

A UNITED STATES
DEPARTMENT OF
COMMERCE
PUBLICATION



National Bureau of Standards
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OCT 28 1969

NBS TECHNICAL NOTE 499

Multiple Scattering Corrections for the Associated-Particle Neutron Time-of-Flight Technique

**U.S. DEPARTMENT OF COMMERCE
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TECHNICAL NOTE 499

ISSUED OCTOBER 1969

Nat. Bur. Stand. (U.S.), Tech. Note 499, 68 pages (Oct. 1969)
CODEN: NBTNA

Multiple Scattering Corrections for the Associated-Particle Neutron Time-of-Flight Technique

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Abstract

The computer code, MAGGIE, for the calculation of multiple scattering and sample attenuation in neutron differential cross-section measurements, has been revised and corrected. The particular case of the scattering geometry required by the associated-particle time-of-flight is considered in detail.

Key words: associated-particle, cross-section, Monte Carlo, multiple scattering, neutron, time-of-flight

MULTIPLE SCATTERING CORRECTIONS FOR THE ASSOCIATED-PARTICLE
NEUTRON TIME-OF-FLIGHT TECHNIQUE

Allan C. B. Richardson

I. INTRODUCTION

Measurements of fast neutron elastic and inelastic differential cross sections have, for many years now, usually been done using one of two time-of-flight techniques. The first of these requires a pulsed source of neutrons, and energy separation of the various neutron groups is then achieved by time correlation of the scattered neutrons with the incident neutron pulse. This technique has the advantage of flexibility of incident neutron energy and intensity, and the disadvantages of a relatively high time-correlated background and a low duty cycle. The other, and less commonly used, technique utilizes the detection of a charged particle from the neutron source reaction to tag the incident neutrons in time and direction. Energy separation of the various scattered neutron groups is then achieved by time correlation with the incident neutrons. This "associated-particle technique" has the advantages of very small time-correlated background, inherent absolute determination of the incident correlated neutron flux, and high duty cycle; but it suffers from limitations on the available neutron intensity and energy. However, at energies and intensities where this technique is applicable, it is the method of choice, since it is capable of yielding results of high accuracy without the ambiguities introduced by the time-correlated backgrounds and the massive shielding required by pulsed source techniques. The source reactions eligible for this method are those involving very light nuclei, and therefore capable of producing a light (and thus energetic) stable recoil nucleus. The $T(d,n)^4\text{He}$ reaction, producing 14-15 MeV neutrons at 90° over a rather wide range of incident particle energies, is most commonly used; other reactions that have been employed are $D(d,n)^3\text{He}$ and $T(p,n)^3\text{He}$, both of which produce lower energy neutrons. We will confine the discussion here to the $T(d,n)^4\text{He}$ case, although the method described is more generally applicable.

The associated-particle technique has scattering sample requirements that are quite different from those for a pulsed source. Instead of a relatively uniform incident neutron flux across the sample, the correlated neutron beam is highly directional. The angular distribution about the neutron beam line is usually well approximated by

$$I = I_0 e^{-\left(\frac{\theta}{\theta_0}\right)^2},$$

with θ_0 typically only a few degrees [1]. This property can be very useful [2]. It provides a high degree of neutron collimation without the need for massive collimators, which, in the case of 14 MeV neutrons, can produce substantial degradation of the initially monoenergetic neutron

beam. However, it forces the use of scattering samples of uniform thickness so that the cylindrical or spherical samples commonly used to simplify multiple scattering corrections are immediately ruled out. Otherwise a detailed knowledge of the neutron beam shape and extremely accurate alignment of this beam with respect to the scattering sample is needed. This alignment problem is further complicated because the center of the neutron beam slowly moves back as the neutron producing target ages with use. In measuring angular distributions of scattered neutrons it is of course also desirable to reduce the amount of scattering sample not directly in the neutron beam, so as to minimize multiple scattering.

A scattering sample in the form of a truncated cone, axis lying along the neutron beam, best satisfies all of these requirements. The origin of this cone is taken sufficiently far behind the source to allow for finite spot size on the neutron producing target, to provide some flexibility in alignment, and also to make allowance for changes due to target aging during a run. In order to make best use of the available neutron intensity, one must also use the thickest sample possible. The limit is set by either the angular resolution required at 90° or the time resolution required. A typical geometry is shown in figure 1.

In either case the resulting samples are sufficiently thick to require a careful multiple scattering correction. None of the analytical techniques [3], useful at energies up to a few MeV, are adequate at 14 MeV, the energy most commonly used for measurements of this type, since at energies above 6-7 MeV the diffraction peaks in the elastic angular distributions become too numerous. The only method of sufficient generality is the Monte Carlo technique. A survey of existing Monte Carlo codes revealed none for this particular geometry, but it was immediately apparent that the code "MAGGIE," developed by Parker, Towle, Sams, et al. at Aldermaston [4,5] contained all of the other elements important to such a calculation. For example, this code easily accommodates the neutron source distribution specified above. In addition, MAGGIE calculates an energy spectrum at each angle, so that false peaks due to double scattering can easily be identified. This sophistication is often useful, for example: the measurement of inelastic scattering from the very weakly excited 7.66 MeV 0^+ level in ^{12}C is easily confused by double scattering from the more easily excited 4.43 MeV 2^+ level in this nucleus.

A corollary need of any Monte Carlo neutronic calculation is an easily accessible, but at the same time sufficiently general, file program for the nuclear data required. This was available in a companion code to MAGGIE, entitled MOULD [6]. We have therefore modified the code MAGGIE so that it now handles the geometry required for associated-particle time-of-flight measurements. During the course of this modification we also corrected a few coding errors found in the original version, and made several additional modifications, principally updating the code to current computer syntax and capabilities.

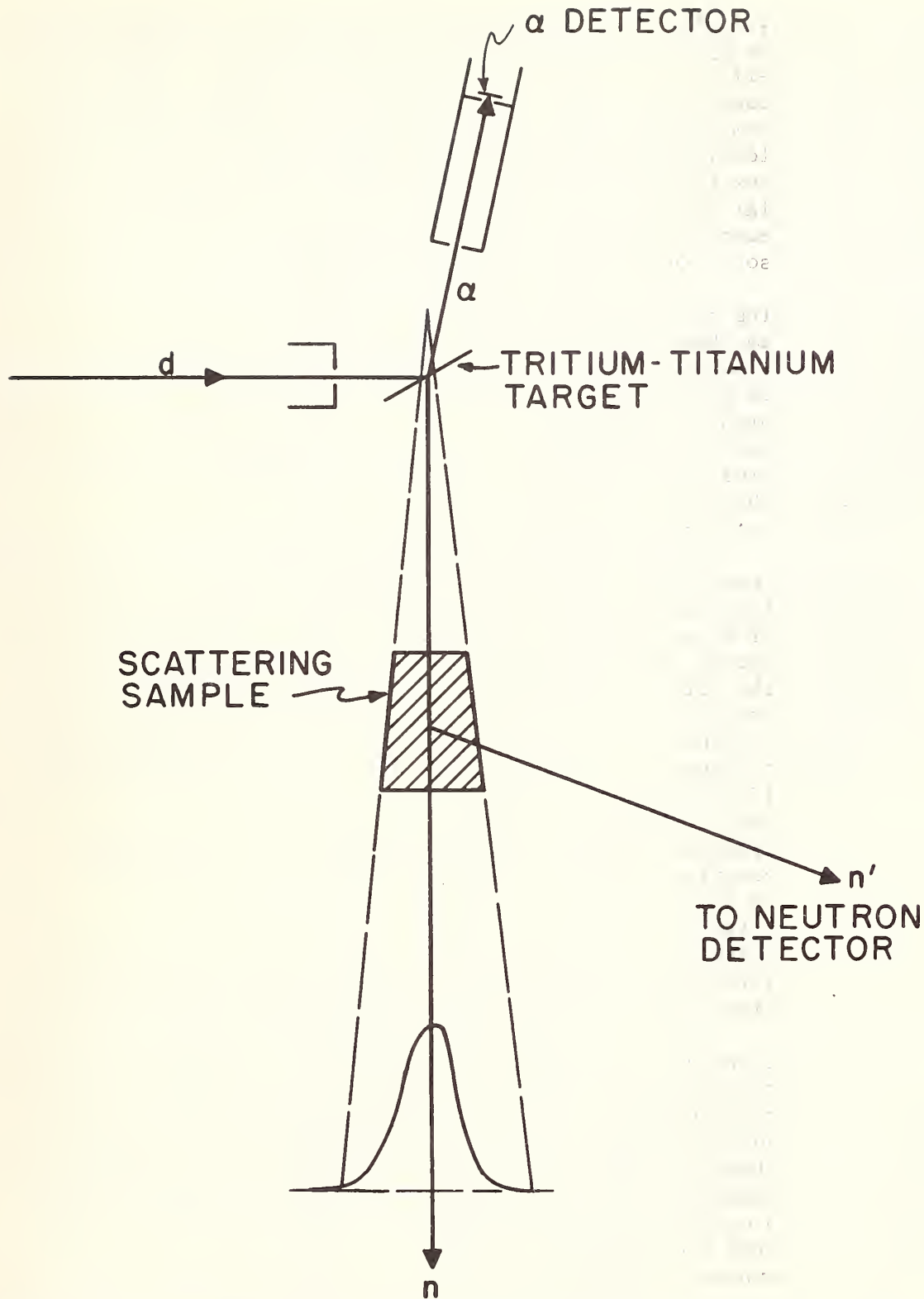


Fig. 1. The experimental geometry for associated-particle measurements at 14 MeV. The associated neutron intensity profile is indicated at the bottom of the figure.

Although a complete description of the code would be out of place here (see references 4 and 5 for details), a general overview of the program will facilitate later discussion of specific details. The point of view adopted is first to sample the scattering sample geometry and the available nuclear interactions using straightforward Monte Carlo techniques. Then, at each collision point, the weighted probability of scattering and escape to each of 33 detector positions (angles) in a half-plane lying to one side of the neutron beam axis is computed using the experimental data for elastic events and each of those inelastic processes of interest. The use of weighted probabilities of scattering to each of the detector positions at the final collision in the sample in place of a completely Monte Carlo approach results in a greatly reduced computation time. This is because the small solid angle subtended by each of the detectors (typically 10^{-3}) makes a final Monte Carlo scattering particularly inefficient. The output angular distributions obtained from applying this procedure to, typically, 1000 interacting neutrons are then reflected about the experimental input data and used as input for a second iteration. Two or three iterations are usually sufficient to obtain a convergent result.

In section II we give a description of the changes made in the code. First, the new coding for the scattering geometry used in the associated particle technique is described. Next we discuss a number of coding errors in the original version. Finally some changes are described that simplify the code and adopt it to FORTRAN V syntax. In section III we give some results obtained using data from the scattering of 14 MeV neutrons on natural carbon. A listing of those subroutines entirely rewritten or having extensive changes is given in Appendix A. In figure 2 the calling sequence for all of the components of the program is shown. A brief description of all of these subroutines appears in Appendix B.

II. DESCRIPTION OF MODIFICATIONS TO THE MONTE CARLO CODE MAGGIE

A. Changes Due to the New Scattering Sample Shape

Four subroutines are affected by changing the shape of the scattering sample. These are: 1) subroutine INPUT - those sections where scatterer parameters are read in and the flux attenuation factor is calculated are affected. The flux attenuation factor is defined as the ratio of incident flux along the axis of the sample to average flux in the entire sample. 2) CRNEU, the subroutine that creates incident neutrons at the entrance face of the scatterer by random sampling of the incident neutron spatial distribution. 3) TRACK, the tracking subroutine. 4) FPATH, the subroutine for calculating the probability of neutron escape from the sample in the direction of each of the assumed detector positions, for each collision point arrived at in TRACK.

It is worth noting here, although not necessary to what follows, that the data used for a) the neutron track lengths in the Monte Carlo sampling of the sample shape from subroutine EGMV, and b) the Monte Carlo sampling of reaction type and angular distribution at each collision by subroutine

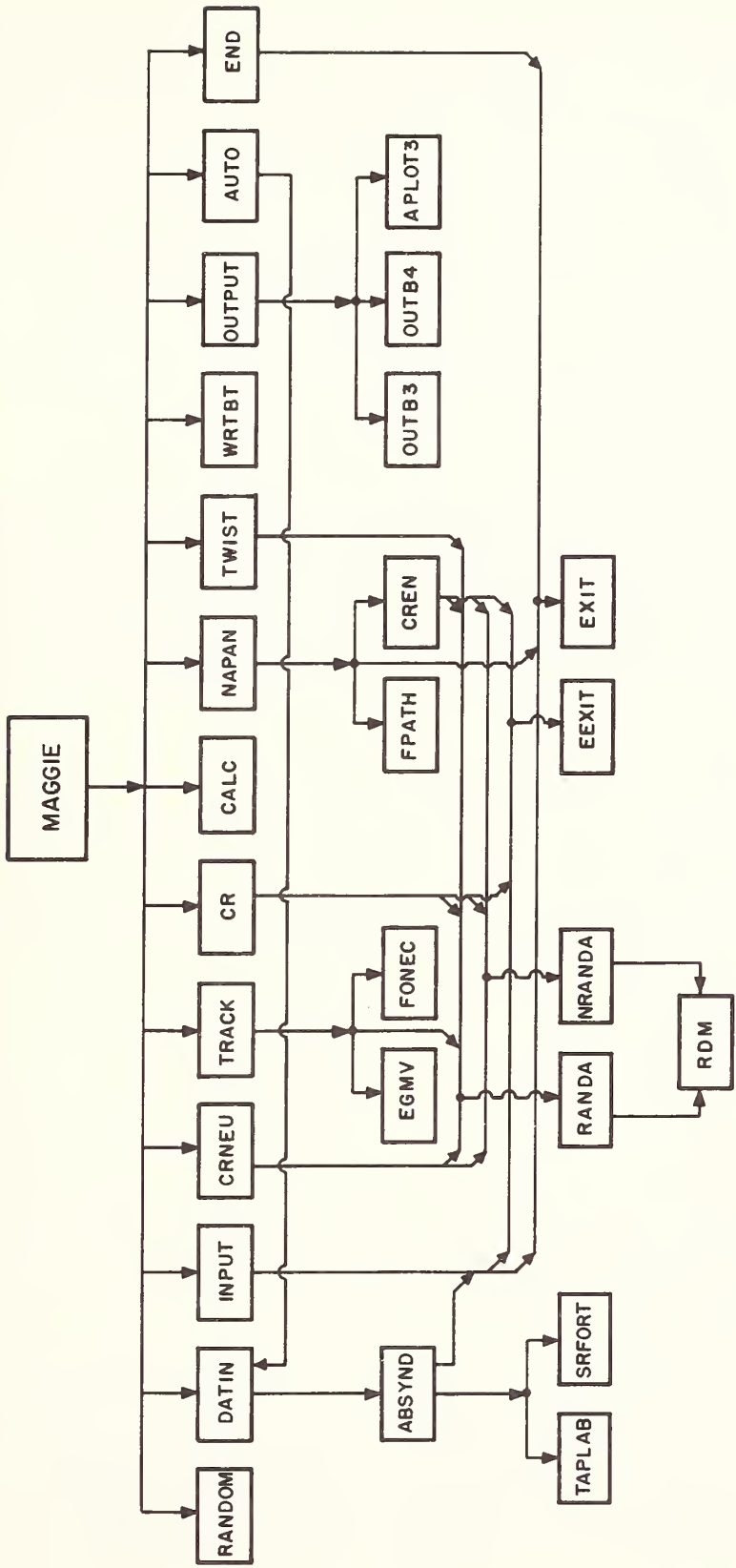


Fig. 2. The calling hierarchy for the present version of code MAGGIE.

CR, as well as c) the transmission probabilities in the direction of the detectors used in subroutines FPATH and NAPAN, are all drawn from the MOULD nuclear data file tape and not from input experimental data. The only experimental data employed are angular distributions used for the calculation of the scores for the relative probability of scattering to the various assumed detector positions in subroutine NAPAN. The elastic angular distribution (for the incident neutron energy only) used here is also automatically updated after each iteration by the subroutine AUTO. Thus it may be necessary to update the angular distributions on the MOULD file tape before and during the course of data-processing.

We now describe the changes due to the new scattering geometry in some detail for each of the subroutines affected.

1. INPUT

a. Cards 0207-0226 are changed to eliminate parameters required for samples consisting of concentric cylindrical shells, with common axes perpendicular to the beam axis, and to substitute those needed to characterize a truncated cone sample, axis lying along the neutron beam, and origin behind the neutron source. The new variables are HITE, FRAD, and ANGLE; the height, entrance face radius, and half-angle of the sample, respectively. HITE is immediately redefined as HITE/2, a more convenient quantity in subsequent calculations.

b. Card 0245, the calculation of the maximum angle subtended by the scattering sample at the neutron source, is changed to conform to the new geometry.

c. Cards 0301-0303 have been replaced. The correct expression for the flux attenuation factor for the new geometry is

$$\frac{F_o}{F} = \frac{K(o)}{\bar{K}} \cdot \frac{C}{N\lambda} \cdot \frac{h(r_1^2 + r_1\Delta r + \Delta r^2/3)}{(1 - \cos\theta_m) d^2},$$

where h is the half-length of the scattering sample, r_1 the radius of its entrance face, Δr the difference in radii of the entrance and exit faces, and d the distance from the source to the entrance face of the sample. The remaining symbols are defined as in the original.

d. The original version of MAGGIE utilized a sample which was not symmetric about the axis of the incident neutron flux. Thus the experimental sampling of scattered neutrons in the detector plane was not truly representative of the scattering into 4π , which is employed in the program to calculate the flux attenuation factors and which are used in turn to infer the integrated cross-sections. This necessitated a small correction which was calculated with the help of a classification of the outgoing

Monte Carlo tracks vs. energy and the angle with respect to the scattering sample axis. Since the present scattering sample is symmetric about the incident neutron flux the experimental sampling of scattered neutrons in the detector plane is representative, and no correction is required. Accordingly, cards 0331-0338 in INPUT, cards 1544-1548 in DATIN, and cards 1735-1743 in MAGGIE are deleted. Card 1734 of MAGGIE is replaced by the statement

20 CONTINUE

The output of this table, by cards 3695-3715 of subroutine OUTB3, is also deleted.

2. CRNEU

Cards 3117-3119, 3140-3141, 3143-3147 and 3151 are replaced as shown in the listing. The new coding creates neutrons randomly scattered over the face of the conical scattering sample, in accordance with the specified spatial distribution. The old reference to multiple materials is not deleted, but the sample is now designated "material one." Similarly, the register containing neutrons which miss the sample is retained, although for this geometry misses occur with very low frequency.

3. TRACK

The entire subroutine has been replaced. The geometry is illustrated in figure 3. The subroutine is entered with the starting point (x_0, y_0, z_0) and the direction cosines $(\cos\theta_x, \cos\theta_y, \cos\theta_z)$ already defined. A random sampling of the neutron mean free path establishes the track length to the next possible collision. The track vector is then extended until it intersects the plane determined by the endface of the sample in the direction of travel of the neutron, and the path length from the starting point to this intersection is computed. A comparison of this intersection point with the radius of the endface establishes whether the track is in the direction of the endface or the curved surface of the sample, and the program branches accordingly. If the track passes through the endface, the path length calculated above is compared with the Monte Carlo track length to the next collision and the coordinates of collision or escape from the surface of the sample as well as the time elapsed along the track are computed in a straight-forward way.

We consider now the procedure used for a neutron headed toward a curved surface of the scattering sample. The coordinates of the endpoint of the previously determined Monte Carlo track length are first found. These are used to calculate the projected distance, perpendicular to the symmetry axis, of this endpoint from the symmetry axis, as well as the length of the radius of the sample lying along this projection. Comparison of these quantities determines whether a collision occurs within the sample or not. If so, the coordinates are already known and the time elapsed along the track is all that remains to be calculated. If not, it is necessary to calculate the coordinates of the neutron's exit from

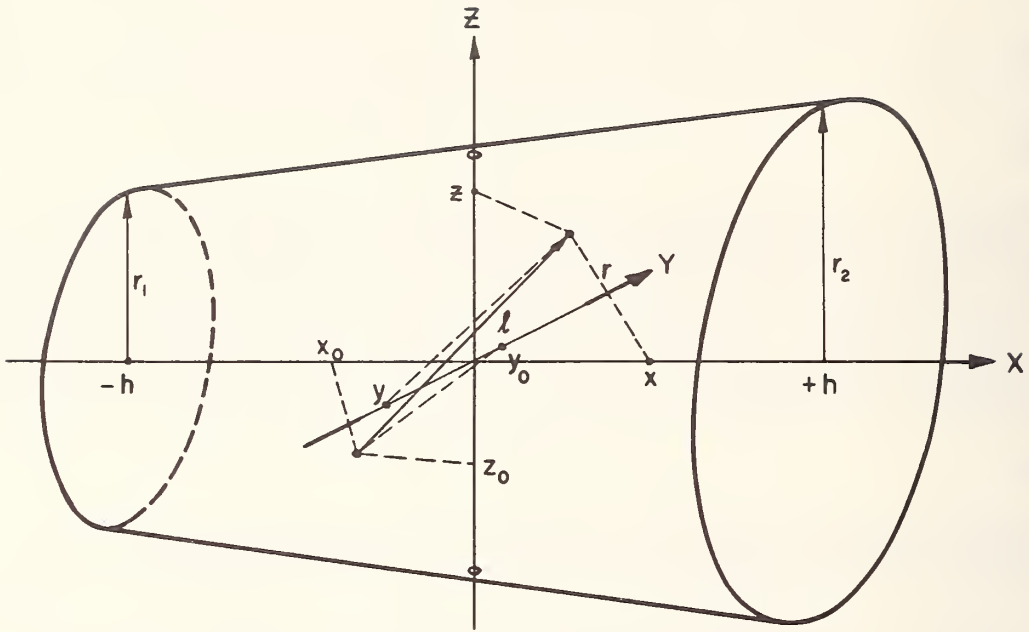


Fig. 3. The geometry for subroutine TRACK. The trajectory shown is that of a neutron which, after collision at point x_0, y_0, z_0 escapes from the curved surface of the sample at point x, y, z .

the surface of the cone. Designate l the path length to the surface and x, y, z the coordinates of the intersection of the track with the surface. We have, referring to figure 3,

$$\begin{aligned}x - x_0 &= l \cos \theta_x \\y - y_0 &= l \cos \theta_y \\z - z_0 &= l \cos \theta_z \quad ,\end{aligned}$$

so that the projected distance, r , from the symmetry axis of the cone to the point (x, y, z) is given by

$$\begin{aligned}r^2 &= y^2 + z^2 \\&= l^2 \sin^2 \theta_x + 2l(y_0 \cos \theta_y + z_0 \cos \theta_z) + y_0^2 + z_0^2 \quad ,\end{aligned} \quad (1)$$

where we have used the identity

$$\cos^2 \theta_x + \cos^2 \theta_y + \cos^2 \theta_z = 1 .$$

Since we know the parameters of the cone we can also calculate r from the x coordinate (again see figure 3.):

$$r = r_1 + (h+x) \tan \theta_0 \quad , \quad (2)$$

where θ_0 is the half-angle of the cone. Eliminating r from eq'ns. (1) and (2) and arranging the terms as a quadratic in l , we obtain

$$\begin{aligned}&l^2 \left[1 - \cos^2 \theta_x (1 + \tan^2 \theta_0) \right] \\&+ 2l \left[y_0 \cos \theta_y + z_0 \cos \theta_z - \cos \theta_x (r_1 \tan \theta_0 + h + x_0) \tan^2 \theta_0 \right] \\&+ \left[y_0^2 + z_0^2 - (r_1 + (h+x_0) \tan \theta_0)^2 \right] = 0.\end{aligned} \quad (3)$$

The positive root of this equation is the required path length to the surface, and the subroutine FONEC is then called to provide the x, y , and

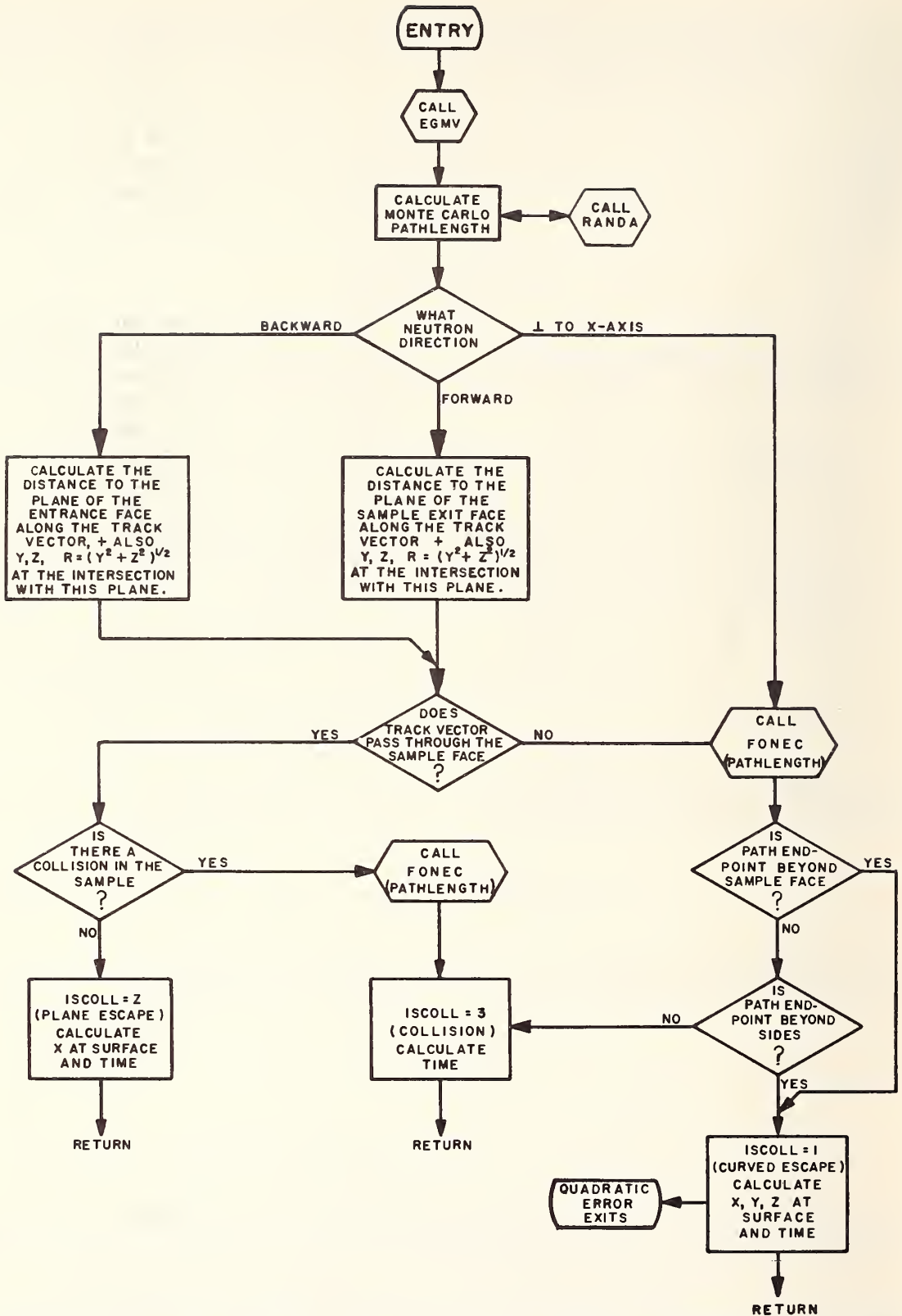


Fig. 4. Subroutine TRACK. The subroutines EGMV and FONEC return the neutron mean free path and the coordinates at the end of a track, respectively. RANDA returns a random number.

z coordinates. Error prints are provided for the cases of two positive, two negative, imaginary, or indeterminate roots as solutions to eq. (3). However none of these situations has occurred in many tens of thousands of tracks, although at least some of them are mathematically possible. Once ℓ has been calculated, determination of the coordinates of escape and the track time are straightforward. The subroutine also returns an index that indicates the fate of the neutron, i.e., curved escape, plane escape, or collision. A block diagram is shown on figure 4.

4. FPATH

This subroutine has also been entirely replaced. The same simplifying assumption is made as in the original coding--that is, the dimensions of the scattering sample are considered to be negligible compared to the flight path to the detector. This has the effect that the angle and flight path to any particular detector position may be considered to be the same from any scattering point in the sample. The quantities to be calculated are the neutron path length, p , in the sample from any point (x, y, z) to the surface of the sample and the scattering angle, θ_d , in the direction of a particular detector, at the angle ψ . The geometry is shown in figures 5 and 6. All paths are assumed to be parallel to the x-y plane, and all detector positions satisfy the condition $y \geq 0$.

We note that the locus of the intersection of a cone with a plane parallel to its axis is an hyperbola, and write, using the notation shown in figure 5,

$$\frac{X^2}{a^2} - \frac{Y^2}{b^2} = 1 \quad ,$$

$$\text{or} \quad Y^2 = X^2 \tan^2 \theta - z \quad , \quad (4)$$

$$\text{since} \quad b = z \quad \text{and} \quad a = z/\tan \theta.$$

Now, referring to figure 6, we may also write

$$Y = p \sin \psi + y$$

$$\text{and} \quad X = p \cos \psi + x' \quad , \quad (5)$$

$$\text{where} \quad x' = R + h + x = \frac{r_1}{\tan \theta} + h + x.$$

Again we reduce the problem to a quadratic in path length, p , after

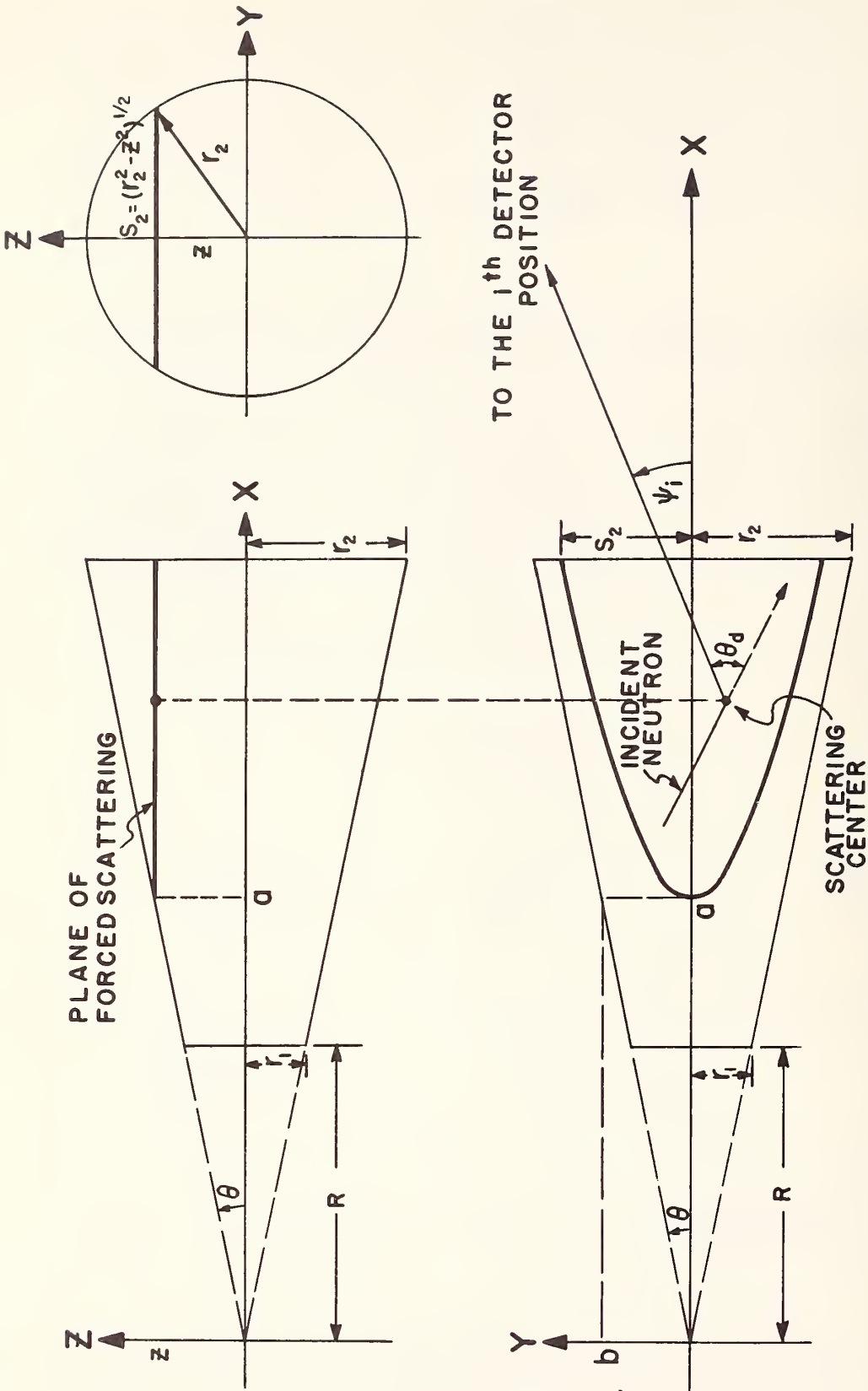


Fig. 5. The geometry for subroutine FPATH. The plane of forced scattering is located at a height $r_1 < z < r_2$ in this instance. The track of the incident neutron does not, in general, lie in this plane as shown.

combining eq's (4) and (5) to eliminate the variables X and Y to obtain:

$$\begin{aligned}
 & p^2(\sin^2\psi - \cos^2\psi \tan^2\theta) \\
 & + 2p(y\sin\psi - x'\cos\psi \tan^2\theta) \\
 & + (y^2 - x'^2 \tan^2\theta + z^2) = 0 .
 \end{aligned} \tag{6}$$

For each value of z there will be a pair of angles, α and β (see figure 6, which define the limits of the hyperbolic curve. These angles are given by

$$\begin{aligned}
 \tan \alpha &= \frac{(r_1^2 - z^2)^{1/2} - y}{h + x} \\
 \tan \beta &= \frac{(r_2^2 - z^2)^{1/2} - y}{h - x} .
 \end{aligned} \tag{7}$$

The subroutine is diagrammed in figure 7. After testing for the special case of exit perpendicular to the beam axis, the program tests, using eq's (7), for exit through the endfaces vs. the curved sides for the exit angle, ψ , and branches accordingly. The path length in the former case is a straightforward calculation, and for curved escape is the positive root of eq. (6). It should be noted that in the limiting case of cylinder ($\tan\theta=0$) this subroutine, unlike all of the others, does not work. A modified subroutine for a cylindrical sample is given in the Appendix following the listing of FPATH for a truncated cone.

B. Correction of Coding Errors [7]

1. INPUT

- a. Statements resulting in LGR=0 in the table look-up for the mean free path (following card 0308) have all been changed to give LGR=1. LGR=0 references not a mean free path, but instead the last tabular value of collision probability in the previous MOULD table. The revised look-up supplies the mean-free path at the lowest tabulated energy for any neutron at or below that energy.
- b. On the card following statement 5001, LRG has been replaced by the correct variable, LGR.
- c. Card 0441 is replaced by

$$IMINM = IMAXM+1 .$$

This change yields sequential storage of supplementary ranges of experimental angular distributions, and thus avoids loss of needed storage space in the array SUPVAL (I,J).

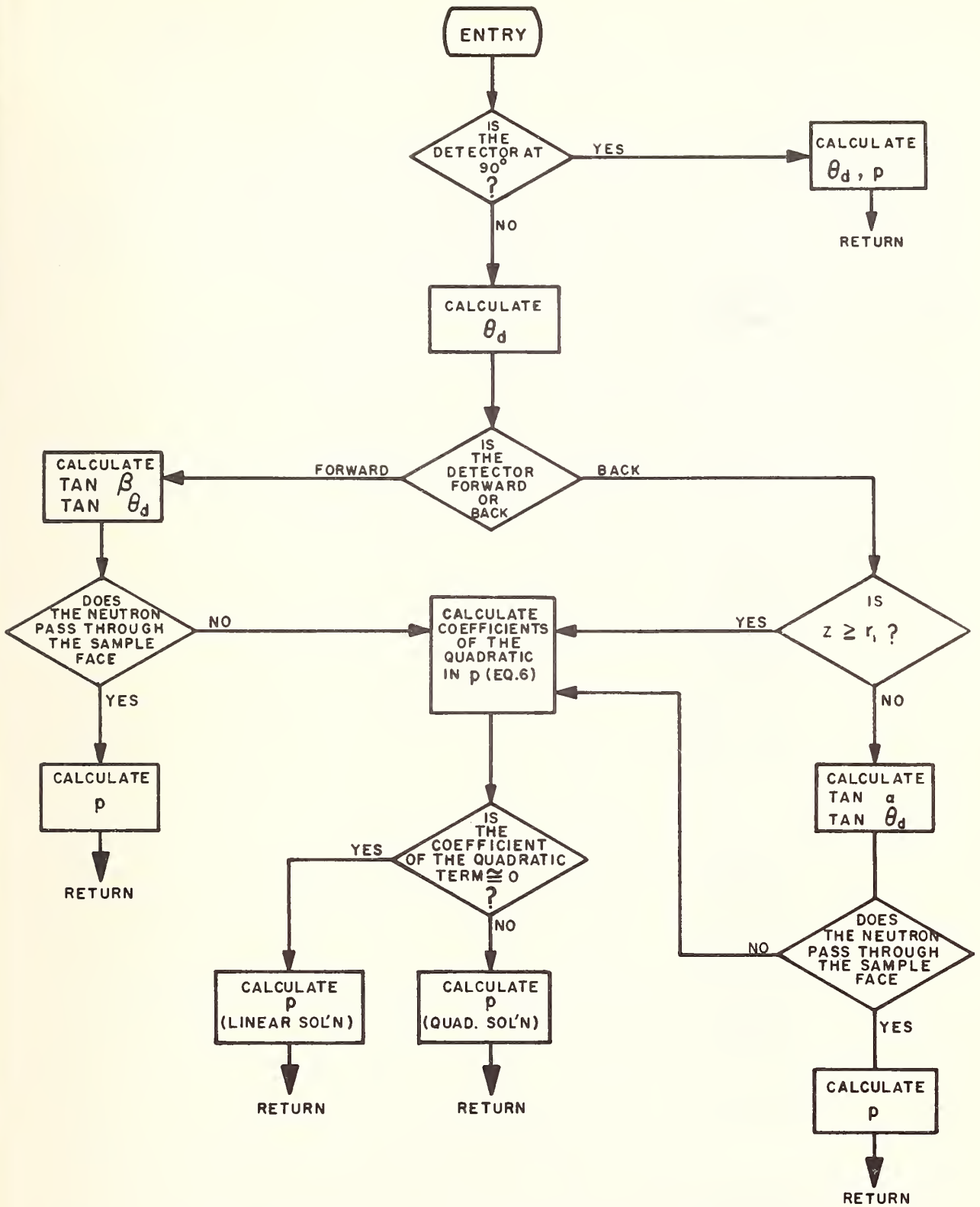


Fig. 7. Subroutine FPATH, for a truncated cone. The subroutine for a cylinder is similar.

2. EGMV

LGROUPE=0 has been eliminated for the reason given in 1a. above.

3. NAPAN

a. The angle used for the interpolation of the experimental distributions, as well as for calculation of the neutron energy after collision is the laboratory angle, calculated in the subroutine FPATH. For this reason the experimental angular distributions must be read in at laboratory angles. Also the center-of-mass energy calculation on cards 3494 and 3497-3501 (in the case of center-of-mass MOULD angular distribution tables) is in error. This has been changed to a calculation in laboratory coordinates by replacing card 3494 by

```
GO TO (36,11,11,11,11,11,11,11,36,11,11,11,11,11,11) NFORM
```

and deleting cards 3497-3501.

b. The sequence of cards 3502-3512, which picks the appropriate experimental angular distribution values, does not refer to the correct angular distributions, nor to the correct neutron energy--that before collision. It has been replaced by the sequence:

```
11      IF(IMAX(J).LE.0) GO TO 15
        IMINM=IMIN(J)
        IMAXM=IMINM+IMAX(J)-1
        DO 50 JJ=IMINM,IMAXM
        IF(EIN-ENVAL(J,JJ))50,51,52

50      CONTINUE

51      KK=JJ
        GO TO 19

52      IF(JJ.LE.IMINM) GO TO 15
        KK=JJ-1
```

c. On the card following statement 5001, LRG has been replaced by the correct variable, LGR.

d. LGR=0 has been eliminated for the reason given in 1a. above.

e. The option that provides for suppression of printout of various scores, through setting the indicators IANALA, IANALB, IANALC, and IANALD negative, should not suppress storage of these scores by NAPAN. The time saved is negligible and the elastic scores, DEB(I,L), are required by the automatic iteration scheme, AUTO. Cards 3528-43 have been revised to eliminate the dependence on these printout indicators.

4. ABSYND

LGROUP=0 has been eliminated for the reason given in 1a. above.

5. CREN

For the same reason set out in section 3a. above, the center-of-mass energy calculation (cards 2960-66) has been replaced by a calculation in laboratory coordinates by deleting cards 2961-66 and replacing card 2960 with

880 GO TO 800 .

6. OUTPUT, OUTB3, OUTB4

The sum over energy, at each detector angle, stored in the vector DES(J,33) is required by subroutine AUTO. So as to allow calculation of this quantity in subroutine OUTB4, even when no B4 printout is requested, a print indicator, IFPR, has been added to the call for subroutine OUTB4. This is accomplished by the changes shown in the listing for cards 3721, 3808A, 3873, and 3975-4002.

7. AUTO

a. The quantity "chi-squared," on card 2084, should be

$$\chi^2 = \sum_1^{33} \frac{(\sigma(\theta)_{\text{exp}} - \sigma(\theta)_{\text{calc}})^2}{\sigma(\theta)_{\text{exp}}},$$

and not as formerly written, with the denominator squared.

b. It should be noted that the normalization of the Monte Carlo output on cards 2065-2068 and 2088 is not, in general, correct; but only holds for detectors positioned at equal increments of $\cos\theta$ between -1 and +1. For this reason, if subroutine AUTO is used, the input detector positions must satisfy this criterion.

8. FPATH

The original version of this subroutine contained two errors. These do not apply to the new version of FPATH reported here, but are listed for completeness. Card 3356 of the original should read

SINOM = SQRT(1-COSOM**2) ,

card 3358 be deleted, and a new card inserted following card 3359:

FP1 = WORK16**2+WORK17**2-FP2**2 .

C. Other Modifications

1. General

a. Common storage assignments have been handled by combining all common statements into one package, processed by the UNIVAC 1108 "Procedure Definition Processor" using the assembler directive FCOPY (Fortran copy). This package, designated CINCl, is then included at the time of assembly in all subroutines making use of common variables by means of the statement

```
INCLUDE CINCl
```

placed immediately following the subroutine name definition. This process, which effects a considerable economy in the size of the source deck and listing, is characteristic of the UNIVAC 1108 Assembler and FORTRAN V; however, equivalent procedures are often available to other systems. In some subroutines dimension statements still appear for those few variables not in common storage. A block of common storage has been specifically assigned to the subroutine ABSYND variables DATA and IDATA. This carries the dummy label BLANK. The details of these changes in memory storage allotment will be obvious upon examination of the listings.

b. The size of many modern computers obviates the need for linkage. We have deleted references to CHAIN(I,J) and incorporated the balance of subroutine PRELUDE into the main program MAGGIE. Similarly, references to CHAIN(I,J) in subroutines DATIN and AUTO have been replaced by the appropriate CALL and/or RETURN statements.

c. The final F in the names of all library functions (such as SIN, COS, EXP, MIN) has been deleted, as it is not compatible with FORTRAN V. In addition, such functions as FLOAT and INT have been eliminated by using the mixed expressions allowed by FORTRAN V.

d. Disc storage read and write statements have been modified in the main program MAGGIE and subroutine WRTBT to conform to 1108 FORTRAN V usage.

e. A number of formats have been altered. Some of the changes are merely different spacing options, but others are required by the new scattering geometry. A list of card numbers of the affected statements follows:

0226	0226A	0226B
0249	0258	0345
0410	0450	0451
1719	1963	1972
2111	2115	3638
3639	3828	3833

On card 2735 in subroutine CR the variable N has been replaced by the correct variable NSECS.

f. A new random number generator, subroutine RANDOM, entry RDM, has been incorporated, and the required changes in the functions RANDA and NRANDA are shown in the listings.

2. MAGGIE

a. All of the subroutine PRELUDE, except for calls to subroutines CLOCK and CHAIN, has been inserted following card 1654. A call to subroutine RANDOM immediately following the read-in of OCT, the starting value of the random number generator, initializes this generator. OCT is printed out following the call to subroutine INPUT.

b. All references to subroutine CLOCK have been deleted.

c. Card 0530 is eliminated by changing cards 1776, 1778, 1787, 1789, and 1792 to read

```
GO TO 1 .
```

3. ABSYND

a. Card 0525, which assigns the logical tape unit carrying the MOULD nuclear data tape, now reads

```
NUCDAT=9 .
```

On cards 0525, 0530, and 1394 the variable name TAPE has been changed to ATAPE to avoid confusing the compiler. These changes may not be required by other installations.

b. Cards 0553 and 0554 set all storage for the variables DATA and IDATA to zero before each ABSYND run.

c. Card 1398 is modified so as to print only those action numbers processed.

4. CR

The variable names NACT and Q have been changed to NACTV and QV, so as to avoid conflict with the array names NACT(I) and Q(I). Similarly, the variable name COS is changed to COZ, so as to make the library function COS available.

5. CREN

a. The variable names NACT and Q have been changed to NACTV and QV, for the reason given above.

b. Cards 2832-2837, 2878-2881, and 2911-2915 have been deleted, since they are not needed. In order to accomodate these deletions card 2910

has been amended to read

IF(NVCOS.LT.1) GO TO 208 ,

and card 2916 to read

$$P = (AS + \sqrt{AS^2 + A * QV * (1 + A) / EIN + A^2 - 1}) / (1 + A) .$$

6. AUTO

a. The multiple elastic sums are not required, and so cards 2069-2072 are deleted.

b. The quantity RATIO is not used, and has been eliminated from cards 2077 and 2078.

c. The do-loop 2083-2085 is redundant, since ELM2(J) is also set at card 2097. Cards 2083-2085 are therefore deleted, and card 2080 is modified to reflect this change.

d. The three do-loops in the sequence 2087-2098 have been combined into a single loop.

e. The variable name EXP(J) has been changed to EZP(J) to avoid conflict with the library function EXP.

f. Often it is desirable to separate an iteration procedure into two or more consecutive runs. Cards 2073-2082 have been modified to allow a run to be made using the partially corrected output of a previous run as input, instead of the experimental data. When this option is to be used the second field of the input card specifying the number of iterations, NTERM, should be non-zero. This is then followed by the cards specifying the experimental values for the angular distribution, FCVAL, which are removed from their usual position and replaced by the partially corrected output of the previous run.

g. In the case of rapidly varying angular distributions such as at 14.1 MeV, the usual iteration procedure is not as rapidly convergent as the "physical" method. For this method the calculated multiple scattering is first subtracted from the experimental input, and then the balance of the iteration (sample attenuation) is performed as usual by reflecting the output for single scattered neutrons about the input. The new coding on cards 2091-2097 reflects this change. For the first two iterations the iteration improves only the multiple scattering, the correction being applied both times to the experimental input. After two iterations multiple scattering is well enough known to allow the iterative procedure full play, so the correction is applied each time to the previous input, instead of the experimental values. This procedure yields much more rapid convergence, three iterations providing better convergence than six of the previous method.

The code has been applied to a typical associated-particle scattering geometry for 14.1 MeV neutrons on carbon. The scattering sample was a truncated cone of half-angle seven degrees, length 3.193 cm., and entrance face radius 2.059 cm., with its center located 20 cm from the neutron source. The experimental input used was simulated using published results for carbon at this energy. The results of a typical run are shown in Table 1, and on figures 8 and 10. The cross-section data used for the Monte Carlo scattering were those of Slaggie and Reynolds [8]. In Table 1 the quantity $\sigma(\theta)$ represents the true angular distribution, and the total and multiple outputs are the results calculated using it and the specified sample shape. The code varies $\sigma(\theta)$ until the total elastic output matches the experimental input to the accuracy required. The calculated angular distribution for neutrons undergoing multiple scattering is shown in figure 8. It can be seen that the result converges quite rapidly. Further iterations resulted only in statistical variations about the values obtained after three iterations; in fact, the largest change after two iterations, that at $\cos \theta = 0.625$, represents only a 1.5% change in the cross section.

In figure 9 we show the part of the correction, exclusive of multiple scattering, that depends upon sample shape. Also displayed for comparison is a calculation of the effect expected due to attenuation in the sample for an isotropic angular distribution. As can be seen, this correction is not independent of the input angular distribution. The quantity plotted, $\Delta\sigma/\sigma$, is the difference between the singly scattered output and the input, divided by the input, for each angle. This geometrical correction is contributed to about equally by varying path length in the sample and the change in the total neutron cross section with energy as a function of the neutron scattering angle. The correction varies from a few percent to about ten percent, and is directly dependent upon the accuracy with which the total neutron cross section is known over the range of energy exhibited by the neutron recoil. On the other hand, the correction for multiple scattering is larger, ranging from about 50% at the backward minimum to 4% at zero degrees, but depends mainly upon the experimental input data. Over most of the practical angular range of measurement it is a quite appreciable 15-20%. Figure 10 shows the input distribution used and the corrected output distribution obtained after three iterations of approximately 2000 interacting neutrons each. The total correction is, of course, largest in regions where the cross section exhibits minima. It varies from a few percent to a maximum of only twenty percent at measurable angles, since the geometrical and multiple scattering corrections tend to be in opposite directions, except at the extreme backward angles. Several examples of spectra showing elastic-inelastic and multiple inelastic effects in the energy spectra are given in the original references [4, 5]; the present version of MAGGIE also yields similar results.

TABLE 1

The scores for a typical Monte Carlo run for carbon at 14.1 MeV.

COUNTER NUMBER	Cos θ	$\sigma(\theta)$	TOTAL ELASTIC OUTPUT	MULTIPLE ELASTIC OUTPUT	EXPERIMENTAL INPUT
1	-1.0000	.023	.044	.021	.044
2	-.9375	.075	.092	.022	.093
3	-.8750	.093	.111	.024	.111
4	-.8125	.100	.120	.026	.121
5	-.7500	.110	.131	.028	.131
6	-.6875	.127	.148	.030	.148
7	-.6250	.152	.170	.031	.171
8	-.5625	.176	.190	.033	.191
9	-.5000	.202	.212	.034	.212
10	-.4375	.225	.231	.035	.232
11	-.3750	.240	.246	.036	.246
12	-.3125	.248	.253	.037	.254
13	-.2500	.245	.251	.037	.251
14	-.1875	.233	.240	.038	.241
15	-.1250	.216	.225	.038	.226
16	-.0625	.192	.205	.038	.205
17	.0000	.173	.183	.038	.183
18	.0625	.145	.155	.037	.155
19	.1250	.125	.135	.037	.135
20	.1875	.114	.125	.037	.125
21	.2500	.107	.123	.039	.122
22	.3125	.110	.127	.041	.126
23	.3750	.124	.142	.044	.140
24	.4375	.154	.170	.049	.167
25	.5000	.203	.215	.055	.212
26	.5625	.281	.285	.063	.281
27	.6250	.412	.404	.074	.400
28	.6875	.628	.605	.087	.600
29	.7500	.969	.929	.104	.925
30	.8125	1.486	1.435	.125	1.435
31	.8750	2.274	2.218	.151	2.219
32	.9375	3.451	3.380	.181	3.387
33	1.0000	5.197	5.045	.216	5.062

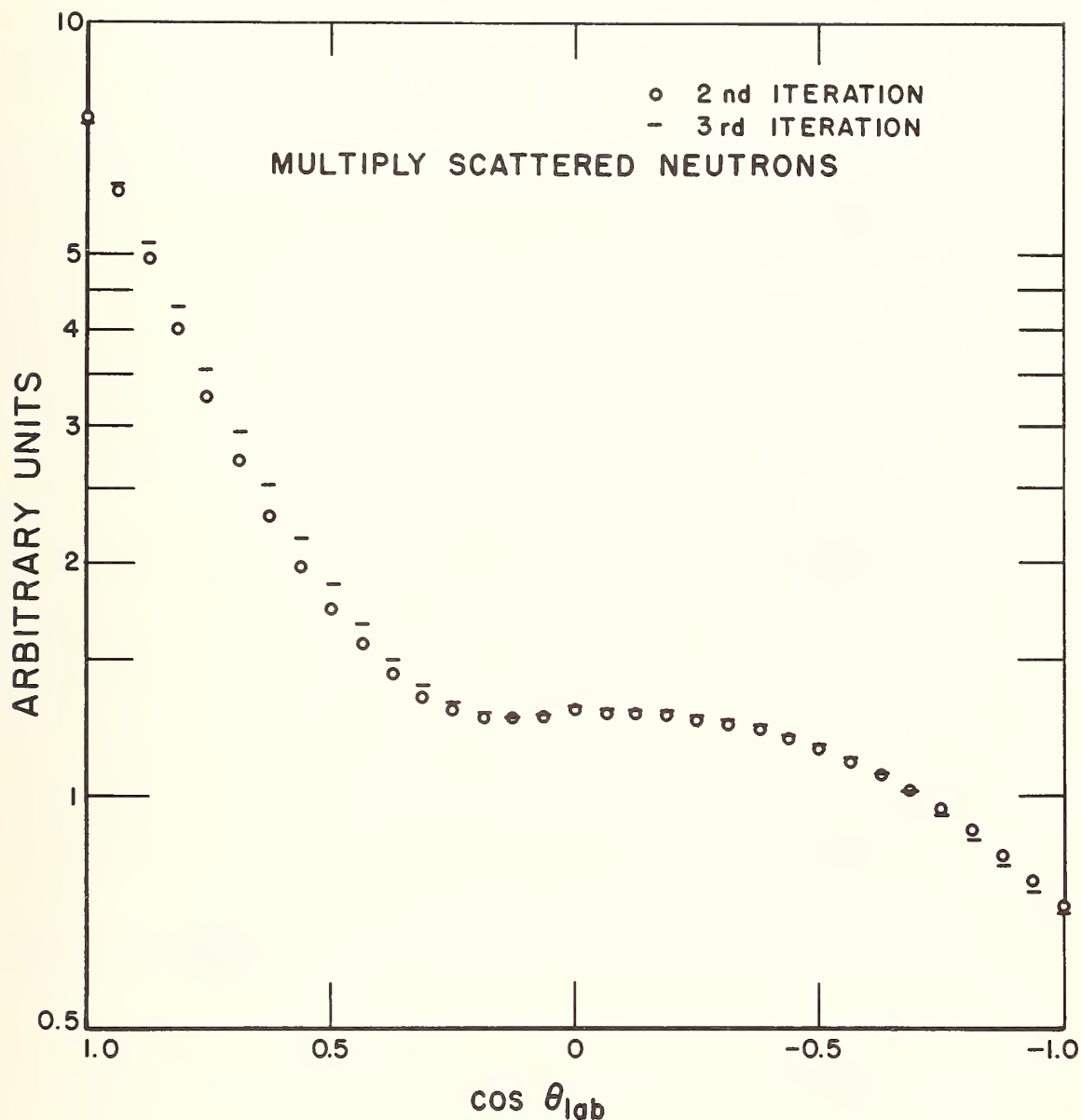


Fig. 8. The angular distribution of multiply scattered neutrons on carbon at 14.1 MeV for two successive iterations. The data for the third iteration are also given in Table 1, after suitable normalization.

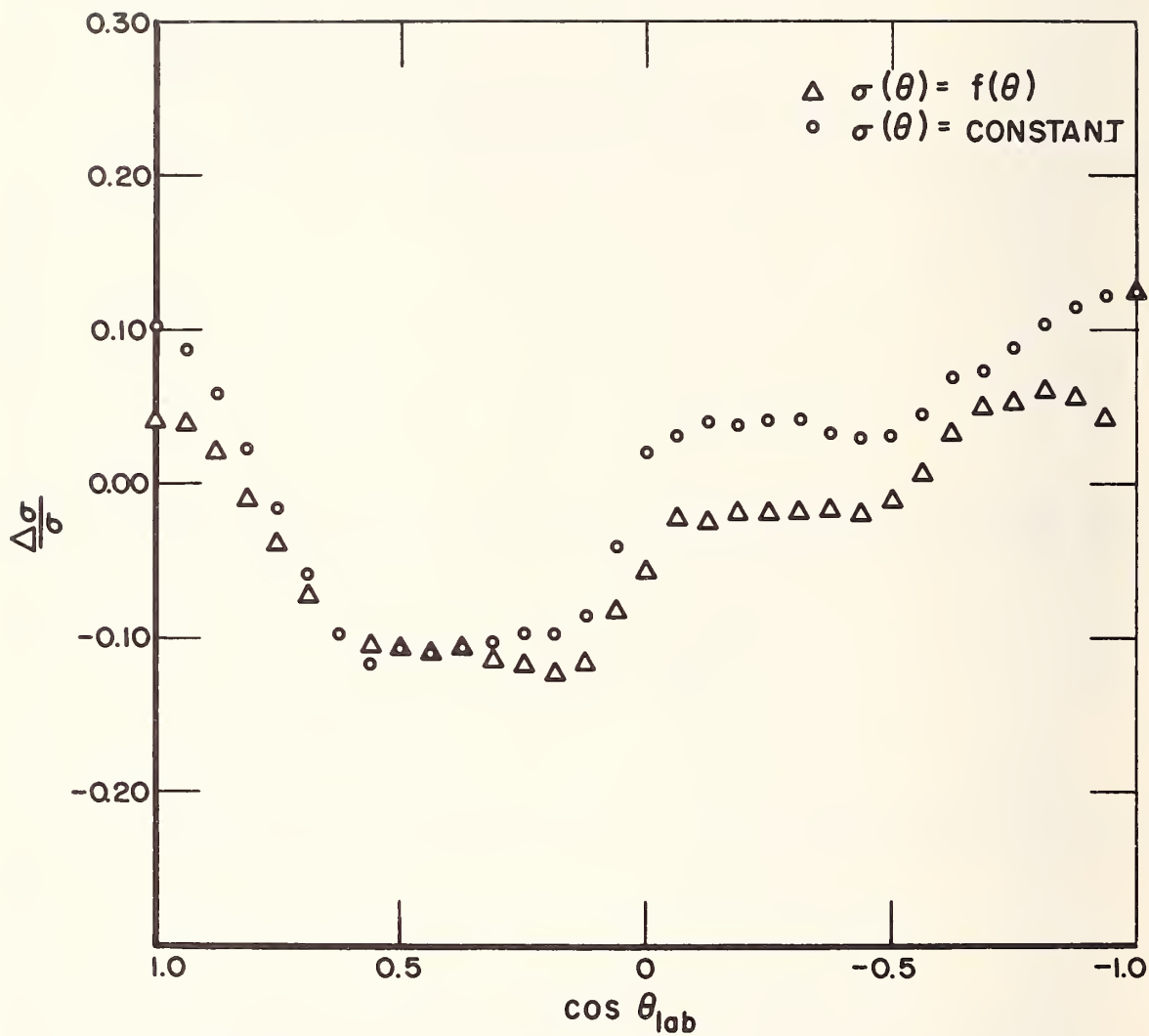


Fig. 9. The correction due to sample shape. The quantity $\Delta\sigma/\sigma$ is the change in the normalized angular distribution after subtracting multiple scattering.

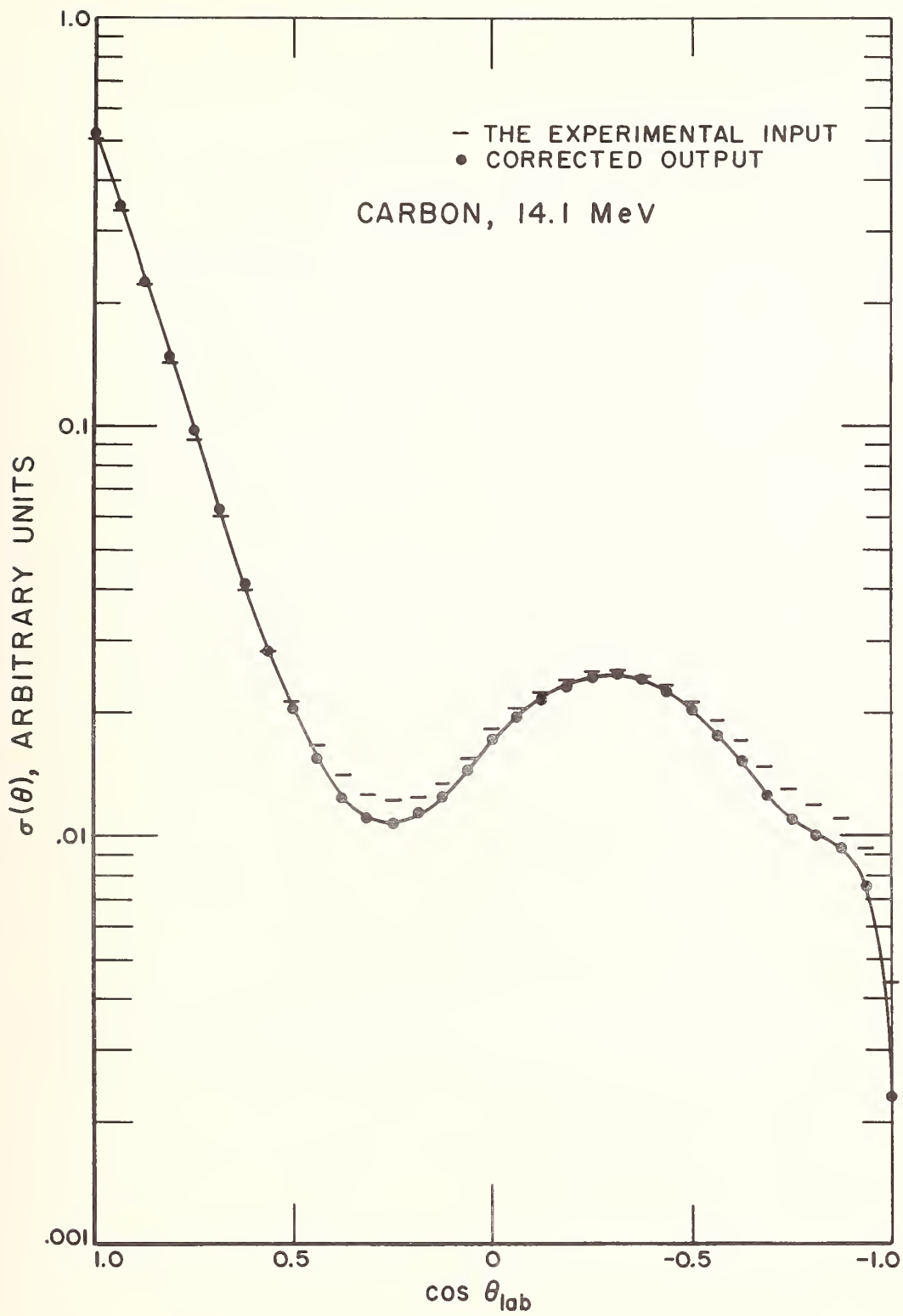


Fig. 10. The experimental input and corrected output for carbon at 14 MeV.

REFERENCES

1. Monier, L. F. C., Tripard, G. E., and White, B. L., Nucl. Instr. and Meth. 45, 282 (1966).
2. Marshak, H., Richardson, A. C. B., and Tamura, T., Phys. Rev. 150, 996 (1966).
3. A useful summary is given by M. Walt in Fast Neutron Physics, Vol. II, Marion, J. B., and Fowler, J. L., Ed. (1960).
4. Parker, J. B., Towle, J. H., Sams, D., and Jones, P. G., Nucl. Instr. and Meth. 14, 1 (1961).
5. Parker, J. B., et al, Nucl. Instr. and Meth. 30, 77 (1964).
6. Parker, K., AWRE report O-70/63; Kerr, W. M. M., AWRE report O-81/64 (1964); Miller, S. M., and Parker, K., AWRE report O-55/65.
7. Some of these coding errors have also been previously noted in private communications. We are indebted to J. B. Parker, Aldermaston, for noting the substitution of REDEN for EIN in subroutine NAPAN, as well as the errors in subroutine FPATH. H. Horstmann and H. Schmid, Geel, have noted, in addition, the difficulty involving the center of mass calculation of neutron recoil energy in subroutines NAPAN and CREN.
8. Slaggie, E. L., and Reynolds, J. T., KAPL-3099, (1966).

APPENDIX A

Listings of new and extensively revised coding.

<u>Name</u>	<u>Page</u>
CINC1	28
MAGGIE 3A	30
INPUT	36
CRNEU	44
TRACK	46
NAPAN	49
FPATH (truncated cone)	53
FPATH (cylinder)	55
OUTPUT	56
AUTO	58
RANDOM	61
RANDA	61
NRANDA	61

A complete listing is available by request from the Center for Radiation Research, Neutron Physics Section, National Bureau of Standards, Washington, D. C. 20234.

```

CINC1* FCOPY
C
C COMMON STORAGE ALLOCATION.
C
COMMON LOC,OCT
DIMENSION DATA(1),IDATA(1)
EQUIVALENCE (DATA,IDATA)
COMMON IDATA,INIT,NEGS,UMAX,XLYMIN,IXMAT,IDICE,NMATS,EMC(257),
1 PATH,SPEED,MAT,EIN,UIIN,LGROUP,ROOTE,WIN,WOUT,NSECS,
2 COSPHI(100),EOUT(100),UOUT(100),ISORAN(100),NOLAW(100),
3 NONUC(100),NOACT(100),NREACT(100),NORNGE(100),ALF,
4 ALFP(20),WMIN,WMAX
C
MAGGIE NUCLEAR DATA AIDS.
C
COMMON NUCLID,NAC,AS,ENER,NFORM,QV,ATOM1(24),Q(24),PART1(24),
1 IXACT1(24),IXNUC1,FNEGS,NUCS,NUCL(24),NACT(24),NAC1(24),
2 NANAL,NENSP,NOCOUI,NOGR,LAWREV(50),LWNO(50),NLWREV
C
BIRTH STORES AND REGISTERS.
C
COMMON BIRTH(5500),IBIRTH(1500)
C
GENERAL WORK STORES.
C
COMMON IWORK1,IWORK2,WORK3,WORK4,WORK5,WORK6,WORK7,WORK8,WORK9,
1 WORK10,WORK11,WORK12,WORK13,WORK14,WORK15,WORK16,WORK17,
2 WORK18,WORK19,ISCOLL
C
DIMENSIONS ETC.
C
COMMON HITE,FRAD,BRAD,ANGLE,TANGLE,DIST,STRTE,CNPHI(64),FLTPTH
C
CONTROL PARAMETERS.
C
COMMON IRECRG,IRGRA,IRGRB,IRGRIN,NOSAMP,SAMPLE,ANEUNO,NSENSE,IFB4
C
NBS
EU.
EU.
EU.
NBS
EU.
EU.
NBS
NBS
NBS
NBS
NBS
EU.
EU.
EU.
NBS
NBS
NBS
EU.
EU.
EU.
NBS
EU.
EU.
EU.
NBS

```



```

C MAGGIE 3A 1566
C 20.7.64 A.D.PURNELL 1570
C 9.68 MODIFIED NBS VERSION FOR FORTRAN V, A.C.B.RICHARDSON NBS1570A
C FOR TRUNCATED CONE SCATTERING SAMPLE, INCLUDING AUTOMATIC ELASTIC NBS1571
C SCATTERING CALCULATION AND GRAPHS. NBS1571A
.
INCLUDE CINC1 NBS
DIMENSION LAWTP(100) 1653
ITER=1 1654
INIT = 16719 NBS
READ(5,51)UCT NBS
CALL RANDOM(OCT) NBS
51 FORMAT(O12) NBS
READ(5,200)(HEAD(I),I=1,12) NBS
WRITE(6,200)(HEAD(I),I=1,12) NBS
200 FORMAT(12A6) NBS
CALL DATIN NBS
CALL INPUT NBS
WRITE(6,52)UCT 1656
52 FORMAT(55HUTHE STARTING VALUE FOR THE RANDOM NUMBER GENERATOR IS ,NBS
1012) NBS
103 NOSAMP=0 1658
IRECRG=0 1659
IRGRA=0 1660
IRGRB=0 1661
ANEUNO=0.0 1662
THERM=0.25000001E-7 1663
IF(IRGRA)2,7,4 1664
PRINT 3,ANEUNO 1665
3 FORMAT(1H1,47HERROR CONDITION-BIRTH REGISTER NEGATIVE ANEUNO=,F6.0 1666
2) 1667
IRGRA=0 1668
IRGRB=0 1669
GO TO 7 1670
4 IRGRA=IRGRA-11 1671
IRGRB=IRGRB-3 1672

```



```

IWORK1=IBIRTH(IRGRB+1)
MAT=IBIRTH(IRGRB+2)
IWORK2=IBIRTH(IRGRB+3)
EIN=IBIRTH(IRGRA+1)
WORK3=BIRTH(IRGRA+2)
WORK4=BIRTH(IRGRA+3)
WORK5=BIRTH(IRGRA+4)
WORK6=BIRTH(IRGRA+5)
WORK7=BIRTH(IRGRA+6)
WORK8=BIRTH(IRGRA+7)
WORK19=BIRTH(IRGRA+8)
UIN=BIRTH(IRGRA+9)
WIN=BIRTH(IRGRA+10)
WORK11=BIRTH(IRGRA+11)
IF(IRGRA-550)13,5,13
IF(IRECKG)13,13,6
IRECRG=IRECRG-11
LOC=LOC-5600
READ(37)IRECRG,(BIRTH(I),I=551,4950),(IBIRTH(I),I=151,1350)
IRGRA=IRGRA+4400
IRGRB=IRGRB+1200
GO TO 13
IF(ANEUNO-SAMPLE)10,8,8
NOSAMP=NOSAMP+1
ANEUNO=0.0
CALL OUTPUT
IF(ITMAG-1)81,82,81
81 IF(NOSAMP-JOBFIN)10,9,9
82 IF(NAC1(1)-2) 9,84,9
84 IF(NOCOU-33)9,85,9
85 CALL AUTO
IF(ITER-10) 103,9,9
9 CONTINUE
CALL END
CALL CRNEU
10 IF(NSNSE)11,13,11

```

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1673
1674
1675
1676
1677
1678
1679
1680
1681
1682
1683
1684
1685
1686
1687
1688
1689
NBS1690
NBS1692
1693
1694
1695
1696
1697
1698
1699
1700
1701
1702
1703
1704
1705
NBS1706
1715
1716
1717

```

```

11 PRINT 12 1718
12 FORMAT(I20H0SERIAL STEP EM LETH WT L M N NBS1719
2 X Y Z TIME FATE ISO. ACT. LAN. LAW NSECS) 1720
13 IF(NSENSE)14,15,14 1721
14 PRINT 33,ANEUNO,IWORK1,EIN,UIN,WIN,WORK3,WORK4,WORK5 1722
15 CALL TRACK 1723
WORK1=WIN 1724
ISCOLM=ISCOLL 1725
GO TO (17,16,22),ISCOLM 1726
PESC=PESC+WORK1 1727
GO TO 18 1728
17 CESC=CESC+WORK1 1729
18 IF(NSENSE)19,20,19 1730
19 PRINT 34,ANEUNO,IWORK1,EIN,UIN,WIN,WORK3,WORK4,WORK 1731
25,WORK10,WORK17,WORK18,WORK19,ISCOLM 1732
20 CONTINUE 1733
505 IF(IWORK1)506,506,507 1744
506 ANOCOL=ANOCOL+WORK1 1745
GO TO 1 1746
507 IF(IWORK1-1)1,508,509 NBS1747
508 I=ICOLL(1) 1748
ONCOL(I)=ONCOL(I)+WORK1 1749
GO TO 1 NBS1750
509 NOIN=0 1751
DO 511 I=1,1WORK1 1752
IF(ICOLL(I)-2)511,511,510 1753
NOIN=NOIN+1 1754
LAWTP(NOIN)=LAWTP(I) 1755
CONTINUE 1756
511 IF(NOIN-2)518,512,518 1757
512 IF(LAWTP(1)-LAWTP(2))513,514,514 1758
513 I=LAWTP(2) 1759
J=LAWTP(1) 1760
GO TO 515 1761
514 I=LAWTP(1) 1762

```

```

515 J=LAWTP(2)
516 IF(I-5)516,516,518
517 IF(J-5)517,517,518
518 TOCOL(I,J)=TOCOL(I,J)+WORK1
519 IF(ICOLL(1)-2)524,519,519
520 J=ICOLL(1)-1
521 NOIN = 0
522 DO 521 I=1,IWORK1
523 IF(ICOLL(I)-2)521,520,520
524 NOIN=NOIN+1
525 CONTINUE
526 IF(NOIN-1)522,522,523
527 DEGENE(J)=DEGENE(J)+WORK1
528 GO TO 1
529 DEGENI(J)=DEGENI(J)+WORK1
530 GO TO 1
531 NOIN=0
532 DO 526 I=1,IWORK1
533 IF(ICOLL(I)-2)526,525,525
534 NOIN=NOIN+1
535 NOIN1=ICOLL(I)
536 CONTINUE
537 IF(NOIN-1)527,529,528
538 ALLEL=ALLEL+WORK1
539 GO TO 1
540 EVDGEN=EVDGEN+WORK1
541 GO TO 1
542 J=NOIN1-1
543 EDGEN(J)=EDGEN(J)+WORK1
544 GO TO 1
C COLSN
22 IWORK1=IWORK1+1
CALL CR
IF(IFB4)23,24,23
1763
1764
1765
1766
1767
1768
1769
1770
1771
1772
1773
1774
1775
NBS1776
1777
NBS1778
1779
1780
1781
1782
1783
1784
1785
1786
NBS1787
1788
NBS1789
1790
1791
NBS1792
1794
1795
1796
1797
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```

23 CALL CALC
24 CALL NAPAN
   WORK1=WOUT
25 IF(IWORK1-100)25,25,26
531 IF(NREACT(1)-2)532,531,532
   ICOLL(IWORK1)=1
5310 IF(IWORK1-1)26,5310,26
   ALAM(1)=ALAM(1)+WORK1
   GO TO 26
532 IF(NREACT(1)-16)534,533,533
533 ICOLL(IWORK1)=2
   IF(IWORK1-1)26,5330,26
5330 ALAM(2)=ALAM(2)+WORK1
   GO TO 26
534 DO 535 I=1,NLWREV
   IF(NOLAW(1)-LAWREV(I))535,536,535
535 CONTINUE
   ICOLL(IWORK1)=3
   LAWTP(IWORK1)=1
   IF(IWORK1-1)26,5350,26
5350 ALAM(3)=ALAM(3)+WORK1
   GO TO 26
536 ICOLL(IWORK1)=LWNO(I)+3
   LAWTP(IWORK1)=LWNO(I)+1
   IF(IWORK1-1)26,5360,26
5360 ALAM(I+3)=ALAM(I+3)+WORK1
26 IF(NSECS)32,507,27
27 WORK10=WOUT
   DO 31 I=1,NSECS
   WORK12=WOUT(I)
   WORK9=WOUT(I)
   IF(WORK12-THERM)28,29,29
28 WORK12=THERM
   WORK9=LOG(THERM/0.25E-7)
29 ANGLE=COSPHI(I)
   CALL TWIST(ANGLE)

```

NBS1833
1834
1835

```

30 CALL WRIBT
   IF(NSENSE) 30, 31, 30
   PRINT
      2WORK14, WORK15, WORK16, WORK17, WORK18, WORK19, ISCOLM, NONUC(I), NOACT(I)
      3,NORNGE(I), NOLAW(I), NSECS
31 CONTINUE
32 GO TO 1
33 FORMAT(1H0, F7.0, I4, 9F7.3, F10.4, I5, 3I6, I5, I7)
34 FORMAT(1X, F7.0, I4, 9F7.3, F10.4, I5, 3I6, I5, I7)
   END
1836
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1838
1839
1840
1841
1842
1843
1844
1845

```



```

0123
0204
0205 NBS0205
0206 NBS0206
      NBS
NBS0207
NBS0219
NBS0224
NBS0225
NBS0226
NBS0226A
NBS0226B
NBS0226C
0227
0228
0229
0230
0231
0232
0233
0234
0235
0236
0237
0238
0239
0240
0241
0242
0243
0244
0245 NBS0245
0247
0248
0249 NBS0249
0250
SUBROUTINE INPUT
GENERAL INPUT FOR TRUNCATED CONE, INCIDENT NEUTRONS ALONG THE X AXIS
      INCLUDE CINCI
      DIMENSION STDIST(25), ANGS(25), AREA(25), NOACA(24)
      READ(5,201) IFB4, NSENSE
      READ(5,202) HITE, FRAD, ANGLE
      WRITE(6,204) HITE, FRAD, ANGLE
204  FORMAT(46H) THE SCATTERER IS A TRUNCATED CONE OF LENGTH =,F7.3,
      127HCMS, ENTRANCE FACE RADIUS =,F7.3,21HCMS, AND HALF ANGLE =,
      2F7.5,8HRADIANS.)
      HITE=HITE/2.
      READ(5,201) JOBFIN, ITMAG
      READ(5,202)
      WRITE(6,2040)
      SAMPLE
      JOBFIN, SAMPLE
2040  FORMAT(14H) THERE WILL BE, I3,16H SAMPLES EACH OF, F8.0,10H NEUTRONS.
      2)
2042  IF (ITMAG-1) 2043,2042,2043
2041  WRITE(6,2041)
2041  FORMAT(1H ,31H AUTOMATIC ITERATION REQUESTED.,10X,17H ONE SAMPLE 0
      1NLY.)
      JOBFIN=1
2043  READ(5,202) STRTE
      IF(STRTE) 205,206,205
205  WRITE(6,207)
      STRTE
207  FORMAT(21H) STARTING ENERGY IS =,F8.3,4H MEV)
      GO TO 208
206  WRITE(6,209)
209  FORMAT(41H) STARTING ENERGY IS FROM FISSION SPECTRUM.)
208  READ(5,202)
      THETAM=COS(ATAN(FRAD/(-DIST-HITE)))
      WRITE(6,210) THETAM, DIST
210  FORMAT(71H) COSINE OF SEMI-VERTICAL ANGLE OF THE CONE JUST ENCLOSED
      26 THE SAMPLE =,F8.6,12H FROM SOURCE, F6.2,22H CMS ALONG THE X AXIS
      3.)

```

```

0251 READ(5,201)NSTDP
0252 READ(5,202)(ANGS(I),STDIST(I),I=1,NSTDP)
0253 WRITE(6,2010)(ANGS(I),STDIST(I),I=1,NSTDP)
0254 FORMAT(31H0START DISTRIBUTION PROVIDED IS /9(6F10.4/))
0255 FORMAT(6I10)
0256 FORMAT(6F10.4)
0257 C DEFINE START DISTRIBUTION TO INCLUDE THE MINIMUM ANGLE WHICH WILL
C JUST ENCLOSE THE SAMPLE (USUALLY ABOUT 10-16 DEGREES) BY UP TO 25 NBS0258
C POINTS.
0259 DO 211 I=1,NSTDP
0260 ANGS(I)=COS (ANGS(I)/57.295828)
0261 NBS0261
0262 CONTINUE
0263 DO 212 I=1,NSTDP
0264 IF(THETAM-ANGS(I))212,213,213
0265 CONTINUE
0266 WRITE(6,2120)
0267 FORMAT(41HINSUFFICIENT START DISTRIBUTION PROVIDED.)
0268 CALL EXIT
0269 AREA(1)=0.0
0270 STDIST(I)=STDIST(I-1)-(STDIST(I-1)-STDIST(I))*(ANGS(I-1)-THETAM)/(
0271 2ANGS(I-1)-ANGS(I))
0272 ANGS(I)=THETAM
0273 LI=I-1
0274 LI=I-1
0275 DO 214 K=1,L
0276 AREA(K+1)=AREA(K)+((STDIST(K+1)+STDIST(K))/2.0)*(ANGS(K)-ANGS(K+1)
0277 2)
0278 CONTINUE
0279 DO 215 K=1,64
0280 AI=K
0281 SEG=(AREA(I)*(2.0*AI-1.0))/128.0
0282 DO 216 L=1,I
0283 IF(AREA(L)-SEG)216,217,218
0284 CONTINUE
0285 WRITE(6,2160)
0286 FORMAT(63H1ERROR IN INPUT. PARTIAL SUM LRSS THAN THE TOTAL. CALL

```



```

0315  AVCOS=(CNPHI(32)+CNPHI(33))/2.0
0316  DO 2150 I=1,LI
0317  IF(ANGS (I+1)-AVCOS)2151,2152,2150
0318  2150 CONTINUE
0319  I=LI
0320  2151 FACT=STDIST(I)-(STDIST(I)-STDIST(I+1))*((ANGS (I)-AVCOS)/(ANGS (I)
0321  2-ANGS (I+1)))
0322  GO TO 2153
0323  2152 FACT=STDIST(I)
0324  2153 FLUXFT=FLUXFT*STDIST(I)/FACT
0325  READ(5,201)NLWREV
0326  READ(5,220)(LAWREV(I),LWNO(I),I=1,NLWREV)
0327  WRITE(6,2193)
0328  2193 FORMAT(38H0LAW REFERENCE NUMBER M.C.LAW NUMBER.)
0329  WRITE(6,2194)(LAWREV(I),LWNO(I),I=1,NLWREV)
0330  2194 FORMAT(I10,I0X,I10)
0331  IF(IFB4)2192,234,2192
0332  2192 READ(5,201)IANALA,IANALB,IANALC,IANALD
0333  READ(5,201)NANAL
0334  READ(5,220) (NUCL(I),NAC1(I),I=1,NANAL)
0335  220 FORMAT(2I10)
0336  WRITE(6,221)NANAL
0337  221 FORMAT(15H054 ANALYSIS ON,I3,10H ACTIONS. )
0338  WRITE(6,222) (NUCL(I),I=1,NANAL)
0339  222 FORMAT(15H0NUCLIDE NOS. ,24I4)
0340  WRITE(6,223) (NAC1(I),I=1,NANAL)
0341  223 FORMAT(16H0K.P.ACTION NOS.,I3,23I4)
0342  M3=M1-NMATS
0343  IXNUC1=IDATA(M3)
0344  M4=M3+4*NMATS
0345  NUCS=IDATA(M4)
0346  DO 3 I=1,NANAL
0347  M5=IXNUC1+NUCL(I)
0348  M6=8*NUCS+M5
0349  NOACA(I)=IDATA(M6)
0350  IXACT1(I)=IDATA(M5)
0351  NBS0345
0352  0346
0353  0347
0354  0348
0355  0349
0356  0350
0357  0351
0358  0352

```



```

MMM=2*NOACA(I)+IXACT1(I)      0359
K=1                             0360
NOACA1=NOACA(I)                0361
DO 4 J=1,NOACA1                0362
L=MMM+K                         0363
IF(IDATA(L)-NAC1(I))7,5,7      0364
K=K+1                           0365
CONTINUE                        0366
WRITE(6,6)                      NAC1(I) 0368
CALL EXIT                       0369
NACT(I)=K                       0370
CONTINUE                        0371
FORMAT(18H K.P.ACTION NUMBER,I4,29H DOES NOT APPEAR IN THE DATA.)
WRITE(6,224)                    (NACT(I),I=1,NANAL) 0372
FORMAT(15HUM.C.ACTION NO.,24I4) 0373
READ(5,202)(FCCOS(I),I=1,33)    0374
33 COSINES WITH 33 VALUES OF THE ANGULAR DISTRIBUTION MUST BE
C SUPPLIED, WITH MONOTONICALLY INCREASING VALUES OF THE COSINE.
DO 225 J=1,NANAL                NBS0376
READ(5,202)(FCVAL(J,I),I=1,33) 0378
AREB=0.0                        0379
DO 228 I=1,32                  0380
AREB=AREB+ABS (FCCOS(I+1)-FCCOS(I))*((FCVAL(J,I)+FCVAL(J,I+1))/2.0NBS0381
2)                               0382
228 CONTINUE                    0383
IA=1                            0384
IB=15                           0385
2280 WRITE(6,226)(FCCOS(I),I=IA,IB) 0386
226 FORMAT(15H0COSINES          ,15F7.3)
WRITE(6,227)(FCVAL(J,I),I=IA,IB) 0388
227 FORMAT(15H DISTRIBUTION    ,15F7.3)
DO 229 I=IA,IB                  0390
FCVAL(J,I)=FCVAL(J,I)/AREB
229 CONTINUE                    0391
WRITE(6,230)(FCVAL(J,I),I=IA,IB) 0392
230 FORMAT(15H UIFF. X SECT.  ,15F7.3) 0394

```

```

0395 IF (IA-1)2300,2300,2301
0396 IA=16
0397 IB=30
0398 GO TO 2280
0399 IF (IA-16)2302,2302,225
0400 IA=31
0401 IB=33
0402 GO TO 2280
0403 CONTINUE
0404 IMINM=1
0405 DO 252 L=1,NANAL
0406 IMIN(L)=IMINM
0407 READ(5,250)IMAX(L)
0408 FORMAT(I10)
0409 PRINT 251,IMAX(L),L
0410 FORMAT(10H)HERE ARE,I2,46H SUPPLEMENTARY RANGES ASSOCIATED WITH ANBS0410
0411 ANALYSIS,I2)
0412 IF (IMAX(L))252,252,253
0413 IMAXM=IMAX(L)-1+IMINM
0414 DO 254 I=IMINM,IMAXM
0415 READ(5,202)ENVAL(L,I)
0416 READ(5,202) (SUPVAL(I,K),K=1,33)
0417 AREB=0.0
0418 DO 258 J=1,32
0419 AREB=AREB+ABS (FCCOS(J+1)-FCCOS(J))*((SUPVAL(I,J)+SUPVAL(I,J+1))/2NBBS0419
0420 2.0)
0421 CONTINUE
0422 IA=1
0423 IB=15
0424 WRITE(6,2580)ENVAL(L,I)
0425 FORMAT(25H)THIS RANGE APPLIES BELOW,F7.4,5H MEV.)
0426 WRITE(6,226) (FCCOS(J),J=IA,IB)
0427 WRITE(6,227) (SUPVAL(I,J),J=IA,IB)
0428 DO 260 J=IA,IB
0429 SUPVAL(I,J)=SUPVAL(I,J)/AREB
0430 CONTINUE

```

```

261 WRITE(6,230) (SUPVAL(I,J),J=IA,IB)
IF(IA-1)261,261,262
IA=16
IB=30
GO TO 259
262 IF(IA-16)263,263,254
263 IA=31
IB=33
GO TO 259
254 CONTINUE
IMIN=IMAXM+1
IF(IMAXM-24)252,252,256
256 WRITE(6,257)
257 FORMAT(27H11MAX HAS EXCEEDED 24-HALT.)
CALL EEXIT
CONTINUE
252 READ(5,201)NOCOU
READ(5,202)
WRITE(6,232)
232 FORMAT(10H0THERE ARE,I3,56H COUNTERS LOCATED AROUND THE EQUATOR OFNBS0450
2 THE SCATTERER AT/(10F9.4))
NENSP
READ(5,201)
NENSP(I),I=1,NENSP)
WRITE(6,233)
233 FORMAT(48H0THE ENERGY SPECTRA OF THE B4 ANALYSIS CELLS ARE/(10F9.3
2))
234 DO 2340 K=1,32
FCCOS(K)=FCCOS(K+1)-FCCOS(K)
2340 CONTINUE
RETURN
END
0431
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NBS0441
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NBS0451
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0463

```

3036		
	SUBROUTINE CRNEU	
	TRUNCATED CONE SCATTERING SAMPLE, INCIDENT NEUTRONS ALONG THE AXIS	NBS
		3116
	INCLUDE CINCL	NBS3121
	DIMENSION TEMP(10),TRIG(4)	3122
	ITEMP=NRANDA(64)	EU.3124
	WORK3=CNPHI(ITEMP)	EU.3126
	TEMP(1)=RANDA(-1)	3127
	TEMP(2)=RANDA(-2)	3128
	TEMP(3)=TEMP(1)*TEMP(1)	3129
	TEMP(4)=TEMP(2)*TEMP(2)	3130
	TEMP(5)=TEMP(3)+TEMP(4)	3131
	IF(TEMP(5))62,62,63	3132
	IF(TEMP(5)-1.0)64,62,62	3133
	TRIG(1)=(TEMP(3)-TEMP(4))/TEMP(5)	NBS3134
	TRIG(2)=(2.0*TEMP(1)*TEMP(2))/TEMP(5)	3135
	TRIG(3)=SQRT(1.0-WORK3*WORK3)	3136
	WORK4=TRIG(1)*TRIG(3)	3137
	WORK5=TRIG(2)*TRIG(3)	3138
	DOES THE NEUTRON HIT THE TARGET	3139
	TEMP(1)=FRAD+(DIST+HITE)*TRIG(3)/WORK3	NBS3140
	IF(TEMP(1))67,67,65	3142
	TEMP(2)=-((DIST+HITE)/WORK3	NBS3143
	WORK6=-HITE	NBS3145
	WORK7=WORK4*TEMP(2)	NBS3146
	WORK8=WORK5*TEMP(2)	NBS3147
	ANEUNO=ANEUNO+1.0	3148
	EIN=STRTE	3149
	UIN=LOG(EIN/0.25E-7)	NBS3150
	MAT=1	NBS3151
	IWORK1=0	3152
	WORK19=0.0	3153
	WIN=1.0	3154

3155
3156
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67
WORK11=0.0
IWORK2=0
RETURN
AMISS=AMISS+1.0
GO TO 61
END

```

SUBROUTINE TRACK
C TRACKS NEUTRONS IN A TRUNCATED CONE WITH AXIS LYING ALONG THE NEUTRON BEAM
INCLUDE CINC1
CALL EGMV
PATHL=(-(LOG(RANDA(1))))*PATH
IF(WORK3)1,6,2
1 TRS=(HITE+WORK6)/WORK3
RAD2=FRAD**2
GO TO 3
2 TRS=(HITE-WORK6)/WORK3
RAD2=BRAD**2
3 WORK17=WORK7+TRS*WORK4
WORK18=WORK8+TRS*WORK5
TRR2=WORK17**2+WORK18**2
IF(RAD2.LT.TRR2)GO TO 6
IF(TRS.GT.PATHL)GO TO 13
ISCOLL=2
IF(WORK3.GT.0)GO TO 4
WORK16=-HITE
GO TO 5
4 WORK16=HITE
5 WORK19=WORK19+TRS/SPEED
RETURN
6 CALL FONEC(PATHL)
IF(HITE.LE.ABS(WORK16))GO TO 7
TRA2=WORK17**2+WORK18**2
RAD2=(FRAD+(HITE+WORK16)*TANGLE)**2
IF(RAD2.GT.TRA2)GO TO 14
7 ISCOLL=1
NRTS=0
RAD=FRAD+(HITE+WORK6)*TANGLE
TA=1.-WORK3**2*(1.+TANGLE**2)
TB=WORK7*WORK4+WORK8*WORK5-WORK3*TANGLE*RAD

```



```

TC=WORK7**2+WORK8**2-RAD**2
IF(ABS(TA).GT.1.E-30)GO TO 8
IF(ABS(TB).LT.1.E-30)GO TO 21
TR=-TC/(2.*TB)
GO TO 12
8 TD=TB*TB-TA*TC
IF(TD)15,10,9
9 TE=-TB+SQRT(TD)
IF(TE.LT.0)GO TO 10
NRTS=1
TR=TE/TA
10 TE=-TB-SQRT(TD)
IF(TE.LT.0)GO TO 11
NRTS=NRTS+1
TR=TE/TA
11 IF(NRTS-1)17,12,19
12 CALL FONEC(TR)
WORK19=WORK19+TR/SPEED
RETURN
13 CALL FONEC(PATHL)
14 ISCOLL=3
WORK19=WORK19+PATHL/SPEED
RETURN
C ERROR EXITS
15 WRITE(6,16)
16 FORMAT(24H IMAGINARY ROOT IN TRACK)
GO TO 23
17 WRITE(6,18)
18 FORMAT(26H NO POSITIVE ROOT IN TRACK)
GO TO 23
19 WRITE(6,20)
20 FORMAT(28H TWO POSITIVE ROOTS IN TRACK)
GO TO 23

```

```
21 WRITE(6,22)
22 FORMAT(20H INDETERMINATE ROOT IN TRACK)
23 WRITE(6,24) WORK3,WORK4,WORK5,WORK6,WORK7,WORK8,PATHL
24 FORMAT(7F10.3)
   CALL EXIT
   END
```

```

SUBROUTINE NAPAN
3371
INCLUDE CINCI
DIMENSION FACT(33),PATHLL(33),ANGLEE(33),KL(33)
NENSP1=NENSP-1
M1=IXMAT+NMATS+MAT
N=M1+4*NMATS
MISSM=IDATA(N)
M2=IDATA(M1)-MISSM
DO 1 I=1,NOCCU
COSOM=COUNT(I)
CALL FPATH(COSOM,PATHL,ANGLEF)
DO 2 K=1,32
IF(FCCOS(K+1)-ANGLEF)2,3,3
CONTINUE
2
K=32
AFCOS=FCCOS(K)
FACT(I)=(ANGLEF-AFCOS)/FCCOSD(K)
PATHLL(I)=PATHL
ANGLEE(I)=ANGLEF
KL(I)=K
CONTINUE
1
DO 7 J=1,NANAL
NFORM=0
NAC=NAC1(J)
NUCLID=NUCL(J)
PART=PART1(J)
ATOM=ATOM1(J)
WSQ=ATOM*ATOM-1.0
DO 6 I=1,NOCCU
K=KL(I)
AS=ANGLEE(I)
FACTOR=FACT(I)
PATHL=PATHLL(I)
ASQ=AS*AS
IF(PART)7,7,8

```

NBS
NBS3451
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NBS3460
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NBS3462
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NBS3466
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NBS3468
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3483
3484

```

8   IF(NAC-2)9,10,9
10  X=SQRT (ASQ+WSQ)
    P=((AS+X)/(ATOM+1.0))*((AS+X)/(ATOM+1.0))
    REDEN=EIN*P
    GO TO 11
9   IF(NFORM)34,34,35
34  CALL CREN
    REDEN=EINER
    IF(REDEN)38,11,11
35  GO TO (36,11,11,11,11,11,11,11,11,11,11,11,11,11,11,11),NFORM
36  P=(AS+SQRT (ASQ+ATOM*QV*(1.0+ATOM)/EIN+WSQ))/(1.0+ATOM)
    REDEN=EIN*P*P
    11 IF(IMAX(J).LE.0)GO TO 15
        IMINM=IMIN(J)
        IMAXM=IMINM+IMAX(J)-1
        DO 50 JJ=IMINM,IMAXM
        IF(EIN-ENVAL(J,JJ))50,51,52
    50 CONTINUE
    51 KK=JJ
    52 IF(JJ.LL.IMINM)GO TO 15
        KK=JJ-1
        GO TO 19
15  SIGMA=(FCVAL(J,K+1)-FCVAL(J,K))*FACTOR+FCVAL(J,K)
    GO TO 16
19  SIGMA=(SUPVAL(KK,K+1)-SUPVAL(KK,K))*FACTOR+SUPVAL(KK,K)
16  DO 20 L=1,NENSP1
    IF(ENSP(L+1)-REDEN)20,21,21
20  CONTINUE
    L=NENSP1
    S=REDEN
    LGR=1
21  IF(S-EMC(1))5000,5000,4999
4999 IF(NEGS-128)4001,4001,5001
5001 M=129
    LGR=M

```

```

3485
NBS3486
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3492
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NBS3494
NBS3495
3496
NBS3503
NBS3504
NBS3505
NBS3506
NBS3507
NBS3508
NBS3509
NBS3510
NBS3511
NBS3512
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3516
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3520
EU.
EU.
NBS EU.
EU.
EU.
EU.NBS

```



```

4200 LGR=M
      IF(S-EMC(M))4210,5000,5000
4210 LGR=M-1
5000 M3=M2+LGR
      PATHB=(-DATA(M3))
      QU=Q(J)
      PTH=PATHL/PATHB
      ADD=QU*SIGMA*EXP(PTH)
      DEAI,L)=DEAI,L)+ADD
      IF(NAC.NE.2)GO TO 24
      IF(IWORK1.LE.1)GO TO 23
      III=IWORK1-1
      DO 22 II=1,III
      IF(ICOLL(II).NE.1)GO TO 24
22 CONTINUE
23 DEB(I,L)=DEB(I,L)+ADD
      IF(IWORK1.LE.1)GO TO 6
      DED(I,L)=DED(I,L)+ADD
      GO TO 6
24 DEC(I,L)=DEC(I,L)+ADD
6 CONTINUE
7 CONTINUE
      GO TO 40
38 PRINT 39,REDEB
39 FORMAT(IH0,F10.4,32HNEGATIVE ENERGY COMPUTED IN CREN)
40 CALL EXIT
      RETURN
      END

```

```

EU. 3524
EU. 3525
EU. 3526
EU. 3527
NBS3529
NBS3530
NBS3531
NBS3532
NBS3533
NBS3534
NBS3535
NBS3537
NBS3539
NBS3540
      3541
NBS3543
      3544
      3545
      3546
      3547
      3548
      3549
      3550
      3551

```



```

SUBROUTINE FPATH(COSOM,PATHL,ANGLEF)
C PATHLENGTH AND ANGLE FOR B4 ANALYSIS, FOR TRUNCATED CONE SAMPLE.
C AXIS ALONG THE INCIDENT NEUTRON BEAM.
INCLUDE CINCI
C NEUTRON OUT AT 90 DEGREES.
IF(ABS(COSOM).GE.1.E-5)GO TO 1
ANGLEF=WORK4
TSQ=TANGLE**2
XPR=FRAD/TANGLE+HITE+WORK16
PATHL=SQRT(XPR*XPR*TSQ-WORK18**2)-WORK17
RETURN
C NEUTRON OUT ENTRANCE FACE
1 SINOM=SQRT(1.-COSOM**2)
ANGLEF=COSOM*WORK3+SINOM*WORK4
IF(COSOM.GT.0)GO TO 2
IF(FRAD.LT.ABS(WORK18))GO TO 3
FTAN=(SQRT(FRAD**2-WORK18**2)-WORK17)/(HITE+WORK16)
TANOM=SINOM/COSOM
IF(FTAN+TANOM.LT.0)GO TO 3
PATHL=- (HITE+WORK16)/COSOM
RETURN
C NEUTRON OUT EXIT FACE
2 BTAN=(SQRT(BRAD**2-WORK18**2)-WORK17)/(HITE-WORK16)
TANOM=SINOM/COSOM
IF(BTAN.LT.TANOM)GO TO 3
PATHL=(HITE+WORK16)/COSOM
RETURN
C NEUTRON OUT THE SIDE, BUT NOT AT 90 DEGREES, NORMAL SOLUTION.
3 TSQ=TANGLE**2
XPR=FRAD/TANGLE+HITE+WORK16

```

```
PA=SIN0w**2-(COS0M**2)*TSQ  
PB=WORK17*SIN0M-XPR**COS0M*TSQ  
PC=WORK17**2-(XPR**2)*TSQ+WORK18**2  
IF(AHS(PA).LT.1.E-30)GO TO 4  
PATHL=(SQRT(PB**2-PA*PC)-PB)/PA  
RETURN
```

```
C LINEAR SOLUTION  
4 PATHL=-PC/(2.*PB)  
RETURN  
END
```

```

C
C SUBROUTINE FPATH(COSOM,PATHL,ANGLEF)
C PATHLENGTH AND ANGLE FOR THE B4 ANALYSIS, FOR A CYLINDRICAL SAMPLE
C WITH ITS AXIS ALONG THE INCIDENT NEUTRON BEAM.
C
      INCLUDE CINC1
      PATHL=SQRT(FRAD**2-WORK18**2)-WORK17
      IF(ABS(COSOM).GT.1E-5)GO TO 1
      ANGLEF=WORK4
      GO TO 3
1  SINOM=SQRT(1-COSOM**2)
   TANOM=SINOM/COSOM
   ANGLEF=COSOM*WORK3+SINOM*WORK4
   IF(COSOM.GT.1E-5)GO TO 2
   FTAN=PATHL/(HITE+WORK16)
   IF((FTAN+TANOM).LT.1E-5)GO TO 3
   PATHL=- (HITE+WORK16)/COSOM
   RETURN
2  FTAN=PATHL/(HITE-WORK16)
   IF(FTAN.LT.TANOM)GO TO 3
   PATHL=(HITE+WORK16)/COSOM
   RETURN
3  PATHL=PATHL/SINOM
   RETURN
END

```

```

SUBROUTINE OUTPUT
OUTPUTS RESULTS AT THE END OF EACH SAMPLE.
C
INCLUDE CINC1
DIMENSION P(33)
PRINT
1 1,(HEAD(I),I=1,12)
   2,NOSAMP
2  FORMAT(1H0,52X,13HSAMPLE NUMBER,I3)
   3,SAMPLE
3  PRINT
   FORMAT(1H0,40X,F8.0,33H NEUTRONS STARTED IN EACH SAMPLE.)
   PART =NOSAMP
   TOT=PART*SAMPLE
   PRINT
   5,TOT
5  FORMAT(1H0,39X,F9.0,35H NEUTRONS HAVE BEEN TRACKED SO FAR.)
   PRINT
   4,PESC,CESC
4  FORMAT(24HUPANE SURFACE ESCAPES =,F11.3/24HOCURVED SURFACE ESCAPE
25=,F11.3)
CALL OUT83
I1=ENSP-1
IF(IFB4.EQ.0)GO TO 7
CALL OUT84(DEA,1,IANALA)
IF(IANALA.LE.0)GO TO 8
DO 6 I=1,I1
P(I)=DEA(IANALA,I)
CALL APL0T3(P,I1,ENSP,IANALA)
8  CALL OUT84(DEB,2,IANALB)
   IF(IANALB.LE.0)GO TO 10
DO 9 I=1,I1
P(I)=DEB(IANALB,I)
9  CALL APL0T3 (P,I1,ENSP,IANALB)
10 CALL OUT84(DEC,3,IANALC)
   IF(IANALC.LE.0)GO TO 12
DO 11 I=1,I1
P(I)=DEC(IANALC,I)
11 CALL APL0T3 (P,I1,ENSP,IANALC)

```

```

3878
3879
NBS
3958
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3972
3973
3974
NBS3975
NBS3937
NBS3978
NBS3979
NBS3980
3982
NBS3984
NBS3985
NBS3986
NBS3987
3989
NBS3991
NBS3992
NBS3993
NBS3994
3996

```

```
12 CALL OUIB4(JED,4,IANALD)
   IF(IANALD.LE.0)GO TO 7
   DO 13 I=1,11
   P(I)=DEU(IANALD,I)
   CALL APLOTS (P,I,ENSP,IANALD)
7 RETURN
  END
```

```
NBS3998
NBS3999
NBS4000
NBS4001
  4003
  4004
  4005
```

1977
1978
1979
1980
1981

SUBROUTINE AUTO
C TO ALLOW AUTOMATIC ITERATION UNDER THE FOLLOWING CONDITIONS.
C NUMBER OF COUNTERS AND INPUT POINTS = 33.
C ONE NUCLIDE ONLY.
C ELASTIC CORRECTION ONLY.
C PHYSICAL ITERATION.
C (DELTA SIGMA)/(SIGMA) IS THE CHANGE IN THE CROSS SECTION EXCLUDING
C MULTIPLE SCATTERING.

```

INCLUDE CINCL
DIMENSION TOTEL(33),CHI(33),ELM2(33),XIP3(33)
SUMEL=(DEB(1,33)+DEB(33,33))/32.0
DO 4 J=2,32
SUMEL=SUMEL+DEB(J,33)/16.0
4 CONTINUE
IF(ITER.NE.1)GO TO 500
READ(5,101)NTERM,IFCONT
101 FORMAT(2I10)
NTERM=NTERM-1
IF(IFCONT.EQ.0)GO TO 100
READ(5,104)(LZP(I),I=1,33)
104 FORMAT(6F10.4)
SUMEZP=(EZP(1)+EZP(33))/32
DO 107 J=2,32
107 SUMEZP=SUMLZP+EZP(J)/16
DO 108 J=1,33
108 EZP(J)=EZP(J)/SUMEZP
GO TO 102
100 DO 105 I=1,33
LZP(1)=FCVAL(1,I)
105 CONTINUE
102 IF(NTERM.LE.1)GO TO 300
106 READ(5,103)(SAMP(I),I=1,NTERM)
103 FORMAT(6F10.2)
300 CHISO=0.0

```

NBS
NBS2062
2065
2066
2067
2068
NBS2073
NBS2077
NBS2078
2079
NBS
NBS
NBS
NBS
NBS
NBS
NBS
NBS
NBS
2074
NBS2075
2076
NBS2080
NBS2081
2082
2086


```

DO 303 J=1,33
TOTAL(J)=DEB(J,33)/SUMEL
CHI(J)=(TOTAL(J)-Ezp(J))*2/Ezp(J)
CHISQ=CHISQ+CHI(J)
DEDNJ=DED(J,33)/SUMEL
ELM2(J)=FCVAL(1,J)
FACTOR=(TOTAL(J)-DEDNJ)/ELM2(J)
IF(ITER.GT.2)GO TO 302
XIP3(J)=(Ezp(J)-DEDNJ)/FACTOR
GO TO 303
302 XIP3(J)=FCVAL(1,J)/FACTOR
303 CONTINUE
214 SUM=(XIP3(1)+XIP3(33))/32.0
DO 207 J=2,32
SUM=SUM+XIP3(J)/16.0
207 CONTINUE
DO 208 I=1,33
FCVAL(1,I)=(XIP3(I))/SUM
208 CONTINUE
DO 220 I=1,33
XIP3(I)=DEB(I,33)-DED(I,33)
SUMSEL=(XIP3(1)+XIP(33))/32
DO 221 I=2,32
221 SUMSEL=SUMSEL+XIP3(I)/16
222 XIP3(I)=(XIP3(I)/SUMSEL-ELM2(I))/ELM2(I)
PRINT 209,ITER
209 FORMAT(1H1,20H AFTER ITERATION NO.,I3,49H . NORMALISATION IS UNIT
1AREA FOR ELASTIC EVENTS.)
216 ITER1=ITER+1
PRINT 210,ITER,ITER,ITER,ITER1
210 FORMAT(1H0,7HCOUNTER,4X,6HCOSINE,11X,3HEXP,7X,10HOUTPUT NO.,I2,4X,
13HNO.,I2,5H CHI ,7X,3HNO.,I2,11X,3HNO.,I2,11X,4X,11HDELTA SIGMA,2X, NBS2112
47HCOUNTER/2X,6HNUMBER,5X,5SHANGLE,10X,5HINPUT,7X,11HALL ELASTIC,6X, 2113
37HSQUARED,8X,5HINPUT,11X,5HINPUT,9X,6H/SIGMA,2X,6HNUMBER) NBS2114
PRINT 111,(J,COUNT(J),Ezp(J),TOTAL(J),CHI(J),ELM2(J),FCVAL(1,J), NBS2115

```

NBS2087
2088
NBS2089
2090
NBS
2097
NBS
NBS
NBS
NBS
NBS
2098
2099
2100
2101
2102
2103
2104
2105
NBS
NBS
NBS
NBS
NBS
NBS
2106
2107
2108
2109
2110
2111
NBS2112
2113
NBS2114
NBS2115

```
IXIP3(J),J,J=1,33)
111 FORMAT(14,7F15.6,I4)
    PRINT 112,CHISO
112 FORMAT(1H0,8HCHISO = ,F10.6)
997 IF(ITER-NTERM)415,416,416
415 SAMPLE=SAMP(ITER)
    ITER=ITER+1
    GO TO 999
416 ITER=10
    GO TO 998
999 CALL DATIN
998 RETURN
    END
```

```
NBS2116
2117
2118
2119
2120
2121
2122
2123
2124
2125
2126
2129
2130
```


A Brief Description of the Components of MAGGIE

1. MAGGIE (Main Program). This program calls the various subroutines required for the analysis, retrieves neutrons from disc storage when required, outputs track parameters if desired, and records the various fates of tracked neutrons.
2. RANDOM (and RDM). This subroutine contains the random number generator, RDM, as a separate entry. RANDOM is called at the beginning of program MAGGIE to enter the starting value for RDM.
3. DATIN serves to set all of the output arrays to zero, and calls ABSYND.
4. ABSYND reads the required nuclear data from the MOULD tape and puts it in encoded sequential storage in the array DATA-IDATA for use during the Monte Carlo tracking.
5. INPUT reads and processes samples and experimental angular distribution data from card input and calculates most of the flux attenuation factor.
6. CRNEU creates random incident neutrons at the entrance face of the sample, in accordance with the input source distribution.
7. TRACK tracks neutrons in the sample, specifying coordinates of collision or escape.
8. EGMV. This subroutine computes the mean free path, velocity, and lethargy group number.
9. FONEC calculates coordinates at the end of a track from initial position, direction cosines, and track length.
10. CR. This subroutine, using random sampling of the information stored by ABSYND from the MOULD data tape, determines all of the parameters of a collision.
11. CALC determines some constants used in subroutine NAPAN.
12. NAPAN. This subroutine calculates and scores the probability of detection at each detector for each collision in the sample.
13. FPATH calculates the path length in the direction of each detector for each collision.
14. CREN is an abbreviated version of CR used by NAPAN that determines only the neutron energy.

15. TWIST chooses new direction cosines after a collision.
16. WRIBT stores any secondary neutrons produced, for recall at the end of the current tracking.
17. OUTPUT, OUTB3, OUTB4 and APLOT3. These subroutines print the results of the calculation.
18. AUTO performs the calculations required for iterative correction of the elastic angular distribution, prints the current output, and calls DATIN to begin the next iteration.
19. END, and its entries EXIT and EEXIT designate normal vs. error exit conditions.
20. TAPLAB returns the tape logical unit label.
21. SRFORT is a subroutine for skipping tape records.

APPENDIX C

Input requirements for MAGGIE-NBS

	<u>FORMAT</u>	<u>VARIABLE</u>	<u>COMMENTS</u>
1.	012	OCT	Octal starting value for RDM.
2.	12A6	HEAD(I)	Arbitrary heading. Column one should be a 1.
3.	I10	NMATS	The number of materials in the sample. This is always = 1.
4.	I10, E10.4	IDATA(MAT5) DATA(MAT4)	Number of different nuclides in the material. Density of the material.
5.	I10, E10.4	IDATA(K2) DATA(K4)	Nuclide reference number (i.e. position on the MOULD tape: see output of MOULD for this). Proportion of this nuclide in the material. This card is repeated for each nuclide in the material.
6.	2I10	IFB4 NSENSE	Positive for B4 output. Positive for track print.
7.	3F10.4	HITE FRAD ANGLE	Length of the samples, in cm. Entrance radius of same, in cm. Half-angle of same, in radians.
8.	2I10	JOBFIN ITMAG	Number of independant samples (= 1 if ITMAG > 0). Positive for automatic iteration.
9.	F10.4	SAMPLE	Number of neutrons in each sample.
10.	F10.4	STRTE	Starting energy in MeV. If the starting energy is set = 0 a fission spectrum is assumed.
11.	F10.4	DIST	Distance from the center of the sample to the source (negative), in cm.
12.	I10	NSTDP	No. of start distribution points to be read in (≤ 25).
13.	6F10.4	ANGS(I)	Three pairs to a card, the angle and start distribution for that angle. The start distribution need not be normalized, and this card is repeated until NSTDP pairs are read in.

- | | | | |
|---|--------|--------------------------------------|--|
| 14. | I10 | NLWREV | The number of Law Reference numbers for inelastic scatter, i.e. the number of Law Reference numbers to different angular distributions on the MOULD output (P.C.N.'s* 4-15 only) (NLWREV ≤ 50). |
| 15. | 2I10 | LAWREV(I)
LWNO(I) | One pair per card, the above Law Reference numbers and the monotonically increasing "Monte Carlo" law numbers allocated. Any inelastic law not given here will be printed in the B3 section results under Law zero. |
| NOTE: The following cards are not required if IFB4 ≤ 0. | | | |
| 16. | 4I10 | IANALA
IANALB
IANALC
IANALD | Four markers, for the tables: A) Complete multiple scatter analysis. B) Elastic events only. C) Inelastic events only. D) Multiple elastic events only. A negative marker suppresses the table. A positive marker = N produces, in addition to the table, a graph for the Nth counter. |
| 17. | I10 | NANAL | Number of actions (index I below) (i.e. different neutron processes, such as elastic, inelastic from the 1st excited state, etc.) to be processed by the sub-routine NAPAN. |
| 18. | 6F10.4 | NUCL(I)
NAC1(I) | The nuclide reference number (as in card 5), and the P.C.N.* for each action. There are NANAL such pairs. |
| 19. | 6F10.4 | FCCOS(J) | Thirty-three values of cosine, including -1 and +1, monotonically increasing. |
| 20. | 6F10.4 | FCVAL(I,J) | Thirty-three values of the angular distribution (index J) corresponding to the above values of cosines. Repeat for each action (index I) in the same order as in 18. |

NOTE: The following three items are repeated, as a group, for each action.

- | | | | |
|-----|-----|---------|--|
| 21. | I10 | IMAX(I) | Number (index L) of supplementary ranges for this action (even if zero). |
|-----|-----|---------|--|

*The "P.C.N.'s (Particular Classification Numbers) are listed starting on p.8 of AWRE report no. O 70/63, "The Aldermaston Nuclear Data Library as at May, 1963", K. Parker.

NOTE: The following two items are repeated for each supplementary range. If, for any action, there is no supplementary range these items are omitted.

22.	F10.4	ENVAL(I, L)	Upper limit of the range in MeV. These must be listed in order of decreasing value.
23.	6F10.4	SUPVAL(L, J)	Thirty-three values (index J) of the distribution (range L) at the cosine values FCCOS(J).
24.	I10	NOCOU	Number of counters (≤ 33).
25.	6F10.4	COUNT(I)	Cosines of counter locations.
26.	I10	NENSP	Number of output energy points (≤ 33). The output will be classified into the bins formed by these points.
27.	6F10.6	ENSP(I)	Values of output energy points in MeV.
28.	2I10	NTERM IFCONT	Number of iterations (≤ 10). If this is a continuation of a previous run, for further iterations this should be non-zero (see Section II.c.6.f).
29.	6F10.4	SAMP(I)	Number of neutrons for each iteration, <u>except</u> the first.

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