52.60 | 320 | 269 | 1.0440 | 53 | 17 |
| 52.62 | 310 | 270 | 1.0445 | 42 | 3 |
| 52.64 | 300 | 270 | 1.0449 | 32 | -6 |
| 52.66 | 290 | 270 | 1.0454 | 21 | -12 |
| 52.68 | 280 | 270 | 1.0458 | 11 | -20 |
| 52.70 | 270 | 270 | 1.0463 | 0 | -18 |
| 52.72 | 270 | 270 | 1.0467 | 0 | -9 |
| 52.74 | 270 | 270 | 1.0471 | 0 | -2 |
| 52.76 | 270 | 270 | 1.0476 | 0 | 3 |
| 52.78 | 270 | 270 | 1.0480 | 0 | 6 |
| 52.80 | 270 | 270 | 1.0485 | 0 | 10 |

```
ALPHA CURVE
PEAK MAX = 51.9298
    STD DEV = 0.1335
PEAK FIGURED FROM 51.72 TO 52.02
LINE CG = 51.8653 STD DEV = 0.0045
ALPHA 1 CURVE
PEAK MAX = 51.8958
    STD DEV = 0.2744
PEAK FIGURED FROM 51.68 TO 51.98
LINE CG = 51.8257 STD DEV = 0.0065
PUNCH PEAK FOR ALPHA = 51.92
SMOOTHING }
```

| ANGLE | INPUT | SMOOTH | 2 | SMOOTH | 4 | SMOOTH | 6 |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| 50.60 | 6 | 2 | 0 | -0 | SMOOTH | 8 |  |
| 50.62 | -4 | -1 | -1 | -1 | -0 |  |  |
| 50.64 | -4 | -3 | -2 | -1 | -1 |  |  |
| 50.66 | 6 | 1 | 2 | 2 | -1 |  |  |
| 50.68 | -4 | 7 | 5 | 4 | 2 |  |  |
| 50.70 | 26 | 4 | 2 | 2 | 3 |  |  |
| 50.72 | -24 | -1 | 3 | 5 | 2 |  |  |
| 50.74 | 26 | 17 | 19 | 20 | 7 |  |  |
| 50.76 | 46 | 43 | 39 | 36 | 20 |  |  |
| 50.78 | 46 | 43 | 41 | 40 | 35 |  |  |
| 50.80 | 26 | 29 | 33 | 35 | 39 |  |  |
| 50.82 | 26 | 30 | 32 | 34 | 36 |  |  |
| 50.84 | 46 | 40 | 37 | 34 | 35 |  |  |
| 50.86 | 35 | 31 | 28 | 26 | 33 |  |  |
|  |  |  |  |  | 25 |  |  |


| 50.88 | 5 | 8 | 12 | 14 | 16 |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 50.90 | -5 | 10 | 12 | 13 | 14 |
| 50.92 | 56 | 29 | 26 | 25 | 25 |
| 50.94 | 15 | 40 | 42 | 41 | 41 |
| 50.96 | 66 | 54 | 53 | 52 | 51 |
| 50.98 | 66 | 57 | 55 | 54 | 54 |
| 51.00 | 35 | 50 | 54 | 56 | 57 |
| 51.02 | 65 | 64 | 65 | 66 | 67 |
| 51.04 | 96 | 85 | 82 | 81 | 80 |
| 51.06 | 75 | 91 | 91 | 89 | 88 |
| 51.08 | 106 | 88 | 87 | 88 | 88 |
| 51.10 | 65 | 80 | 82 | 84 | 84 |
| 51.12 | 86 | 83 | 83 | 83 | 82 |
| 51.14 | 96 | 87 | 84 | 82 | 82 |
| 51.16 | 75 | 81 | 83 | 84 | 85 |
| 51.18 | 85 | 91 | 94 | 97 | 98 |
| 51.20 | 126 | 120 | 118 | 117 | 116 |
| 51.22 | 136 | 137 | 132 | 129 | 127 |
| 51.24 | 136 | 128 | 128 | 128 | 128 |
| 51.26 | 106 | 122 | 126 | 130 | 152 |
| 51.28 | 156 | 149 | 150 | 152 | 153 |
| 51.30 | 187 | 191 | 187 | 185 | 184 |
| 51.32 | 228 | 215 | 213 | 212 | 210 |
| 51.34 | 207 | 224 | 225 | 226 | 225 |
| 51.36 | 248 | 236 | 235 | 235 | 236 |
| 51.38 | 248 | 251 | 251 | 251 | 251 |
| 51.40 | 268 | 278 | 279 | 280 | 281 |
| 51.42 | 340 | 323 | 324 | 325 | 326 |
| 51.44 | 360 | 380 | 381 | 381 | 380 |
| 51.46 | 452 | 439 | 435 | 432 | 429 |
| 51.48 | 483 | 474 | 471 | 471 | 471 |
| 51.50 | 483 | 502 | 510 | 514 | 517 |
| 51.52 | 575 | 571 | 574 | 576 | 576 |
| 51.54 | 657 | 652 | 643 | 637 | 633 |
| 51.56 | 698 | 685 | 680 | 677 | 675 |
| 51.58 | 688 | 700 | 708 | 714 | 718 |
| 51.60 | 760 | 775 | 782 | 787 | 790 |
| 51.62 | 923 | 905 | 899 | 895 | 893 |
| 51.64 | 1016 | 1014 | 1008 | 1004 | 1001 |
| 51.66 | 1088 | 1090 | 1095 | 1097 | 1097 |
| 51.68 | 1170 | 1184 | 1187 | 1189 | 1190 |
| 51.70 | 1314 | 1295 | 1291 | 1290 | 1289 |
| 51.72 | 1386 | 1394 | 1394 | 1393 | 1392 |
| 51.74 | 1489 | 1493 | 1496 | 1497 | 1498 |
| 51.76 | 1613 | 1607 | 1607 | 1607 | 1608 |
| 51.78 | 1716 | 1721 | 1720 | 1719 | 1719 |
| 51.80 | 1830 | 1825 | 1825 | 1824 | 1824 |
| 51.82 | 1912 | 1918 | 1917 | 1917 | 1917 |
| 51.84 | 2006 | 1998 | 1998 | 1999 | 1999 |
| 51.86 | 2058 | 2069 | 2071 | 2071 | 2072 |
| 51.88 | 2141 | 2133 | 2132 | 2131 | 2131 |
| 51.90 | 2173 | 2173 | 2172 | 2171 | 2171 |
| 51.92 | 2184 | 2187 | 2187 | 2187 | 2187 |
| 51.94 | 2185 | 2181 | 2180 | 2181 | 2181 |
| 51.96 | 2155 | 2156 | 2157 | 2159 | 2160 |
| 51.98 | 2114 | 2120 | 2123 | 2123 | 2122 |
| 52.00 | 2074 | 2069 | 2064 | 2061 | 2059 |
| 52.02 | 1982 | 1971 | 1969 | 1968 | 1969 |
| 52.04 | 1828 | 1844 | 1852 | 1858 | 1861 |
| 52.06 | 1726 | 1734 | 1736 | 1736 | 1736 |
| 52.08 | 1623 | 1599 | 1587 | 1579 | 1574 |
| 52.10 | 1376 | 1373 | 1369 | 1365 | 1363 |
| 52.12 | 1097 | 1114 | 1124 | 1130 | 1134 |
| 52.14 | 922 | 924 | 929 | 933 | 938 |


| 52.16 | 809 | 803 | 799 | 797 | 797 |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $52 \cdot 18$ | 705 | 703 | 701 | 699 | 698 |
| 52.20 | 612 | 618 | 620 | 621 | 621 |
| 52.22 | 561 | 557 | 558 | 559 | 560 |
| 52.24 | 509 | 511 | 511 | 510 | 510 |
| 52.26 | 468 | 466 | 465 | 465 | 464 |
| 52.28 | 416 | 415 | 414 | 413 | 412 |
| 52.30 | 354 | 356 | 356 | 356 | 356 |
| $52 \cdot 32$ | 302 | 300 | 300 | 301 | 302 |
| 52.34 | 250 | 255 | 257 | 258 | 259 |
| $52 \cdot 36$ | 230 | 228 | 228 | 228 | 228 |
| 52.38 | 209 | 206 | 205 | 204 | 203 |
| 52.40 | 178 | 180 | 180 | 180 | 180 |
| 52.42 | 157 | 156 | 156 | 157 | 157 |
| 52.44 | 136 | 138 | 138 | 138 | 139 |
| 52.46 | 126 | 125 | 125 | 125 | 125 |
| 52.48 | 115 | 115 | 115 | 115 | 115 |
| 52.50 | 105 | 105 | 105 | 105 | 105 |
| 52.52 | 95 | 95 | 95 | 95 | 95 |
| 52.54 | 84 | 84 | 84 | 84 | 84 |
| 52.56 | 74 | 74 | 74 | 74 | 74 |
| 52.58 | 63 | 63 | 63 | 63 | 63 |
| 52.60 | 53 | 53 | 53 | 53 | 53 |
| 52.62 | 42 | 42 | 42 | 42 | 42 |
| 52.64 | 32 | 32 | 32 | 32 | 32 |
| 52.66 | 21 | 21 | 21 | 21 | 21 |
| 52.68 | 11 | 10 | 10 | 10 | 10 |
| 52.70 | 0 | 2 | 2 | 3 | 3 |
| 52.72 | 0 | -0 | -0 | -0 | -0 |
| 52.74 | 0 | -0 | -0 | -1 | -1 |
| 52.76 | 0 | 0 | 0 | -0 | -0 |
| 52.78 | 0 | 0 | 0 | 0 | 0 |
| $52 \cdot 80$ | 0 | 0 | 0 | 0 | 0 |

## Punched Card Output For "UNFOLD"

| 2187. | 2180. | 2157. | 2123. | 2064. | 1969. | 1852. | 1736. | 1587. | 1369. |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1124. | 929. | 799. | 701. | 620. | 558. | 511. | 465. | 414. | 356. |
| 300. | 257. | 228. | 205. | 180. | 156. | 138. | 125. | 115. | 105. |
| 95. | 84. | 74. | 63. | 53. | 42. | 32. | 21. | 10. | 2. |
| -0. | -0. | 0 . | 0. | 0. |  |  |  |  |  |
| 2187. | 2172. | 2132. | 2071. | 1998. | 1917. | 1825. | 1720. | 1607. | 1496 |
| 1394. | 1291. | 1187. | 1095. | 1008. | 899. | 782. | 708. | 680. | 643. |
| 574. | 510. | 471. | 435. | 381. | 324. | 279 。 | 251. | 235. | 225. |
| 213. | 187. | 150. | 126. | 128. | 132. | 118. | 94. | 83. | 84. |
| 83. | 82. | 87. | 91. | 82. | 65. | 54. | 55. | 53. | 42. |
| 26. | 12. | 12. | 28. | 37. | 32. | 33. | 41. | 39. | 19. |
| 3. | 2. | 5. | 2. | -2. | -1. | 0 。 |  |  |  |

The listing which follows includes two versions on the subroutine "ANG", one for the two phase $\mathrm{Fe}-\mathrm{Ni}$ alloys and one for Ag . Of course, only one version should appear in the program deck. The program has been written in FORTRAN.

PROGRAM PREP
DIMENSION $Y(300), Y 1(300), Y C(300), N S(10)$
COMMON $X(300), B(300), C(300), N$
DOUBLE PRECISION PK,PKI,CG,CGI,SDPK,SDPKI,SDCG,SDCGI
READ 6
IF (EOF,60) 999,100
READ 2, N,AI,AINC,APK,V,FN,IG
READ 1, IPUN,NP,NUMP
IF (IPUN-2) 121,120,121
READ 1,(NS(I),I=1,NP)
DO $101 \mathrm{~J}=1, \mathrm{~N}$
$X(J)=A I+(J-1) * A I N C$
READ 3, R1,R2
READ 4, $(Y(J), J=1, N)$
CALL BACK $(Y, A, Z, D)$
CALL ANG(IG,V,FN)
CALL RACH(Y,Y1,APK,A1)
CALL CENTROID(Y,CG,SDCG,A,Z,D)
CALL CENTROID(YI,CGI,SDCGI,A,Z,D)
CALL PEAK (Y,PK,SDPK,R1,R2)
NN = XFIXF((PK - Al*360.13.1415927)/AINC)
R3 $=$ R1-NN*AINC
$R 4=R 2-N N * A I N C$
CALL PEAK (Y1,PK1,SDPK1,R3,R4)
DO $103 \mathrm{~J}=1 \mathrm{~N}$
IF (APK-X(J))103,102,103
$102 \quad \mathrm{NF}=\mathrm{J}$
GO TO 104
103 CONTINUE
104 DO $105 \mathrm{~J}=1, \mathrm{~N}$
$Y 1(J)=(Y 1(J)-B(J)) * C(J)$
$105 \quad Y C(J)=(Y(J)-B(J)) * C(J)$
PRINT 14
PRINT 6
PRINT $7,(X(J), Y(J), B(J), C(J), Y((J), Y 1(J), J=1, N) ~$
PRINT $8,(X(J)$
PRINT 9, PK,SDPK
PRINT 10, R1,R2
PRINT 11, CG,SDCG
PRINT 12. PKI, SDPK1
PRINT 10, R3,R4
PRINT 11,CGI,SDCGI
IF (IPUN) 107,107,108
IF (IPUN-2) 110,109,109
109 CALL SMOOTH(YC,NUMP,NS,NP,NCEN)
GO TO 107
$110 \quad P P=A I+(N F-1) * A I N C$
PRINT 13, PP
DO $106 \mathrm{~K}=1$, NF
$C(K)=Y C(N F+1-K)$
PUNCH 5, (YC(J), J=NF,N)
PUNCH 5, (C(J),J=1,NF)
107 GO TO 111
1 FORMAT (10I5)
FORMAT (I10,5F10.2,I10)
FORMAT (2F10.2)
FORMAT (1OF7.0)
FORMAT (10(F6.0.\$.\$))
FORMAT 150 H
FORMAT (///3X,\$2 THETA\$, $3 X, \$ R E C$ CNT $\$, 6 X, \$ B A C K \$, 3 X, \$ A N G$ )
1 \$ALPHA\$, $3 X$, \$ALPHA $1 \$ / 1$
FORMAT (F10.2,2F10.0,F10.4,2F10.0)
FORMAT (///\$ ALPHA CURVE $\$ / \$$ PEAK MAX $=\$, F 9.4 / \$$ STD DEV $=\$$,
1F9.4)
FORMAT (1X, \$PEAK FIGURED FROM $\$, F 7.2,2 X, \$$ TO\$,F7.2)
FORMAT ( $/ \$$ LINE CG $=\$$ F9.4. $\$$ STD DEV $=\$, F 9.4$ )
FORMAT $(/ / / \$$ ALPHA 1 CURVE $\$ / \$$ PEAK MAX $=\$, F 9.4 / \$$ STD DEV $=\$$,
1F9.4)
FORMAT (///1X,\$PUNCH PEAK FOR ALPHA $=\$, F 7.2$ )
FORMAT (//////)
END

## SUBROUTINE BACK $(Y, A, Z, D)$

```
    COMMON X(300), B(300),C(300):N
    DIMENSION Y(300)
    A = (Y(1)+Y(2)+Y(3)+Y(4)+Y(5))/5.
    Z = YY(N-4)+Y(N-3)+Y(N-2)+Y(N-1)+Y(N))/5.
    D = X(N) - X(1)
    DO 200 J=1,N
    B(J)=(Z-A)*X(J)/D+(A*X(N)-Z*X(1))/D
    RETURN
    END
```


## SUBROUTINE ANG(IG,V,FN)

ANGULAR CORRECTION - FOR FE-NI ALLOY SAMPLES AND CO RADIATION COMMON X(300), $B(300), C(300), N$
$P I=3.1415927$
$D A=\cos (31.2333 * P I / 180) * *$.
DO $313 \mathrm{~J}=1, \mathrm{~N}$
A
SL $=\operatorname{SIN}((X(J)+V) * P I / 360) /$.
IF (IG-1) 306,306,301
IF (SL-0.12) 302,302,303
$F=1 .-0.142 * S L-6.44 * S L * * 2$
GO TO 311
C-FCC
IF (SL-0.35) 304,304,305
$F=1.0645004-1.4746725 * S L+0.68002691 * S L * * 2-0.53337307 * S L * * 3 \quad$ D-FCC
GO TO 311
$\mathrm{F}=1.1781009-2.1006716 * S L+1.4750092 * \mathrm{SL} * * 2-0.28333892 * \mathrm{SL} * * 3$
GO TO 311
E-FCC
IF (SL-0.12) 307,307,308
$F=1 .-0.184 * S L-7.08 * S L * * 2$
GO TO 311
$C-B C C$
IF (SL-0.35) 309,309,310
$\mathrm{F}=1.0637002-1.6056701 * S L+0.86001593 * S L * * 2-0.53335684 * S L * * 3$
D-BCC
GO TO 311
$F=1.1987014-2.2833411 * S L+2.0200144 * S L * * 2-0.56667538 * S L * * 3 \quad E-B C C$

```
    BB=}\operatorname{cos((X(J)+V)*PI/360.)
    CC= SIN((X(J)+V)*PI/180.)
    D= Cos((X(J)+V)*PI/180.)
    C(J)=(A**2)*BB/(F*F*(10+DA*D**2))
    E=BB*BB*LOGF ((1.+A)/BB)/A
        H
    RM = (75.1*FN + 59.5*(1.-FN))*0.365
    ABS = (I\bullet-E)/(PI*RM)+CC/(2.*PI*RM*RM)+(-. 25+.375*BB*BB*(1\bullet+E))/
    I
1(PI*RM**3)
SUBROUTINE ANG(IG,V,FN)
ANGULAR CORRECTION - FOR AG SAMPLES AND CO RADIATION
COMMON \(\mathrm{X}(200), \mathrm{B}(200), C(200), N\)
\(P I=3.1415927\)
\(D A=\operatorname{COS}(31.2333 * P I / 180) * *\).
DO \(313 \mathrm{~J}=1, \mathrm{~N}\)
A
\(S L=S I N((X(J)+V) * P I / 360) /\).1.79021 .
IF (SL-0.12) 307,307,308
\(F=1 .-0.068 * S L-5.96 * S L * * 2\)
GO TO 311
IF (SL-0.35) \(309,309,310\)
\(\mathrm{F}=1.0685171-1.2033810 * S L-1.2685714 * S L * * 2+3.0666667 * S L * * 3\)
GO TO 311
\(F=1.2061-2.4973333 * S L+2.805 * S L * * 2-1.2166667 * S L * * 3\)
\(B B=\cos ((X(J)+V) * P I / 360\).
\(C C=\operatorname{SIN}((X(J)+V) * P I / 180\).
\(D=\cos ((X(J)+V) * P I / 180\).
\(C(J)=\left(A^{* *} 2\right) * B B /\left(F * F^{*}\left(1 \bullet+D A * D^{*} * 2\right)\right)\)
\(E=B B * B B * \operatorname{LOGF}((1 \bullet+A) / B B) / A\)
\(A B S=(1,-E) /(P I * R M)+C C /(2 \cdot * P I * R M * R M)+(-\cdot 25+.375 * B B * B B *(1 \bullet+E)) /\) 1(PI*RM**3)
\(C(J)=C(J) / A B S\)
\(C N=C(1)\)
DO \(312 \mathrm{~J}=1, \mathrm{~N}\)
\(C(J)=C(J) / C N\)
RETURN
END
SUBROUTINE RACH(Y,Y1,APK,P1)
RACHINGER SEPARATION - FOR CO RADIATION
COMMON \(\mathrm{X}(300), \mathrm{B}(300), C(300), N\)
DIMENSION Y(300), Y1(300)
\(\mathrm{P} 1=1.78892 / 1.79021 * S I N(A P K * 3.1415927 / 360\).
\(\mathrm{P} 2=1.79278 / 1.79021 * \operatorname{SIN}(A P K * 3.1415927 / 360\).
\(P_{1}=\operatorname{ATAN(P1/SQRT(1.-P1**2))}\)
\(P 2=\operatorname{ATAN}(P 2 / S Q R T(1 .-P 2 * * 2))\)
\(D=(P 2-P 1) * 360.13 \cdot 1415927\)
DO \(606 \mathrm{I}=2, \mathrm{~N}\)
```

```
        IF (X(I)-X(I)-D) 606,605,605
        DO \(600 \mathrm{~J}=1\), I
        \(Y 1(J)=B(J)\)
        \(I A=I+1\)
        GO TO 607
        606
607
601
602
604
608
603
```

$C$
$C$
C
SUBROUTINE CENTROID(Y,CG,SDCG,A,Z,D)
COMMON X(300), B(300), C(300), N
DIMENSION Y(300)
DOUBLE PRECISION S(6),P1
DO $500 \mathrm{~J}=1,6$
$S(J)=0$.
P1 $=0$.
DO $501 \mathrm{~J}=1, \mathrm{~N}$
$S(1)=S(1)+Y(J) * X(J) * C(J)$
$S(2)=S(2)+(X(J)-X(1)) / D * C(J) * X(J)$
$S(3)=S(3)+(X(J)-X(N)) / D * C(J) * X(J)$
$S(4)=S(4)+Y(J) * C(J)$
$S(5)=S(5)+(X(J)-X(1)) / D * C(J)$
$S(6)=S(6)+(X(J)-X(N)) / D * C(J)$
$C G=(S(1)-Z * S(2)+A * S(3)) /(S(4)-Z * S(5)+A * S(6))$
DO $502 \mathrm{~J}=1, \mathrm{~N}$
$P 1=P 1+Y(J) *(X(J)-C G) *(X(J)-C G) * C(J) * C(J) /(S(4)-Z * S(5)+A * S(6)) /$
1(S(4)-Z*S(5)+A*S(6))
SDCG $=P 1+A / 5 * *(S(3)-C G * S(6)) *(S(3)-C G * S(6)) /(S(4)-Z * S(5)+A * S(6))$
1/(S(4)-Z*S(5)+A*S(6))
SDCG=SDCG+Z/5•*(CG*S(5)-S(2))*(CG*S(5)-S(2))/(S(4)-Z*S(5)+A*S(6))/ / / (5)
$1(S(4)-2 * S(5)+A * S(6))$
SDCG = DSQRT(SDCG)
RETURN
END
SUBROUTINE PEAK (Z,PK,SDPK,R1,R2)

PARABOLIC FIT TO PEAK MAX
COMMON X(300), BG(300), COR(300), N
DIMENSION $A(99,3), B(99), C(3,3), D(3), R(200), Y(3), Z(300), E(3,3)$,
1 AA(3)

```
    TYPE DOUBLE A,B,C,D,R,Y,AA,PK,SDPK
    DO 406 J=1,N
    IF (RI-X(J)) 403,403,406
    IRI = J
    GO TO 411
4 0 6 ~ C O N T I N U E ~
4 1 1 ~ D O ~ 4 0 4 ~ J = I R I , N
    IF (R2-X(J)) 407,405,404
405 IR2 = J
    GO TO 412
407 IR2 = J-1
    GO TO 412
404 CONTINUE
412 NP = IR2-IR1+1
    DO 401 I=IRI,IR2
    A(I-IRI+1,I)=1.0
    A(I-IRI+1,2)=X(I)
    A(I-IRI+1,3)= X(I)*X(I)
    B(I-IRI+I)=(Z(I)-BG(I))*COR(I)
    D(1)=0.0
    Y(1)=0.0
    SS=0.
    DO 1 I=1,NP
    SS=B(I)*B(I)+SS
    D(1)=A(I,I)*A(I,I)+D(I)
    1 Y(1)=A(I,I)*B(I)+Y(I)
    Y(1)=Y(1)/D(1)
    IR=0
    DO 5 K=2,3
    IR=IR+1
    R(IR)=0.0
    DO 2 I=1,NP
    2R(IR)=A(I,K-1)*A(I,J)+R(IR)
    R(IR)=R(IR)/D(K-I)
    DO 3 I=1,NP
    3 A(I,J)=A(I,J)-A(I,K-1)*R(IR)
    D(K)=0.0
    Y(K)=0.0
    DO 4 I=1,NP
    B(I)=B(I)-A(I,K-1)*Y(K-1)
    Y(K)=A(I,K)*B(I)+Y(K)
    4D(K)=A(I,K)*A(I,K)+D(K)
    5 Y(K)=Y(K)/D(K)
        DO 55 I= 1,3
        AA(I)=Y(I)*D(I)**0.5
        SS=SS-AA(I)*AA(I)
        FN=NP-I
        IF(FN) 51,51,52
    SD = O.
        GO TO 53
52 SD=SQRT(SS/FN)
5 3
    DO 55 J=1,I
55 E(I,J)=SD
        IRS=-3
        DO 8 K=1,3
        IRS=IRS-K+4
        IR=IRS
        DO 8 JJ=1,K
        J=K-JJ+1
        C(K,J)=Y(J)
        IF(JJ-1) 8,8,6
        6O 7 I=2,JJ
        C(K,J)=C(K,J)-C(K,K-I+2)*R(IR)
```

        E(K,J)=E(K,J)+ABSF(R(IR)*E(K,K-I+2))
    7 IR=IR-1
    8 IR=IR-3 +K
        PK = -C( 3,2)/(2.*C(3,3))
        SDPK = DSQRT((PK*E(3,2)/C(3,2))**2 + (PK*E(3,3)/C(3,3))**2 -
    1 2*(PK**2)*E(3,2)*E(3,3)/C(3,2)/C(3,3))
        RETURN
        END
    ```
    SUBROUTINE SMOOTH(YC,NUMP, NS, NP, NCEN)
    DIMENSION \(S(11,500), D 1(500), D 2(500), N S(10), Y C(300)\)
    COMMON X(300), B(300), C(300), IIJM
    \(X 1=X(N C E N)\)
    DO \(22 \mathrm{~N}=1\), NUM
    \(S(1, N+3)=Y C(N)\)
    NUM \(=\) NUM +3
    DO \(20 \quad \mathrm{~N}=1.3\)
    \(D 2(N)=0\).
\(20 \mathrm{D} 2(\mathrm{NUM}+\mathrm{N})=0\).
    DO \(5 \mathrm{~N}=4\), NUM
    Dl \((N)=S(1, N)\)
    \(K=0\)
    DO \(12 I=1, N P\)
    J=NS (I)
    DO \(9 \mathrm{~N}=4\), NUM
    \(D 2(N)=D 1(N)\)
    \(K=K+1\)
    DO \(7 \mathrm{~L}=4\), NUM
    D1 \((L)=(D 2(L-3)-18 * * D 2(L-2)+63 * * D 2(L-1)+164 * * D 2(L)+63 * * D 2(L+1)\)
    1-18.*D2(L+2)+D2(L+3))/256.
    IF \((K-J) 8,10,8\)
    \(M=I+1\)
    DO \(11 \mathrm{~N}=4\), NUM
    \(S(M, N)=D 1(N)\)
    CONTINUE
    IF (NUMP) 13,15,13
    NCEN \(=\) NCEN +3
    DO \(18 \mathrm{~N}=1\), NP
    IF (NUMP-NS(N)) \(18,17,18\)
    NUMP \(=N+1\)
    GO TO 19
18 CONTINUE
19 PUNCH 105,(S(NUMP,N),N=NCEN,NUM)
    \(M=\) NCEN -3
    DO \(14 \mathrm{~N}=1, \mathrm{M}\)
    \(K=N C E N-N+1\)
    \(D 2(N)=S(N U M P, K)\)
    PUNCH 105, (D2(N), N=1,M)
    NUMP \(=\) NUMP -1
    PRINT 109, XI,NS(NUMP)
    PRINT 107, ( \(N S(N), N=1, N P\) )
    \(M=N P+1\)
    DO \(16 \mathrm{~N}=4\), NUM
    PRINT \(108, X(N-3),(S(J, N), J=1, M)\)
    FORMAT (10(F6.0,\$.\$))
    FORMAT (////1X,\$ANGLE\$,5X,\$INPUT\$,3X,10(\$SMOOTH \$,I2,1X))
    FORMAT ( \(1 \mathrm{X}, \mathrm{F} 7.2,11 \mathrm{~F} 10.0\) )
    FORMAT (//1X,\$PUNCH PEAK FOR ALPHA \(=\$, F 7.2 / 1 X, \$ S M O O T H I N G \mathbb{S}, I 3)\)
    END
10. Appendix: Corrections Necessary When Using Other Radiations, Samples and Geometries.
Specific program steps are referred to by the letter identification which appears in the identification columns.
A. Other radiations

In "ANG":
The value of the diffraction angle in the monochromator must be inserted into the calculation of DA (step A). In this case, for Co \(\mathrm{K} \alpha\) radiation and a quartz crystal, \(2 \alpha=31.2333^{\circ}\). It is also necessary to use the \(K \propto\) wavelength to determine \(S L\) (step \(B\) ); here \(\lambda=1.79021 \AA\) 。

In "RACH":
The values of the \(K \alpha_{;} K \alpha_{1}\), and \(K \alpha_{2}\) wavelengths (i. e., \(\lambda, \lambda_{1}\), and \(\lambda_{2}\) ) are necessary to evaluate \(P 1\) and \(P 2\) (steps \(A\) and B) as follows:
\[
\begin{aligned}
& \mathrm{Pl}=\frac{\lambda_{1}}{\lambda} \sin \left[\operatorname{APK}\left(\frac{\pi}{360}\right)\right] \\
& \mathrm{P} 2=\frac{\lambda_{2}}{\lambda} \sin \left[\operatorname{APK}\left(\frac{\pi}{360}\right)\right]
\end{aligned}
\]
B. Other samples

In " \(A N G\) ":
The parameters of the power series fit to the atomic scattering factors must be used in steps C, D, and E. RM in step H is the factor \(\mu \mathrm{r}\) in section 3 and must be evaluated as noted there.
C. Other geometries

In "ANG":
When no monochromator is used, DA (step A) should be set equal to one. For a flat sample, eliminate \(C C, E\), and \(R M\) (steps \(F, G\), and \(H\) respectively) and set \(A B S\) (step \(I\) ) equal to one.

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[^1]:    Using the Debye-Scherrer technique, it is possible to calibrate the diffractometer by examining the reflection on both positive and negative sides of the diffraction cone ${ }^{[6,8]}$. This gives a correction factor to the angle read from the instrument.

