

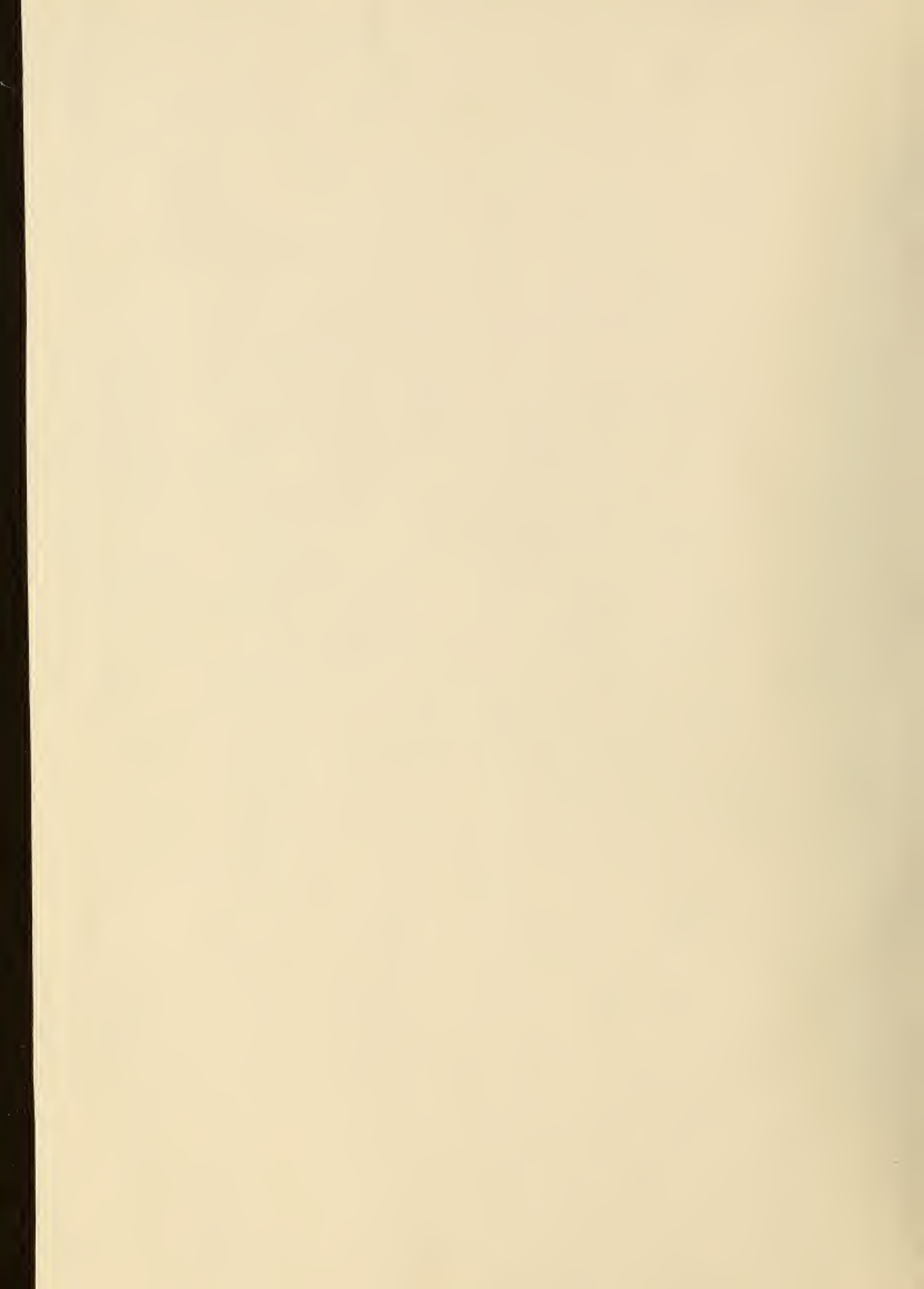
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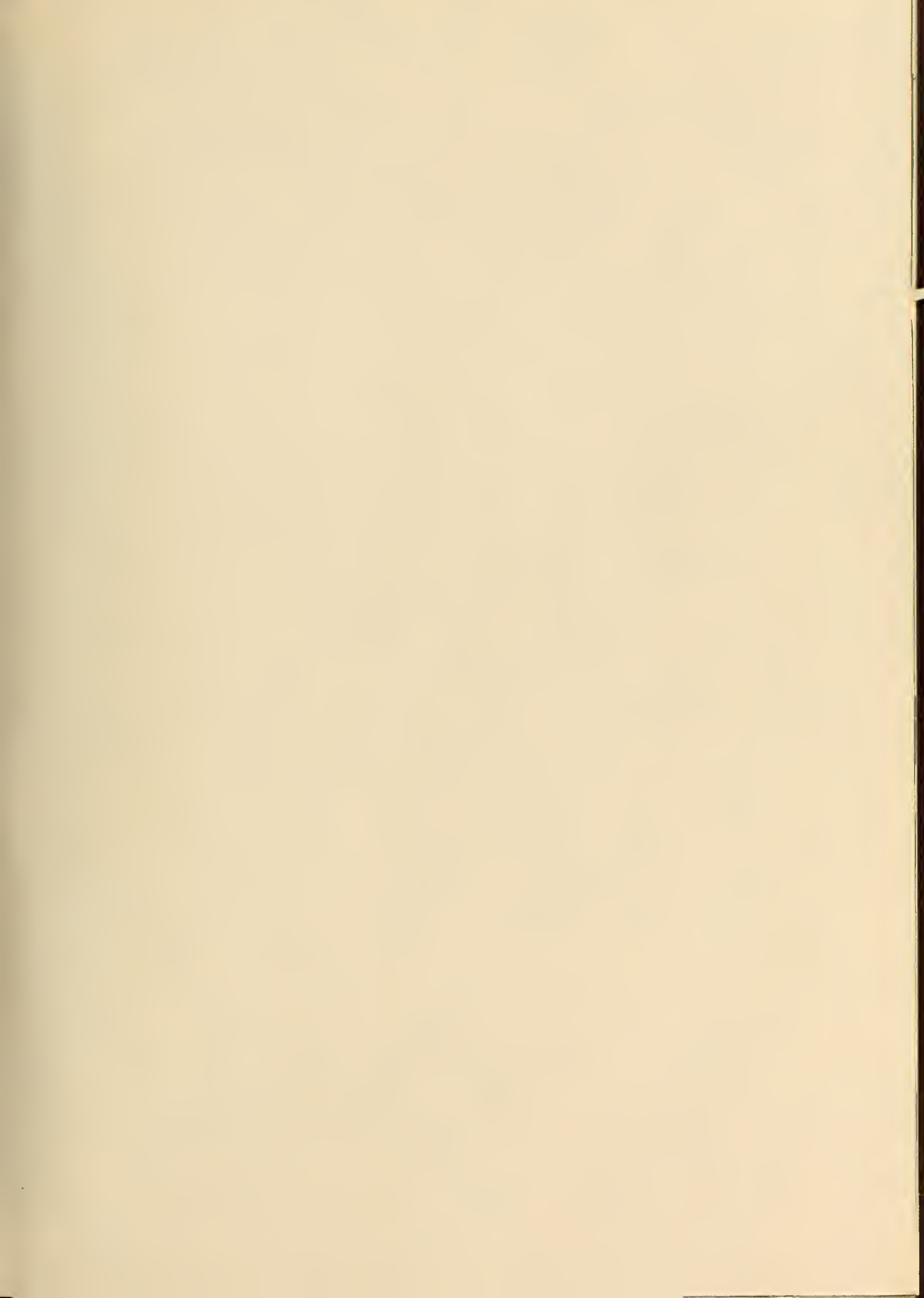


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COMPUTATION OF THE PERMEABILITY AND PERMITTIVITY OF A RELATIVELY SMALL RING SAMPLE IN A TOROIDAL COIL

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U. S. DEPARTMENT OF COMMERCE
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Computation of the Permeability and
Permittivity of a Relatively
Small Ring Sample in a
Toroidal Coil

Eric G. Johnson, Jr.

The derivation and the FORTTRAN machine program are presented for a formula which gives the impedance of a partially filled toroid coil. This formula shows the relationship of the complex permeability and permittivity of the ring material, the dimensions of the coil and sample, and the frequency of the applied electromagnetic field. Limitations of the formula due to proximity of the ring and coil, range of permeability and permittivity of the ring, and uniformity of the coil winding are considered. The cost of the program in terms of computer time is also considered. The principal value of the formula is to accurately determine the permeability and permittivity of ferrites.

Key Words: Complex permeability and permittivity, FORTRAN, impedance, toroid coil.

1. Introduction

In this technical note we develop the theory and the associated computer program for the impedance, Z , of a toroid sample with rectangular cross section imbedded in a toroid shaped frame of similar cross section around which a wire is wound. In order that the problem be amenable to theoretical solution it is necessary to require that the wire coil and the ferrite sample both have a rectangular cross section. The expected range of validity of this theory can be estimated by the requirement that $2\pi \frac{Df}{c} \mu \epsilon < 1$. Here D is the largest length in meters of the toroid cross section, the f is the frequency in hertz, and the ϵ , μ are the relative permittivity and permeability, respectively, of the sample. The c is the velocity of light which is in M. K. S. units 3×10^8 m/sec.

2. Procedure

Assume a ferrite sample with the following boundary conditions:

1. The experimental circumstances are such that the surface of a rectangular cross-sectioned toroid can be assumed to be a spacially uniform current sheet.

2. A second rectangular cross-sectioned toroid consisting of the ferrite material inside the current sheet toroid is to be constructed so that its size is electrically small compared with the current sheet toroid. This means that the ferrite sample is small enough that non-uniformity of a uniform current sheet will produce no essential difference in the determination of the permeability and permittivity of the ferrite sample.

3. Two complex frequency dependent parameters -- permeability μ and permittivity ϵ are used to characterize the ferrite.

4. The excitation mechanism of the toroid system is such that no angular field dependence exists. This means that the basic electromagnetic field equations will be reduced to a single equation. The basic equations that are used are

$$\begin{aligned} \nabla \times \underline{\underline{E}} &= -i \omega \mu \underline{\underline{H}}, \\ \text{and} \quad \nabla \times \underline{\underline{H}} &= i \omega \epsilon \underline{\underline{E}}. \end{aligned} \tag{1}$$

(natural units are used, namely $c = 1$)

5. The coordinates of interest are given in Figure 1.

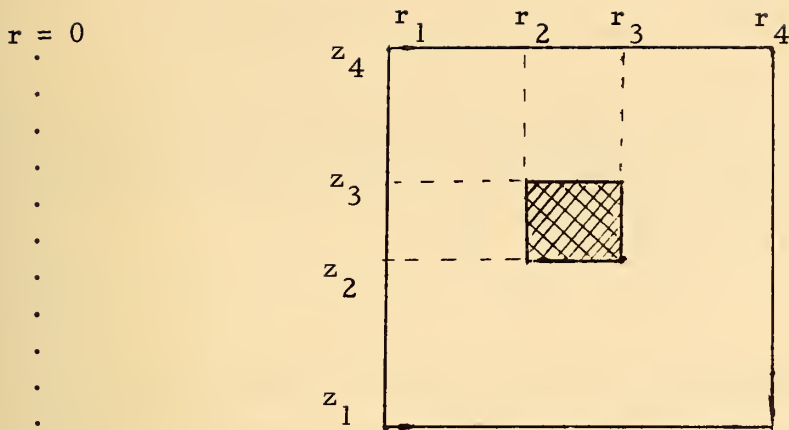


Figure 1

The r_1 , r_2 , r_3 , and r_4 are the radial coordinates. The z_1 , z_2 , z_3 , and z_4 are the "z" axis coordinates. The crosshatched area characterizes the ferrite sample.

By defining the function

$$P(r, z) = [\theta(r - r_2) - \theta(r - r_3)] [\theta(z - z_2) - \theta(z - z_3)],$$

$$\begin{aligned} \text{with } \theta(x) &= 1 && \text{if } x \geq 0 \\ \text{and } \theta(x) &= 0 && \text{if } x < 0 \end{aligned} ,$$

the permeability and permittivity may be written as $\mu = 1 + (\mu_M - 1)P(r, z)$, and $\epsilon = 1 + (\epsilon_M - 1)P(r, z)$. The μ_0 , ϵ_0 which are the vacuum permeability and permittivity, respectively, are set equal to one. The μ_M , ϵ_M are the ferrite's permeability and permittivity, respectively. The μ and ϵ give the permeability and permittivity for the entire cross section of Figure 1.

Upon reducing (1) with assumptions 1 to 5 and letting $f = rH_\phi$, the following differential equation results:

$$\frac{r}{\mu} \partial_r \left(\frac{1}{r\epsilon} \partial_r f \right) + \frac{1}{\mu} \partial_z \left(\frac{1}{\epsilon} \partial_z f \right) + \omega^2 f = 0. \quad (2)$$

Here $\left(\partial_r \equiv \frac{\partial}{\partial r} \right)$, $\left(\partial_z \equiv \frac{\partial}{\partial z} \right)$.

In order to get rapidly convergent results in the machine calculation, we extract the quasi-static part from this equation, and to make the function continuous to its first derivative, we write

$$f = I \left(1 + \epsilon g(r, z) \right). \quad (3)$$

This transformation assumes that the discontinuity of the material surface tends under the condition of the experimental system to form a surface current,

$$V = i \omega \int \frac{\mu f}{r} dr dz = i \omega I \left[\int \frac{\mu}{r} (1 + \epsilon g) dr dz \right], \quad (4)$$

and that the total current is given by $I_{\text{total}} = 2\pi H_{\varphi} r$; where $r = r_1$ and H_{φ} is the magnetic field intensity along the azimuthal direction. Since $I_{\text{total}} = 2\pi f(r_1, z) = Nj$, where j is the current in each wire loop, we deduce that $I = \frac{Nj}{2\pi}$. Therefore, the measured impedance may be written as

$$Z = \frac{NV}{j} = i \frac{\omega N^2}{2\pi} \left[\ln \left(\frac{r_4}{r_1} \right) (z_4 - z_1) + (\mu_M - 1) \ln \left(\frac{r_3}{r_2} \right) (z_3 - z_2) + \int dr dz \frac{\mu \epsilon}{r} g(r, z) \right]. \quad (5)$$

Because natural units are used, Z is dimensionless. The first and second terms represent the quasi-static approximation. The last term gives the corrections. To evaluate Z we need g . The differential equation for g is obtained by substituting (3) into (2). The following equation is obtained:

$$r \partial_r \frac{1}{r} \partial_r g + \partial_z^2 g + \omega^2 \mu \epsilon g = -\mu \omega^2. \quad (6)$$

Assume the following form for the solution:

$$g = \sum_n R_n(r) Z_n(z) \quad (7)$$

in which $g = 0$ on the surface of the current toroid.

3. Analysis

Once the above steps have been taken, the rest of the analysis uses in the usual manner normal modes and perturbation theory.¹ Because of the cylindrical symmetry, the best orthogonal set to use is Bessel functions for the radial dependence. This gives integrable results.

The Bessel functions used are implied by the equation

$$\left(r \frac{\partial}{\partial r} \frac{1}{r} \frac{\partial}{\partial r} + k_n^2 \right) R_n = 0, \quad (8a)$$

which can be obtained from the homogeneous version of (6), with (7) substituted for g and with the boundary conditions $R_n(r_1) = R_n(r_4) = 0$. The dual function to R_n is given by $\bar{R}_n = \frac{1}{r} \bar{R}_n$, where \bar{R}_n satisfies

$$\bar{R}_n \left[r \overleftarrow{\frac{\partial}{\partial r}} \frac{1}{r} \overleftarrow{\frac{\partial}{\partial r}} + k_n^2 \right] = 0, \quad (8b)$$

where the $\overleftarrow{\frac{\partial}{\partial r}}$ denotes differentiation to the left. The orthogonality condition is

$$\int_{r_1}^{r_4} \bar{R}_n R_{n'} dr = \delta_{nn'}, \quad (8c)$$

where $\delta_{nn'} = 1$ $n = n'$
 $= 0$ $n \neq n'$.

$Z_n(z)$ obeys the following differential equation:

$$\left(\frac{\partial^2}{\partial z^2} + \omega^2 - k_n^2 \right) Z_n + \sum A_{nn'} P(z) Z_{n'} = -B_n - D_n P(z). \quad (9)$$

¹P. M. Morse and H. Feshbach, Methods of Theoretical Physics (McGraw Hill, New York, 1953), Vol. II, Chapter 9, pp. 1001-1025.

This is derived by taking (6), multiplying by \bar{R}_n , and integrating from r_1 to r_4 . In (9),

$$A_{nn'} \equiv \int_{r_2}^{r_3} dr \bar{R}_n(r) R_{n'}(r) \omega^2 (\mu_M \epsilon_M - 1),$$

$$B_n \equiv \omega^2 \int_{r_1}^{r_4} \bar{R}_n(r) dr,$$

$$D_n \equiv \omega^2 (\mu_M - 1) \int_{r_2}^{r_3} \bar{R}_n(r) dr,$$

and

$$P(z) \equiv \theta(z - z_2) - \theta(z - z_3).$$

Define the variables $\lambda_n^2 = -\omega^2 + k_n^2$, and $\bar{\lambda}_n^2 = \lambda_n^2 - A_{nn'}$, or

$\lambda_n^2(z) = \lambda_n^2 + (\bar{\lambda}_n^2 - \lambda_n^2)P(z)$, and use the Green's function defined by

$$G_n \left[\frac{\partial}{\partial z}^2 - \lambda_n^2(z) \right] = -\delta(z - \bar{z})$$

($\delta(z - \bar{z})$ is the Dirac delta function)

to solve the differential (9), the resulting integral equation is

$$Z_n(\bar{z}) = \sum_{n' \neq n} \int_{z_2}^{z_3} A_{nn'} G_n(\bar{z}, z) Z_{n'}(z) dz + \int_{z_1}^{z_4} G_n(\bar{z}, z) B_n dz$$

$$+ D_n \int_{z_2}^{z_3} G_n(\bar{z}, z) dz. \tag{10}$$

By considering only first order mode correlation, $Z_n(\bar{z})$ may be approximately written as

$$Z_n(\bar{z}) \cong \left(\int_{z_1}^{z_4} G_n(\bar{z}, z) dz B_n + D_n \int_{z_2}^{z_3} G_n(\bar{z}, z) dz \right) \quad (11)$$

$$+ \sum_{n' \neq n} \int_{z_2}^{z_3} A_{nn'} G_n(\bar{z}, \bar{z}) d\bar{z} \left[\int_{z_1}^{z_4} G_{n'}(\bar{z}, z) B_{n'} dz + \int_{z_2}^{z_3} D_{n'} G_{n'}(\bar{z}, z) dz \right].$$

This approximation implies that the sample is electrically small. G_n is $G_n(z, \bar{z}) = A_n C_n(z_>) S_n(z_<)$, where $z_>$ is the larger of z, \bar{z} pair and $z_<$ is the smaller of z, \bar{z} pair. A_n is the normalization. It is such that (12) is true:

$$A_n C_n S_n - C_n S_n = -1, \quad (12)$$

where $z = \bar{z}$.

The functions S_n, C_n are defined as follows:

$$\left. \begin{aligned} S_n &= \exp[\lambda_n(z - z_1)] - \exp[-\lambda_n(z - z_1)] \\ C_n &= C_n^1 \exp[\lambda_n(z - z_2)] + C_n^2 \exp[-\lambda_n(z - z_2)] \end{aligned} \right\} z_2 \geq z \geq z_1$$

$$\left. \begin{aligned} S_n &= \bar{S}_n^1 \exp[\bar{\lambda}_n(z - z_2)] + \bar{S}_n^2 \exp[-\bar{\lambda}_n(z - z_2)] \\ C_n &= \bar{C}_n^1 \exp[\bar{\lambda}_n(z - z_3)] + \bar{C}_n^2 \exp[-\bar{\lambda}_n(z - z_3)] \end{aligned} \right\} z_3 \geq z \geq z_2$$

$$\begin{aligned} S_n &= S_n^1 \exp[\lambda_n(z - z_3)] + S_n^2 \exp[-\lambda_n(z - z_3)] \\ C_n &= \exp[\lambda_n(z - z_4)] - \exp[-\lambda_n(z - z_4)]. \end{aligned}$$

Here,

$$\lambda_n = + \left[-\omega^2 + k_n^2 \right]^{\frac{1}{2}}$$

and

$$\bar{\lambda}_n = + \left[\lambda_n^2 - A_{nn} \right]^{\frac{1}{2}}.$$

(13)

The coefficients C_n^1 , C_n^2 , \bar{C}_n^1 , \bar{C}_n^2 and S_n^1 , S_n^2 , \bar{S}_n^1 , \bar{S}_n^2 are defined with the aid of the function f_1 :

$$f_1(\lambda_n, \bar{\lambda}_n, 3, 4) \equiv \frac{1}{2} \left[\left(\frac{\lambda_n}{\bar{\lambda}_n} + 1 \right) a_{34}^n + \left(\frac{\lambda_n}{\bar{\lambda}_n} - 1 \right) a_{43}^n \right].$$

Thus

$$\bar{C}_n^1 \equiv f_1(\lambda_n, \bar{\lambda}_n, 3, 4)$$

$$\bar{C}_n^2 \equiv f_1(\lambda_n, -\bar{\lambda}_n, 3, 4)$$

$$\bar{S}_n^1 \equiv f_1(\lambda_n, \bar{\lambda}_n, 2, 1)$$

(14)

and

$$\bar{S}_n^2 \equiv f_1(\lambda_n, -\bar{\lambda}_n, 2, 1).$$

By defining a second function

$$f_2(\lambda_n, \bar{\lambda}_n, 2, 3, \bar{C}_n^1, \bar{C}_n^2) \equiv \frac{1}{2} \left[\bar{C}_n^1 \left(1 + \frac{\bar{\lambda}_n}{\lambda_n} \right) a_{23}^n + \bar{C}_n^2 \left(1 - \frac{\bar{\lambda}_n}{\lambda_n} \right) a_{32}^n \right],$$

the C_n^1 , etc., are:

$$C_n^1 = f_2(\lambda_n, \bar{\lambda}_n, 2, 3, \bar{C}_n^1, \bar{C}_n^2)$$

$$C_n^2 = f_2(-\lambda_n, \bar{\lambda}_n, 2, 3, \bar{C}_n^1, \bar{C}_n^2)$$

$$S_n^1 = f_2(\lambda_n, \bar{\lambda}_n, 3, 2, \bar{S}_n^1, \bar{S}_n^2)$$

$$S_n^2 = f_2(-\lambda_n, \bar{\lambda}_n, 3, 2, \bar{S}_n^1, \bar{S}_n^2).$$

(15)

The a_{34}^n , etc. are defined as

$$a_{34}^n = \exp[\lambda_n(z_3 - z_4)] = (a_{43}^n)^{-1}$$

$$a_{23}^n = \exp[\bar{\lambda}_n(z_2 - z_3)] = (a_{32}^n)^{-1}$$

and

$$a_{12}^n = \exp[\lambda_n(z_1 - z_2)] = (a_{21}^n)^{-1}.$$

Using the above definitions, the normalization A_n is

$$A_n = 1/2 \lambda_n \left[C_n^1 a_{12}^n + C_n^2 (a_{12}^n)^{-1} \right]$$

and the impedance reduces to

$$Z = i \frac{\omega N^2}{2\pi} \left[(z_4 - z_1) \ln(r_4/r_1) + (\mu_M - 1)(z_3 - z_2) \ln(r_3/r_2) + \sum_{n=1} \left\{ \begin{array}{l} [T_n(r_4) - T_n(r_1)] (V_n + W_n) \\ (\mu_M - 1) [T_n(r_3) - T_n(r_2)] W_n \end{array} \right\} N_n \right]. \quad (16)$$

Here we have further defined

$$V_n \equiv \int_{z_1}^{z_4} Z_n dz,$$

$$W_n \equiv \int_{z_2}^{z_3} Z_n dz,$$

and have noted that the T_n is a solution to the equation

$$\left[\frac{d}{dr}{}^2 + \frac{1}{r} \frac{d}{dr} + k_n^2 \right] T_n = 0. \quad \text{The } R_n \text{ is related to } T_n \text{ by } R_n = r \frac{dT_n}{dr} N_n. \quad \text{The}$$

k_n is fixed by the condition that at $r = r_1$ and r_4 , $\frac{dT_n}{dr} = 0$. The N_n is

the normalization factor needed to have (8c) true. We can write T_n in terms of Bessel functions of the zero and first order as

$$T_n = \left[J_0(k_n r) Y_1(k_n r_1) - J_1(k_n r_1) Y_0(k_n r) \right] \quad (17a)$$

The coefficients $A_{nn'}$, B_n , D_n , and N_n are given in terms of the Bessel functions of zero and first order as

$$B_n = N_n \omega^2 \left[T_n(r_4) - T_n(r_1) \right],$$

$$D_n = N_n \omega^2 (\mu_n - 1) \left[T_n(r_3) - T_n(r_2) \right],$$

and

$$A_{nn'} = \omega^2 \frac{(\mu_M - 1)}{(k_{n'}^2 - k_n^2)} N_n N_{n'} \left[-rk_n^2 T_n \frac{d}{dr} T_{n'} + rk_{n'}^2 T_{n'} \frac{d}{dr} T_n \right] \Bigg|_{r_2}^{r_3} \quad (17b)$$

where $n \neq n'$, when

$n = n'$, we have

$$A_{nn} = \omega^2 (\mu_M - 1) N_n^2 \left[r \left\{ T_n \frac{\partial}{\partial r} T_n + k_n^2 \left(\frac{\partial}{\partial k} T_n \right) \frac{\partial}{\partial r} T_n - T_n \frac{\partial}{\partial r} \left(\frac{\partial}{\partial k} T_n \right) \right\} \right] \Bigg|_{r_2}^{r_3}$$

The N_n^2 is given by

$$N_n^2 = 1 / \left[k_n^2 \left(r_1 T_n(r_1) \frac{\partial}{\partial r} \frac{\partial}{\partial k} T_n(r_1) - r_4 T_n(r_4) \frac{\partial}{\partial r} \frac{\partial}{\partial k} T_n(r_4) \right) \right].$$

If we use the recurrence relationships

$$\frac{d}{dr} J_0(k_n r) = -k_n J_1(k_n r),$$

$$\frac{d}{dr} Y_0(k_n r) = -k_n Y_1(k_n r),$$

$$\frac{d}{dr} J_1(k_n r) = k_n J_0(k_n r) - \frac{1}{r} J_1(k_n r),$$

and

$$\frac{d}{dr} Y_1(k_n r) = k_n Y_0(k_n r) - \frac{1}{r} Y_1(k_n r),$$

we find that N_n^2 and A_{nn} become

(18)

$$N_n^2 = 1 / \left[k_n^2 \left(T_n(r) r \right) \left(r T_n + r_1 \left(Y_0(r_1) J_1(r) - Y_1(r) J_0(r_1) \right) \right) \right] \left[\frac{r}{r_1} \right]^4,$$

$$A_{nn} = \omega^2 (\mu_M - 1) N_n^2 \left[r \left\{ T_n \frac{d}{dr} T_n + \frac{d}{dr} T_n \left[r d_r T_n k_n r_1 \left(Y_0(k_n r_1) J_0(k_n r) - Y_0(k_n r) J_0(k_n r_1) \right) \right] + \frac{T_n}{2} k_n^2 \left[r T_n + r_1 \left(Y_0(r_1) J_1(r) - Y_1(r) J_0(r_1) \right) \right] \right\} \right] \left[\frac{r}{r_2} \right]^3 \quad (19)$$

In order to evaluate V_n and W_n , we note the definitions

$$B'_n \equiv T_n(r_4) - T_n(r_1),$$

(20)

and

$$D'_n \equiv T_n(r_3) - T_n(r_2),$$

and explicitly insert into the impedance, the G_n for V_n and W_n . The impedance finally becomes

$$Z = \frac{N_n^2 i \omega}{2\pi} \left[(z_4 - z_1) \ln(r_4/r_1) + (\mu_M - 1) (z_3 - z_2) \ln(r_3/r_2) + \right. \quad (21) \\ \left. + \sum_{n=1}^{\infty} A_n N_n^2 B'_n \left\{ B_n (R_n^1 + R_n^3) + D_n (R_n^2 + R_n^4) \right\} + \right. \\ \left. + \sum_{n=1}^{\infty} A_n N_n^2 B'_n \sum_{n' > n} N_{n'}^2 A_{nn'} A_{n'} \left\{ B_{n'} (R_{nn'}^5 + R_{nn'}^7 + R_{n'n}^5 + R_{n'n}^7) \right. \right. \\ \left. \left. + D_{n'} (R_{nn'}^6 + R_{nn'}^8 + R_{n'n}^6 + R_{n'n}^8) \right\} \right]$$

$$\begin{aligned}
& + (\mu_M - 1) \sum_{n=1}^{\infty} A_n D'_n N_n^2 (D_n R_n^4 + B_n R_n^3) \\
& + (\mu_M - 1) \sum_{n=1}^{\infty} D'_n A_n N_n^2 \sum_{n' > n} A_{nn'} A_{n'n} N_{n'}^2 \left\{ B_{n'} (R_{nn'}^7 + R_{n'n}^7) \right. \\
& \qquad \qquad \qquad \left. + D_{n'} (R_{nn'}^8 + R_{n'n}^8) \right\}
\end{aligned}$$

Here we have explicitly extracted the normalization factors N_n , $N_{n'}$, from B_n , D_n , and $A_{nn'}$. The R_n^1 , R_n^2 , etc., are defined as follows:

$$\begin{aligned}
R_n^1 &= \int_{z_1}^{z_2} C_n(\bar{z}) d\bar{z} \int_{z_1}^{\bar{z}} S_n(z) dz + \int_{z_1}^{z_2} S_n(\bar{z}) d\bar{z} \int_{\bar{z}}^{z_2} C_n dz \\
&+ \int_{z_3}^{z_4} C_n d\bar{z} \int_{z_3}^{\bar{z}} S_n dz + \int_{z_3}^{z_4} S_n d\bar{z} \int_{\bar{z}}^{z_4} C_n dz + \left(\int_{z_1}^{z_2} S_n d\bar{z} \right) \\
&\left(\int_{z_3}^{z_4} C_n dz + \int_{z_2}^{z_3} C_n dz \right) + \left(\int_{z_3}^{z_4} C_n d\bar{z} \right) \left(\int_{z_1}^{z_2} S_n dz + \int_{z_2}^{z_3} S_n dz \right)
\end{aligned}$$

$$R_n^2 = \left(\int_{z_1}^{z_2} S_n(\bar{z}) d\bar{z} \right) \left(\int_{z_2}^{z_3} C_n dz \right) + \left(\int_{z_3}^{z_4} C_n d\bar{z} \right) \left(\int_{z_2}^{z_3} S_n dz \right)$$

$$\begin{aligned}
R_n^3 &= \int_{z_2}^{z_3} C_n d\bar{z} \int_{z_2}^{\bar{z}} S_n dz + \int_{z_2}^{z_3} S_n d\bar{z} \int_{\bar{z}}^{z_3} C_n dz + \left(\int_{z_2}^{z_3} C_n d\bar{z} \right) \left(\int_{z_1}^{z_2} S_n dz \right) \\
&+ \left(\int_{z_2}^{z_3} S_n d\bar{z} \right) \left(\int_{z_3}^{z_4} C_n dz \right)
\end{aligned}$$

$$R_n^4 = \int_{z_2}^{z_3} C_n d\bar{z} \int_{z_2}^{\bar{z}} S_n dz + \int_{z_2}^{z_3} S_n d\bar{z} \int_{\bar{z}}^{z_3} C_n dz$$

$$R_{nn'}^5 = \left(\int_{z_1}^{z_2} S_n dz \right) \left[\int_{z_2}^{z_3} C_n(\bar{z}) d\bar{z} \left\{ C_{n'} \int_{z_2}^z S_{n'} dz + S_{n'} \int_z^{z_3} C_{n'} dz + C_n \left(\int_{z_1}^{z_2} S_{n'} dz \right) \right. \right. \\ \left. \left. + S_{n'} \left(\int_{z_3}^{z_4} C_{n'} dz \right) \right\} \right] + \left(\int_{z_3}^{z_4} C_n dz \right) \left[\int_{z_2}^{z_3} S_n(\bar{z}) d\bar{z} \left\{ C_{n'} \int_{z_2}^{\bar{z}} S_{n'} dz \right. \right. \\ \left. \left. + S_{n'} \int_{\bar{z}}^{z_3} C_{n'} dz + C_{n'} \left(\int_{z_1}^{z_2} S_{n'} dz \right) + S_{n'} \left(\int_{z_3}^{z_4} C_{n'} dz \right) \right\} \right]$$

$$R_{nn'}^6 = \left(\int_{z_1}^{z_2} S_n dz \right) \left[\int_{z_2}^{z_3} C_n d\bar{z} \left(C_{n'} \int_{z_2}^{\bar{z}} S_{n'} dz + S_{n'} \int_{\bar{z}}^{z_3} C_{n'} dz \right) \right] \\ + \left(\int_{z_3}^{z_4} C_n dz \right) \left[\int_{z_2}^{z_3} S_n d\bar{z} \left(C_{n'} \int_{z_2}^{\bar{z}} S_{n'} dz + \int_{\bar{z}}^{z_3} C_{n'} dz \right) \right]$$

$$R_{nn'}^7 = \int_{z_2}^{z_3} C_n(\bar{z}) d\bar{z} \int_{z_2}^{\bar{z}} S_n(z) dz \left[C_{n'} \int_{z_2}^z S_{n'} d\bar{z} + S_{n'} \int_z^{z_3} C_{n'} d\bar{z} + \right. \\ \left. + C_{n'} \left(\int_{z_1}^{z_2} S_{n'} d\bar{z} \right) + S_{n'} \left(\int_{z_3}^{z_4} C_{n'} d\bar{z} \right) \right] + \int_{z_2}^{z_3} S_n(\bar{z}) dz \int_{\bar{z}}^{z_3} C_n(z) dz$$

$$\cdot \left[C_{n'} \int_{z_2}^z S_{n'} d\bar{z} + S_{n'} \int_z^{z_3} C_{n'} d\bar{z} + C_{n'} \left(\int_{z_1}^{z_2} S_{n'} d\bar{z} \right) + S_{n'} \left(\int_{z_3}^{z_4} C_{n'} d\bar{z} \right) \right]$$

$$R_{nn'}^8 = \int_{z_2}^{z_3} C_n(\bar{z}) d\bar{z} \int_{z_2}^{\bar{z}} S_n(z) dz \left[C_{n'} \int_{z_2}^z S_{n'} d\bar{z} + S_{n'} \int_z^{z_3} C_{n'} d\bar{z} \right]$$

$$+ \int_{z_2}^{z_3} S_n(\bar{z}) d\bar{z} \int_{\bar{z}}^{z_3} C_n(z) dz \left[C_{n'} \int_{z_2}^z S_{n'} d\bar{z} + S_{n'} \int_z^{z_3} C_{n'} d\bar{z} \right]$$

By making use of common factors in R_n^1 , etc., the above may be reduced to

$$R_n^1 = R_n^2 + 2F_{n12} G_{n34} + L_n^1 \quad (22)$$

$$R_n^2 = F_{n12} G_{n23} + F_{n23} G_{n34} \quad (23)$$

$$R_n^3 = R_n^4 + R_n^2 \quad (24)$$

$$R_{nn'}^5 = R_{nn'}^6 + H_{nn'}^1 F_{n12} F_{n'12} + H_{nn'}^2 G_{n34} G_{n'34} + F_{n12} G_{n'34} H_{nn'}^3$$

$$+ G_{n34} F_{n'12} H_{nn'}^4 \quad (25)$$

$$R_{nn'}^7 = R_{nn'}^8 + F_{n'12} L_{nn'}^4 + G_{n'34} L_{nn'}^5 \quad (26)$$

$$R_{nn'}^6 = L_{nn'}^2 F_{n12} + L_{nn'}^3 G_{n34} \quad (27)$$

The remaining definitions are

$$L_n^1 = \int_{z_1}^{z_2} C_n d\bar{z} \int_{z_1}^{\bar{z}} S_n dz + \int_{z_1}^{z_2} S_n d\bar{z} \int_{\bar{z}}^{z_2} C_n dz + \int_{z_3}^{z_4} C_n d\bar{z} \int_{z_3}^{\bar{z}} S_n d\bar{z}$$

$$+ \int_{z_3}^{z_4} S_n d\bar{z} \int_{\bar{z}}^{z_4} C_n dz$$

$$H_{nn'}^1 = \int_z^{z_3} C_n C_{n'} dz$$

$$H_{nn'}^2 = \int_{z_2}^{z_3} S_n S_{n'} dz$$

$$H_{nn'}^3 = \int_{z_2}^{z_3} C_n S_{n'} dz$$

$$H_{nn'}^4 = \int_{z_2}^{z_3} S_n C_{n'} dz$$

$$L_{nn'}^2 = \int_{z_2}^{z_3} \left[C_n d\bar{z} \left(C_{n'} \int_{z_2}^{\bar{z}} S_{n'} dz + S_{n'} \int_{\bar{z}}^{z_3} C_{n'} dz \right) \right]$$

$$L_{nn'}^3 = \int_{z_2}^{z_3} \left[S_n d\bar{z} \left(C_{n'} \int_{z_2}^{\bar{z}} S_{n'} dz + S_{n'} \int_{\bar{z}}^{z_3} C_{n'} dz \right) \right]$$

$$L_{nn'}^4 = \int_{z_2}^{z_3} C_n d\bar{z} \int_{z_2}^{\bar{z}} S_n C_{n'} dz + \int_{z_2}^{z_3} S_n d\bar{z} \int_{\bar{z}}^{z_3} C_n C_{n'} dz$$

$$L_{nn'}^5 = \int_{z_2}^{z_3} \left[C_n \int_{z_2}^{\bar{z}} S_n S_{n'} dz + S_n \int_{\bar{z}}^{z_3} C_n S_{n'} dz \right] d\bar{z}$$

$$F_{n12} = \int_{z_1}^{z_2} S_n dz$$

$$F_{n23} = \int_{z_2}^{z_3} S_n dz$$

$$G_{n23} = \int_{z_2}^{z_3} C_n dz ,$$

and

$$G_{n34} = \int_{z_3}^{z_4} C_n dz .$$

The results of integration of the above definitions are listed as follows.

The single integral expressions become

$$F_{n12} = \frac{1}{\lambda_n} \left[a_{n12} + (a_{n12})^{-1} - 2 \right] \quad (28)$$

$$F_{n23} = \frac{1}{\lambda_n} \sum_i \bar{S}_n^i (-1)^{i+1} \left[(a_{n23})^{(-1)^i} - 1 \right] \quad (29)$$

$$G_{n23} = \frac{1}{\lambda_n} \sum_i \bar{C}_n^i (-1)^{i+1} \left[1 - (a_{n23})^{(-1)^{i+1}} \right] \quad (30)$$

$$G_{n34} = \frac{1}{\lambda_n} \left[2 - a_{n34} - (a_{n34})^{-1} \right] . \quad (31)$$

If we define

$$\lambda_{ml}(i, j) \equiv \left[(-1)^{i+1} \bar{\lambda}_n + (-1)^{j+1} \bar{\lambda}_{n'} \right] \quad (32a)$$

$$P_1(i, j) \equiv \left[1 - (a_{n23})^{(-1)^{i+1}} (a_{n'23})^{(-1)^{j+1}} \right] / \lambda_{ml}(i, j) ,$$

and

$$P_2(i, j) \equiv \left[(a_{n'23})^{(-1)^j} - (a_{n23})^{(-1)^{i+1}} \right] / \lambda_{ml}(i, j) , \quad (32b)$$

then the remaining single integral becomes

$$\begin{aligned}
 H_{nn'}^1 &= \sum_{i,j} \bar{C}_n^i \bar{C}_{n'}^j P_1(i,j) \\
 H_{nn'}^2 &= \sum_{i,j} \bar{S}_n^i \bar{S}_{n'}^j P_1(i-1, j-1) \\
 H_{nn'}^3 &= \sum_{i,j} \bar{C}_n^i \bar{S}_{n'}^j P_2(i,j) = H_{n'n}^4 \\
 H_{nn'}^4 &= \sum_{i,j} \bar{S}_n^i \bar{C}_{n'}^j P_2(i-1, j-1) = H_{n'n}^3.
 \end{aligned} \tag{33}$$

The double integrations give

$$\begin{aligned}
 R_n^4 &= \sum_{i,j} \frac{\bar{C}_n^i \bar{S}_n^j}{\bar{\lambda}_n} \left[\frac{(-1)^{j+i}}{\bar{\lambda}_n} \left[-2 + (a_{n23})^{(-1)^{i+1}} + (a_{n23})^{(-1)^j} \right] \right. \\
 &\quad \left. + (z_2 - z_3) (a_{n23})^{(-1)^j} \left[(-1)^j - (-1)^i \right] \right]
 \end{aligned} \tag{34}$$

$$\begin{aligned}
 L_n^1 &= \frac{1}{\lambda_n} \sum_{i,j} \left[C_n^i (-1)^{j+1} \left\{ (-1)^{j+i} \left[-2 + (a_{n12})^{(-1)^{i+1}} + (a_{n12})^{(-1)^j} \right] / \lambda_n \right. \right. \\
 &\quad \left. \left. + (z_1 - z_2) (a_{n12})^{(-1)^j} \left[(-1)^j - (-1)^i \right] \right\} \right. \\
 &\quad \left. + S_n^j (-1)^{i+1} \left\{ (-1)^{j+i} \left[-2 + (a_{n34})^{(-1)^{i+1}} + (a_{n34})^{(-1)^j} \right] / \lambda_n \right. \right. \\
 &\quad \left. \left. + (z_3 - z_4) (C_{n34})^{(-1)^j} \left[(-1)^j - (-1)^i \right] \right\} \right]
 \end{aligned} \tag{35}$$

Upon further definition

$$P_4(i) \equiv \left[1.0 - (a_{n23})^{(-1)^{i+1}} \right] / \bar{\lambda}_n (-1)^{i+1} \tag{36}$$

and

$$P_7(i, j) \equiv P_4(i) (a_{n', 23})^{(-1)^j} ; \quad (37)$$

the two-fold integrations give

$$L_{nn'}^2 = \sum_{i j \ell} \frac{\bar{C}_n^\ell \bar{C}_{n'}^i \bar{S}_{n'}^j}{\bar{\lambda}_{n'}} \left[- \left\{ (-1)^j - (-1)^i \right\} P_7(\ell, j) + (-1)^j P_1(\ell, i) \right. \\ \left. + (-1)^{i+1} P_2(\ell, j) \right] \quad (38)$$

$$L_{nn'}^3 \equiv \sum_{i j \ell} \frac{\bar{S}_n^\ell \bar{C}_{n'}^i \bar{S}_{n'}^j}{\bar{\lambda}_{n'}} \left[- \left\{ (-1)^j - (-1)^i \right\} P_7(\ell-1, j) + (-1)^j \right. \\ \left. \cdot P_2(\ell-1, i+1) + (-1)^{i+1} P_1(\ell-1, j-1) \right]. \quad (39)$$

If we define

$$\lambda_{m2}(i, j, \ell) \equiv \bar{\lambda}_n \left[(-1)^{j+1} + (-1)^{i+1} \right] + \bar{\lambda}_{n'} (-1)^{\ell+1} , \quad (40)$$

$$\lambda_{m3}(\ell, i, j) \equiv \bar{\lambda}_{n'} \left[(-1)^{i+1} + (-1)^{j+1} \right] + \bar{\lambda}_n (-1)^{\ell+1} , \quad (41)$$

and

$$R_3(i, j, \ell) \equiv \left[(a_{n23})^{(-1)^j} - (a_{n23})^{(-1)^{i+1}} (a_{n', 23})^{(-1)^{\ell+1}} \right] / \lambda_{m2}(i, j, \ell), \quad (42)$$

the remaining two-fold integrations are

$$L_{nn'}^4 = \sum_{i j \ell} \bar{C}_n^i \bar{S}_n^j \bar{C}_{n'}^\ell \left[\left(P_3(i, j, \ell) - P_7(i, \ell+1) \right) / \lambda_{m1}(j, \ell) \right. \\ \left. + \left(-P_3(i, j, \ell) + P_4(j-1) \right) / \lambda_{m1}(i, \ell) \right] \quad (43)$$

$$L_{nn'}^5 = \sum_{i,j,\ell} \bar{C}_n^i \bar{S}_n^i \bar{S}_n^\ell \left[\left(P_3(j-1, i+1, \ell-1) - P_4(i) \right) / \lambda_{m1}(j, \ell) \right. \\ \left. + \left(-P_3(j-1, i+1, \ell-1) + P_7(j-1, \ell) \right) / \lambda_{m1}(i, \ell) \right] \quad (44)$$

The triple integration gives

$$R_{nn'}^8 = \sum_{i,j,\ell,m} \frac{\bar{C}_n^i \bar{C}_{n'}^j \bar{S}_n^\ell \bar{S}_{n'}^m}{\lambda_{n'}} \\ \cdot \left[(-1)^m \left[\left(P_5(i, j, \ell, m) - P_7(i, j+1) \right) / \lambda_{m3}(\ell, j, m) \right. \right. \\ \left. \left. \left(-P_3(i, \ell, j) + P_7(i, j+1) \right) / \lambda_{m1}(\ell, j) \right] \right. \\ (-1)^j \left[\left(-P_5(i, j, \ell, m) + P_7(i, j+1) \right) / \lambda_{m3}(\ell, j, m) \right. \\ \left. \left(P_3(\ell-1, i+1, m-1) - P_4(i) \right) / \lambda_{m1}(\ell, m) \right] \\ (-1)^m \left[\left(-P_5(i, j, \ell, m) + P_7(\ell-1, m) \right) / \lambda_{m3}(i, j, m) \right. \\ \left. \left(P_3(i, \ell, j) - P_4(\ell-1) \right) / \lambda_{m1}(i, j) \right] \\ (-1)^j \left[\left(P_5(i, j, \ell, m) - P_7(\ell-1, m) \right) / \lambda_{m3}(i, j, m) \right. \\ \left. \left. \left(-P_3(\ell-1, i+1) + P_7(\ell-1, m) \right) / \lambda_{m1}(i, m) \right] \right]. \quad (45)$$

Here the definition

$$P_5(i, j, \ell, m) \equiv \frac{\left[(a_{n23})^{(-1)^\ell} (a_{n'23})^{(-1)^m} - (a_{n23})^{(-1)^{i+1}} (a_{n'23})^{(-1)^{j+1}} \right]}{\bar{\lambda}_n \left[(-1)^{i+1} + (-1)^{\ell+1} \right] + \left[(-1)^{j+1} + (-1)^{m+1} \right] \bar{\lambda}_{n'}} \quad (46)$$

is used.

4. Conclusion

The derivation of the formula for the impedance of a toroid sample in a toroid has been completed. In (21), one can note that the Z is a dimensionless quantity. To make it have the dimensions of ohms, we must multiply Z by Z_0 , where $Z_0 = \sqrt{\frac{\mu_0}{\epsilon_0}} = 376.7\Omega$ is the free space impedance. In addition, if one examines (9) and (10), one can see that under the circumstance that the ferrite sample has $r_1 = r_2$ and $r_3 = r_4$, that $A_{nn'} = 0$ when $n' \neq n$. This means that the formula (21) is now mathematically exactly correct. However, the physical effects of the closeness of the winding on the outside toroid becomes a major source of error. This means that the condition under which the large disk toroid may be used must be such that the effects of the winding can be neglected. Such condition can be accurately determined only if one used a theory for which the non-uniformity of the windings has been explicitly accounted for. Rather than to have to develop such a theory, we avoid such needs by using the present formula in a way that will make them sufficiently accurate. This may be done by mathematically varying the sample size relative to the outside current sheet toroid. When the double sum correction in the impedance formula (21) is less than some figure of merit, say 2%, of the total correct impedance, one can be reasonably sure that the resulting impedance values are known to 0.5%.*

* Note - The criteria used to establish the 0.5% figure is developed in the FORTRAN subroutine IMPED. Here the double sum series is truncated by the requirement that the last correction to the double sum correction be less than 10%. The single sum series is truncated by the requirement that the last correction to the single sum part be less than 0.5%. Because the two series have alternating signs and apparently converge like $1/n^2$ for the double sum and like $1/n^3$ for the single sum, the error made in the evaluation of the impedance is determined by the value of the largest finally accepted correction - namely 0.5%.

In the theory and experiment the effects of the reactive coupling of the wire coil to the sample can be minimized by placing the sample as far as possible from the wire coil--namely the exact center of the coil. To make explicit comparison of the theory with the experiment, one must have the following operational procedure: first, the experiment is done with and without the sample. The two resulting impedances are then subtracted to produce an effective impedance Z_{eff} . The reason for this subtraction is to attempt the removal of the effects of external fringing in the impedance results. These effects are due to the coil arrangement which allows magnetic flux to exist outside the toroidal cross section. For this operation to be meaningful the assumption is made that the fringing field is the same for the system, with or without the sample. This condition is the same as the theoretical assumption that one can use a uniform current sheet to characterize the internal impedance of the toroid. However, to have a more precise relationship to the experimental procedure, one must produce a Z_{eff} (for the theory) in exactly the same manner as was the Z_{eff} produced for the experiment. This means that we compute the case where the ferrite core has $\mu = \epsilon = 1$ to obtain Z_{vac} and then calculate for the case of μ, ϵ equal the desired values to get Z_t .^{*} The effective impedance is given by $Z_{\text{eff}} = Z_t - Z_{\text{vac}}$. This result may now be compared with the experimental effective impedance. To get a reliable comparison

* Note - The experiment produces two numbers, while we need four numbers to fix the theoretical impedance. It should be obvious that to fix the four numbers it is necessary to do a second measurement under known conditions, so that we have sufficient experimental data to get the μ and ϵ . The second measurement is to use a parallel plate capacitor.

one should have several sample sizes for the same material. By assuming that the material is sufficiently homogenous, this comparison will give a check on the accuracy of the above correspondence.

The appendix I gives the FORTRAN and FLOW chart needed to calculate the correct impedance. The main program, TOROID, gives the definitions and indicates the needed input. The number appearing on the far right of the listing corresponds roughly to the number of the same equation in the text. This program takes about 0.1 minute per point. If one is to do very many points, say 6000 or more, it would be desirable to make the program more efficient. The dimensions of the system must be such that
$$\left(\frac{Z_4 - Z_1}{r_4 - r_1} \right) 3n \leq \ln \left[\text{Largest number allowed} \right]$$

in the computing machine.] The n is the largest number of modes needed to satisfy the figure of merit. (Usually $n \leq 7$.)

Because of the transcendental nature of the Equation (21), we cannot solve for the μ , ϵ , explicitly. This means that the best one can do is to prepare a table of impedance in which the values of μ and ϵ are varied over ranges of interest. Such a table will be forthcoming when the needs of the various personnel have been determined.

To aid in the use of this program, sample output is included in appendix II. The reader should note that the 1.5×10^6 Hz case strongly violates the region of validity of the theory.

APPENDIX I

```

'EQUIP,5=60
'EQUIP,6=61
'FTN,X,*
C MAIN PROGRAM FOR TOROID CALCULATION TOROID
• CTOROID
TYPE COMPLEX CZII,CFIX,CZ,CZI,CZIII,ZC,CR,CV,DUM,DELV,W
COMMON R,Z,V,FIX, FR ,ZI,ZII,IT,IC,ITT ,FR2,ICYLE ,DUM,ZC TOROID
DIMENSION IC(6),FR2( 5004)
DIMENSION ZII(2),FIX(2),Z(8),CZ(4),ZI(2),ZIII(2),ZC(58),CR(5),
1 R(10),CV(2),V(4),DUM(10),DELV(2),W(2),UM(20)
EQUIVALENCE (CZII,ZII),(CFIX,FIX),(CZ,Z),(CZI,ZI),(CZIII,ZIII),
1 (CR,R),(CV,V),(DUM,UM)
DO 980 I=1,5796
980 R(I)=0.
1 READ(5,2)(R(I),I=1,8,2),(Z(I),I=1,8,2),TURN,IC(5),IC(6)
2 FORMAT( 4E18.5/4E18.5/E18.5,4I4)
IF(EOF,5)100,999
999 CONTINUE
IT=-1 TOROID
101 READ INPUT TAPE 5,102,ER,IC(4),IS,ITES
IF(EOF,5)100,200
200 CONTINUE
FR=ER*(2.0944E-10)
102 FORMAT(E18.5,3I4)
C 1THE INPUT IS THE RADIUS VARIABLES,THE Z AXIS VARIABLES,THE MAGNETITOROID
C 2C PERMIBILITY,AND DIELECTRIC PERMITIVITY,TURN COUNT,AND FREQUENCY TOROID
C 3ALL DIMENSIONS ARE ASSUMED TO BE IN CM. THE FREQUENCY IS HERTZ.
C FREE SPACE MU AND EPSILON IS ONE. THE INPUT FORMAT IS FLOATINGTOROID
C V(1)+IV(3) IS THE DIELECTRIC PERMITIVITY AND V(2)+IV(4) IS THE
C COMPLEX PERMIBILITY
DO 4 I=2,8,2
4 Z(I) = R(I) = 0.
FIX(1)=0. TOROID
6 WRITE(6,7)IS,ER,TURN,(R(I),I=1,8,2),(Z(I),I=1,8,2)
7 FORMAT( 44H1IMPEDANCE RESULTS FOR TOROID GEOMETRY PAGE I4/ 1TOROID
18H THE FREQUENCY IS E18.4,25HHZ. THE TURNS COUNT IS F4.0,/ 2TOROID
27H THE RADIAL DIMENSIONS ARE 4E18.4 / 2TOROID
37H THE Z AXIS DIMENSIONS ARE 4E18.4 // 9TOROID
41H THE LENGTH MEASUREMENTS ARE IN CM. IMPEDANCES ARE IN OHMS. THE
5E REST ARE DIMENSIONLESS. //115H MU
6 EPSILON QUASI STATIC IMP CORRECT IMP
7 DBL SUM IMP /125H U1 +I U2 E1 +I E2
8 Z1 +I Z2 ZC1 +I ZC2 DZ1
9 +I DZ2 CNT )
ITT=1
9 IS=IS+1 TOROID
IF(IT) 71,71,13
71 IV=IC(4) TOROID
72 FIX(2)= FR*(59.95 )*(TURN**2)
DO 13 I1=1,IV TOROID
211 FORMAT(4E18.5)
READ(5,211)V(1),V(3),V(2),V(4)
CALL IMPED TOROID
CZII=CFIX*CZII

```

	CZI=CZII+CFIX*CZI	
	CZIII=ZC(1)*CFIX	
	WRITE(6,8)V(3),V(4),V(1),V(2),ZIII,ZI,ZII,ICYLE	
	ITT=ITT+1	TOROID
8	FORMAT(E13.4,9E12.4,I4)	
	IF(ITT-46)13,6,6	
13	CONTINUE	
	IF(ITES)1,1,103	TOROID
103	IT=0	TOROID
	GO TO 101	TOROID
100	END	
	SUBROUTINE IMPED	IMPED
	TYPE COMPLEX CFR2,R,Z,VE,Z12,Z23,Z34,ZA1,ZA,SV,FR3,ZC,AD,AE,DUM,	
1	AN,XI,XJ,A4,B1,D1,A5,A12,A34,A23,CB,ZB,V,FIX,ZI,ZII,	
2	F12,G34,ALAM	
	DIMENSION IC(6),TD(100,2),WK(100),T(100,2),ADI(100),TW(8)	IMPED
	DIMENSION FR2(2),R(5),Z(4),ZC(58),DUM(10),AN(100,5),A4RL(2),B1(100	
1),D1(100),A5RL(2),A12(100,2),A34(100,2),A23(100,2),	
2	UM(20),	
2	CB(100,8),V(2),F12(100),ALAM(100,2),G34(100)	
	EQUIVALENCE (CFR2,FR2),(A4,A4RL),(XI,XIRL),(XJ,XJRL),(A5,A5RL),	
	1(DUM,UM)	
	COMMON R,Z,V,FIX,FR,ZI,ZII,IT,IC,ITT,FR2,FR3,B1,D1,AN, F12,G34,	IMPED
	1ALAM,CB,A23,A34,A12,ICYLE,DUM,ZC,TD,WK,T,ADI,Z12,Z23,Z34	IMPED
	2,TW,SV,VE	
	ICYLE=0	IMPED
	IF OVERFLOW FAULT 10,10	
10	IF(IT)1,101,2	IMPED
1	Z34=Z(3)-Z(4)	IMPED
	Z12=Z(1)-Z(2)	IMPED
	Z23= Z(2)-Z(3)	IMPED
	ZA1=Z23*CLOG(R(3)/R(2))	IMPED
	ZA= (Z(4)-Z(1))*CLOG(R(4)/R(1))	IMPED
	FR2(2)=0.0	IMPED
	IA=IC(5)	IMPED
101	FR2(1)=(FR)**2	IMPED
	CALL TXX(1,1)	IMPED
	IT=1	IMPED
2	SV=V(2)-(1.0,0.0)	IMPED
	VE= V(1)*V(2)-(1.0,0.0)	IMPED
	FR3=CFR2*VE	
	N2=1	IMPED
	CALL SCX(N2,1)	IMPED
102	IF(N2)42,42,201	IMPED
201	DO 3 I=2,58	IMPED
3	ZC(I)=(0.0,0.0)	IMPED
	ZC(1)=ZA-ZA1*SV	IMPED
	DO 40 N=1,IA	IMPED
	N1=N+1	IMPED
	N2=N	IMPED
	IB=IC(6)+N	IMPED
	IF (IB .GT. 100) IB=100	
301	DO 303 L=25,28	IMPED
303	ZC(L)=(0.0,0.0)	IMPED

	DO 30 M=N1,IB	IMPED
	ICYCLE=ICYCLE+1	IMPED
	M2=M	IMPED
	CALL TXX(N2,M2)	IMPED
	CALL SCX(N2,M2)	IMPED
103	IF(N2)42,42,304	IMPED
304	AD= AN(M,1)*DUM(8)	IMPED
	AE= AN(M,2)*DUM(8)	IMPED
	ZC(35)= DUM(4)*AD*AN(N,1)	IMPED
	ZC(36)= DUM(5)*AE*AN(N,1)	IMPED
	ZC(37)=DUM(6)*AE*AN(N,2)	IMPED
	ZC(38)= DUM(7)*AE*AN(N,2)	IMPED
	ZC(26)=ZC(35)+ZC(36)*SV+(ZC(37)+ZC(38)*SV)*VE	IMPED
	ZC(25)=ZC(25)+ZC(26)	IMPED
1000	PRINT 1001,M ,AD,AE,AN(M,1),AN(M,2),SV,ZC(25),ZC(26),(ZC(I),I=35,38	
	X8), (DUM(I),I=4,8)	
1001	FORMAT(6H 1001 I2,5C(E12.4,E12.4)/(8X,5C(E12.4,E12.4)))	
	XI=CABS(ZC(26))	
	XJ=CABS(ZC(25))	
	XK=XJRL*(0.1)-XIRL	
	IF(XK)30,31,31	
30	CONTINUE	IMPED
	ITT=ITT+2	PTEST
	A4=FR3*CFR2*FIX*ZC(26)	
	WRITE OUTPUT TAPE 6,11,N2,A4RL	
11	FORMAT(54H WARNING CONVERGENCE IS SUCH THAT ON DBLE SUM CYCLE	PTEST
1	14/25H THE LAST IMPEC COR. WAS E14.4, 4H +I E14.4)	PTEST
31	ZC(5)=ZC(5)+ZC(25)	
	ZC(12)= AN(N,1)*B1(N)*DUM(1)	IMPED
	ZC(13)= AN(N,1)*D1(N)*DUM(2)	IMPED
	ZC(14)= AN(N,2)*D1(N)*DUM(3)	IMPED
	ZC(4)=ZC(12)+ZC(13)*SV+(ZC(13)+ZC(14)*SV)*VE	
	ZC(2)=ZC(2)+ZC(4)	
	XI=CABS(ZC(4))	
	XJ=CABS(ZC(2))	
	XK=XJRL*(0.005)-XIRL	
	IF(XK)40,41,41	
40	CONTINUE	IMPED
	ITT=ITT+2	
	A5=CFR2*FIX*(ZC(4)+FR3*ZC(25))	
	WRITE OUTPUT TAPE 6,22,IA,A5RL	
22	FORMAT(57H WARNING TOTAL CONVERGENCE IMPLIES THE SINGLE SUM COUNT	PTEST
1	14/25H THE LAST IMPED COR. WAS E14.4, 4H +I E14.4)	PTEST
41	ZI=ZC(1)+CFR2*ZC(2)	
	?II=FR3*CFR2*ZC(5)	
	IF C OVERFLOW FAULT 43,42	
43	WRITE OUTPUT TAPE 6,44	
44	FORMAT(47H ACCUMULATOR OVERFLOW NEXT LINE HAS NO MEANING)	
42	RETURN	
	END	IMPED
	SUBROUTINE TXX(N,M)	TXX
	TYPE COMPLEX TWO,ONE,AVE,ZRO,CX22,CR,B1,D1,AN,ATEM,A12,A34,G34,F12	
1	,ALAM	
	DIMENSION X22(2),CR(5),R(10),B1(100),B1RL(200),D1(100),D1RL(200),	

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1      AN(100,5),A12(100,2),A34(100,2),F12(100),G34(100),
2      ALAM(100,2),ANRL(1)
DIMENSION IC(6),TD(100,2),WK(100),T(100,2),ADI(100)          IMPED
1,TB(2)      ,T11(4),T12(2)
EQUIVALENCE (CX22,X22),(CR,R),(B1,B1RL),(D1,D1RL),(AN,ANRL)
COMMON R,Z,V,FIX,FR,ZI,ZII,IT,IC,ITT,FR2,FR3,B1,D1,AN, F12,G34, IMPED
1ALAM,CB,A23,A34,A12 ,ICYCLE,DUM,ZC,TD,WK,T,ADI ,Z12,Z23,Z34 IMPED
2,ONE,TWO,ZRO,AVE
COMPLEX Z,V,FIX,ZI,ZII,FR2,FR3,CB,A23,DUM,ZC,Z12,Z23,Z34
DIMENSION Z(4),V(2),CB(100,8),A23(100,2),DUM(10),ZC(58)
IF(IT )1,101,2
1      N1=1          TXX
      TWO=(2.0,0.0)  TXX
      ONE=(1.0,0.0)  TXX
      AVE=(0.5,0.0)  SCX
      ZRO=(0.0,0.0)  SCX
      X22(2)=0.0     TXX
      R(9)=1.5/(R(7)-R(1))
      R(10)=0.1 *R(9)
101    N2=1
2      IF(M-N1)41,3,3
3      N1=N1+1      TXX
      B1RL(2*M)=D1RL(2*M) =ANRL( 2*M+400 ) = 0.
      WK(M)= BESSEL(5,1)          TXX
      WK1=WK(M)                   TXX
      X2=BESSEL(1,4)              TXX
      X3=BESSEL(1,2)              TXX
      X4= BESSEL(1,1)             TXX
      X1= BESSEL(1,3)             TXX
      DO 33 I=2,3                 TXX
      K=I-(-1)**I                 TXX
      J=I-1                       TXX
      TD(M,J)=(BESSEL(I,4)*X1-BESSEL(I,3)*X2)*WK1 TXX
      TB(J)=BESSEL(K,1)*X2-BESSEL(K,2)*X1
      T(M,J)=BESSEL(I,1)*X2-BESSEL(I,2)*X1
      T12(J)=BESSEL(I,1)*X3-BESSEL(I,2)*X4
      T11(I)=BESSEL(I,3)*X3-BESSEL(I,4)*X4
33     T11(K)=BESSEL(K,3)*X3-BESSEL(K,4)*X4
      AN2=0.0                      TXX
      B1(M)=TB(2)-TB(1)            TXX
      D1(M)=T(M,2)-T(M,1)         TXX
      WK2=WK(M)**2                 TXX
      ADI(M)=0.0                   TXX
      DO 35 I=2,3                 TXX
      IF(I.EQ.2)340,341
340    K=KK=I      $ II=3 $ GO TO 342
341    K=4 $ KK=7 $ II=5
342    J=I-1
6      AN2=AN2-(-1.0)**I*(TB(J)*R(KK))*(TB(J)*R(KK)+R(1)*T11 ( K )) TXX
35     ADI(M)=ADI(M)-(-1.0)**I*R(II)*(T(M,J)*WK2*(R(II)*T(M,J)+R(1)*T11 ( TXX
1I ))+TD(M,J)*(2.0*T(M,J)+R(II)*TD(M,J)+ WK(M)*R(1)*T12 ( J))) TXX
5      ADI(M)=0.5*ADI(M)          TXX
4      AN(M,3)= 1.0/(WK2*AN2)    TXX 18
      ADI(M)=ADI(M)*AN(M,3)      TXX
3
2

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	X22(1)= WK2-FR2	TXX
	GO TO 39	
41	IF(M-N2)4,38,38	
38	X22(1)=WK(M)**2-FR2	
39	ALAM(M,1)= CSQRT(CX22)	
	N2=N2+1	
	A12(M,1)= CEXP(ALAM(M,1)*Z12)	
	A34(M,1)= CEXP(ALAM(M,1)*Z34)	
	A34(M,2)=ONE/A34(M,1)	TXX
	A12(M,2)= ONE/A12(M,1)	TXX
	G34(M)=(TWO-A34(M,1)-A34(M,2))/ALAM(M,1)	TXX
	F12(M)=(A12(M,1)+A12(M,2)-TWO)/ALAM(M,1)	TXX
	ATEM =AN(M,3)/ALAM(M,1)	
	AN(M,1)=B1(M)*ATEM	
	AN(M,2)=D1(M)*ATEM	
	AN(M,4)=ONE	
4	RETURN	TXX
	END	TXX
	FUNCTION BESSEL(I,K)	
	DIMENSION R(10),B(12,4)	
	DIMENSION S(19)	
	COMMON R,S,IT	BESSEL
	IF(IT)100,1010,1010	BESSEL
100	A5= R(9)	
	R(10)=5.0*R(10)	
	IT=1	
	ISV=2	
	IM=1	
	A4= 1.0E+04	BESSEL
1010	GO TO (10,10,10,10,1),I	BESSEL
1	DEL=R(10)	BESSEL
14	R(9)=R(9)+DEL	
141	X=R(9)*R(1)	
	Y= X/3.0	BESSEL
	IF(Y-1.0) 101,110,110	BESSEL
101	I1=5	BESSEL
	GO TO 3	BESSEL
15	X=R(9)*R(7)	
	Y= X/3.0	BESSEL
	IF(Y-1.0) 102,103,103	BESSEL
102	I1=6	BESSEL
	GO TO 3	BESSEL
103	I1=5	BESSEL
	GO TO 4	BESSEL
16	ANULL=B(5,3)*B(6,4)-B(6,3)*B(5,4)	BESSEL
	GO TO 71	BESSEL
20	ANULL=(B(5,3)*COSF(PH1)+B(5,4)*SINF(PH1))*(-F1)	BESSEL
	GO TO 71	BESSEL
110	I1=6	BESSEL
	Y=1.0/Y	BESSEL
	GO TO (41,400), ISV	
21	PH2=PH1	BESSEL
	F2=F1	
	I1=7	BESSEL


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X=R(9)*R(7)
Y= 3.0/X
GO TO (41,400), ISV
22 ANULL= PH1-PH2
ANULL=F2*F1*SINF (ANULL)
GO TO 71
2 Y= X/3.0
IF(Y-1.0)3,3,4
3 Y= Y**2
IF ( X ) 910,910,912
910 X4=0. $ GO TO 914
912 X4=0.63661977*LOGF(X/2.0)
914 B(I1,3)= X*(0.5-Y*(0.56249985-Y*(0.21093573-Y*(0.03954289-Y*(0.000
1443319-Y*(0.00031761-Y*(0.00001109))))))
B(I1,4)= X4*B(I1,3) -((0.6366198-Y*(0.2212091+YBESSEL
1*(2.1682709-Y*(1.3164827-Y*(0.3123951-Y*(0.0400976-Y*(0.0027873)))))))/ X
GO TO (6,6,6,6,15,16),I1
4 Y=1.0/Y
400 X1= 1.0/SQRTF(X)
F1= X1 *(0.79788456+Y*(0.00000156+Y*(0.01659667+Y*(0.00017105
1-Y*(0.00249511-Y*(0.00113653-Y*(0.00020033))))))
41 PH1= X-(0.78539816-Y*(0.12499612+Y*(0.00005650-Y*(0.00637879-Y*(
10.00074348+Y*(0.00079824-Y*(0.00029166))))))
GO TO(42,42,42,42,20,21,22),I1
42 B(I1,3)= F1*SINF(PH1)
B(I1,4)=-F1*COSF(PH1)
GO TO 7
6 B(I1,1)=1.0-Y*(2.2499997-Y*(1.2656208-Y*(0.3163866-Y*(0.0444479-Y*(
1(0.0039444-0.0002100*Y))))))
B(I1,2)= X4*B(I1,1)+0.36746691+Y*(0.60559366-Y*(0.74350384-Y*(
10.25300117-Y*(0.04261214-Y*(0.00427916-Y*(0.00024846))))))
GO TO 8
7 F0= X1*(0.79788456-Y*(0.00000077+Y*(0.00552740+Y*(0.00009512-Y*(
10.00137237-Y*(0.00072805-Y*(0.00014476))))))
PH0=X-(0.78539816+Y*(0.04166397+Y*(0.00003954-Y*(0.00262573-Y*(
10.00054125+Y*(0.00029333-Y*(0.00013558))))))
B(I1,1)=F0*COSF(PH0)
B(I1,2)=F0*SINF(PH0)
GO TO 8
71 IF(SIGNF(ANULL,A4)-ANULL)741,74,741
74 IM=1
740 A4=ANULL
GO TO 14
741 GO TO (742,740),IM
742 R(9)=R(9)+(DEL*ANULL)/(A4-ANULL)
DEL=(-0.5)*DEL
A4=ANULL
ANULL=ABSF(ANULL)
6 IF(ANULL-1.0E-05)75,75,141
5 75 B(5,1)=R(9)
GO TO (77,751),ISV
4 751 IF(R(9)*R(1)/3.0-1.0)77,76,76
3 76 ISV=I
2

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F1=1.0
77 DO 8 J=1,4
    I1=J
    JJ=J+J-1
    X=R(JJ)*R(9)
    GO TO 2
8 CONTINUE
    IM=2
    R(9)=R(9)+ A5
10 BESSEL =B(I,K)
    RETURN
END
SUBROUTINE SCX(N,M)
TYPE COMPLEX A12,A23,A34,ADD,AH,AH1,AH2,AH3,AH4,AK1,AK2,AK4,AK5,
1 AK6,AK7,AK8,AL2,AL3,AL4,AL5,ALAM,ALM,ALM2,ALM3,AN,AVE
2 ,B1, CB, D1,DIV,DUM, F12,F23,FIX,FR2,FR3, G23,G34,ONEC
3 ,R1,R2,RD1,RD2,RS1,RS2, TWO, V, X1,X2,XA,XAL,XP1,XP2,
4 XP3,XP4,XP5,XP7,XPP,XT,XY, Y2, Z,Z23,ZB,ZC,ZI,ZII,ZRO
5 ,Z12,Z34
DIMENSION DUMRL(20),R(10),DIVRL(2),
1 A12(100,2),A23(100,2),A34(100,2),ALAM(100,2),ALM(2,2),
2 ALM2(2,2,2),ALM3(2,2,2),AN(100,5),B1(100),CB(100,8),
3 D1(100),DUM(10),F12(100),G34(100),XP1(2,2),XP2(2,2),
3 XP3(2,2,2),XP4(2),XP5(4,4),XP7(2,2),Z(4),ZC(58),V(2)
DIMENSION IC(6),TD(100,2),WK(100),T(100,2),ADI(100)
EQUIVALENCE(DUM,DUMRL),(DIV,DIVRL)
1 ,(ONEC,ONE)
COMMON R,Z,V,FIX,FR,ZI,ZII,IT,IC,ITT,FR2,FR3,B1,D1,AN, F12,G34,
1ALAM,CB,A23,A34,A12 ,ICYLE,DUM,ZC,TD,WK,T,ADI ,Z12,Z23,Z34
2,ONEC,TWO,ZRO,AVE
7001 IF( ICYLE ) 1,1,2
1 N3=1
IT=IT+1
2 IF(M-N3)4,3,3
3 IF DIVIDE CHECK 31,31
31 N3=N3+1
IF(IT-2)12,11,12
11 AN(M,5)=ONE
GO TO 311
12 AN(M,5)=ONE/AN(M,4)
IF DIVIDE CHECK 310,311
310 AN(M,5)=(1.0,0.0)
311 ADD=ADI(M)
ALAM(M,2)= CSQRT(ALAM(M,1)**2-ADD*FR3)
A23(M,1)= CEXP( ALAM(M,2)*Z23)
A23(M,2)= ONE/ A23(M,1)
IF DIVIDE CHECK 5,32
32 R1= ALAM(M,1)/ALAM(M,2)
R2= ONE/R1
RS1= R1+ONE
RD1= R1-ONE
RS2= R2+ONE
RD2= R2-ONE
CB(M,1)= AVE*( RS1*A34(M,1)+RD1*A34(M,2))

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	CB(M,2)=(-AVE)*(RD1*A34(M,1)+RS1*A34(M,2))	SCX
	CB(M,3)=(AVE)*(RS1*A12(M,2)+RD1*A12(M,1))	SCX
	CB(M,4)=(-AVE)*(RD1*A12(M,2)+RS1*A12(M,1))	SCX
	CB(M,5)=(AVE)*(CB(M,1)*RS2*A23(M,1)-CB(M,2)*RD2*A23(M,2))	SCX 15
	CB(M,6)=(AVE)*(CB(M,2)*RS2*A23(M,2)-CB(M,1)*RD2*A23(M,1))	SCX
	CB(M,7)=(AVE)*(CB(M,3)*RS2*A23(M,2)-CB(M,4)*RD2*A23(M,1))	SCX
	CB(M,8)=(AVE)*(CB(M,4)*RS2*A23(M,1)-CB(M,3)*RD2*A23(M,2))	SCX
	XA= AVE/ (CB(M,5)*A12(M,1)+CB(M,6)*A12(M,2))	
	IF DIVIDE 5,34	SCX
34	AN(M,4)=CSQRT(XA)	
	AN(M,5)=AN(M,5)*AN(M,4)	
	DO 35 I=1,8	
35	CB(M,I)=CB(M,I)*AN(M,4)	
	F12(M)=F12(M)*AN(M,5)	
	G34(M)=G34(M)*AN(M,5)	
4	IF(M-N-1)500,41,430	SCX
41	G23=(CB(N,2)*(A23(N,2)-ONE)-CB(N,1)*(A23(N,1)-ONE))/ALAM(N,2)	SCX
	F23=(CB(N,3)*(A23(N,2)-ONE)-CB(N,4)*(A23(N,1)-ONE))/ALAM(N,2)	SCX
	AK1=ZRO	SCX
	AK4=ZRO	SCX
	AK2=F12(N)*G23+F23*G34(N)	SCX 23
	DO 43 I=1,2	SCX
	I1=I+4	SCX
	DO 43 J=1,2	SCX
	J2=J+6	SCX
	J1=J+2	SCX
	K=3-J	SCX
	XY=(-ONE)**J-(-ONE)**I	SCX
	AK1=AK1-(((ONE)**J)*CB(N,I1)*(A12(N,K)*Z12*XY+	SCX 35
	1((-ONE)**(I+J))*(A12(N,I)+A12(N,K)-TWO)/ALAM(N,1))SCX
	2+((-ONE)**I)*CB(N,J2)*(A34(N,K)*Z34*XY+	SCX
	3((-ONE)**(I+J))*(A34(N,I)+A34(N,K)-TWO)/ALAM(N,1))SCX
43	AK4=AK4+CB(N,I)*CB(N,J1)*(A23(N,K)*Z23*XY+	SCX 34
	1((-ONE)**(I+J))*(A23(N,I)+A23(N,K)-TWO)/ALAM(N,2))	SCX
	DUM(3)=AK4/ALAM(N,2)	SCX
	DUM(2)=DUM(3)+AK2	SCX 24
	AK1=AK1*AN(N,4)	
	DUM(1)=DUM(2)+AK2+TWO*F12(N)*G34(N)+AK1/ALAM(N,1)	SCX 22
	WK2=WK(N)**2	SCX
430	WK2P=WK(M)**2	SCX
	DUMRL(15)=((R(5)*T(N,2)*TD(M,2)-R(3)*T(N,1)*TD(M,1))*WK2	
	1-(R(5)*T(M,2)*TD(N,2)-R(3)*T(M,1)*TD(N,1))*WK2P)/(WK2-WK2P)	SCX
4308	DUMRL(16)=0.0	
	DO 431 I=4,7	SCX
431	DUM(I)=ZRO	SCX
	DO 499 I1=1,2	SCX
	GO TO(432,433),I1	SCX
432	N1=N	SCX
	N2=M	SCX
6	DO 4321 I2=1,2	SCX
	DO 4321 I3=1,2	SCX
5	ALM(I2,I3)=((-ONE)**(I2+1))*ALAM(N1,2)+((-ONE)**(I3+1))*ALAM(N2,2)	SCX 32A
4	XP1(I2,I3)=(ONE-A23(N1,I2)*A23(N2,I3))/ALM(I2,I3)	SCX32A
	IP=3-I3	SCX
3		
2		

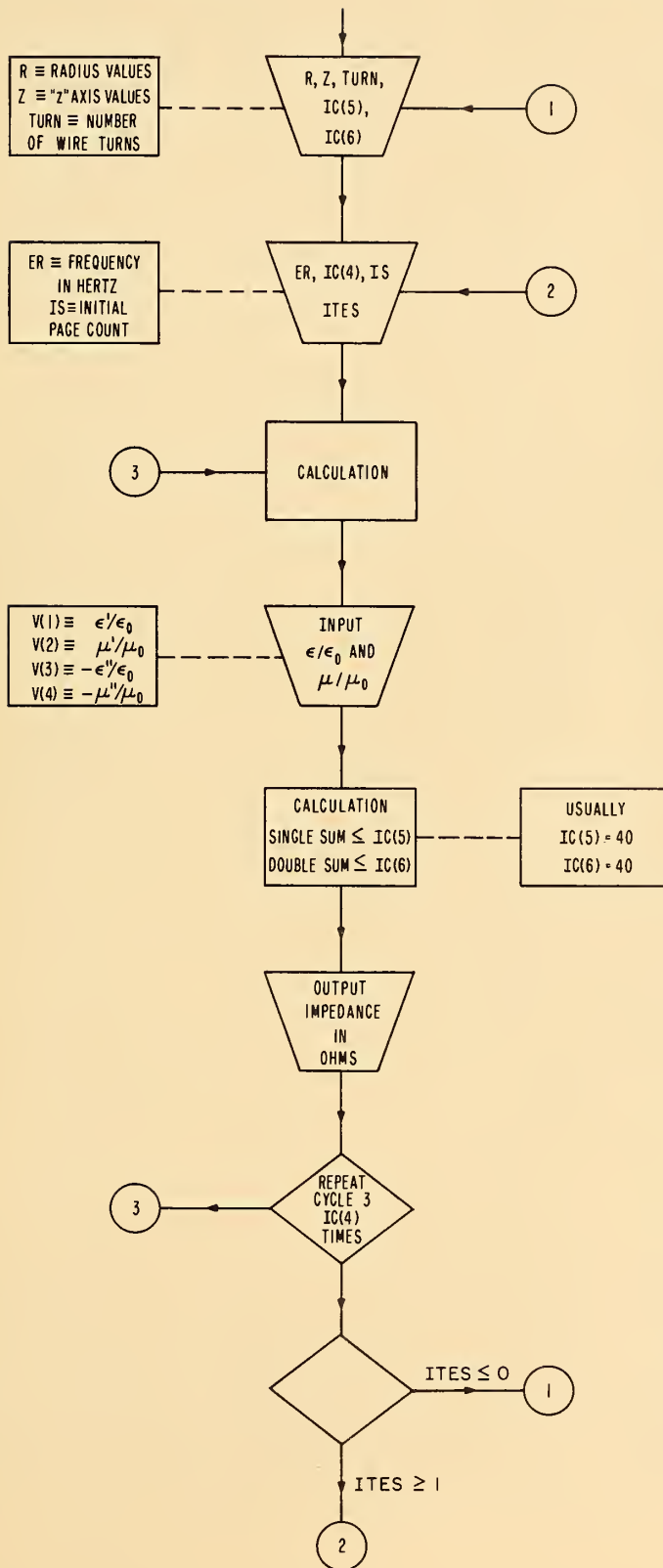
4321	XP2(I2,I3)=(A23(N2,IP)-A23(N1,I2))/ALM(I2,I3)	SCX
	AH4=ZRO	SCX
	AH1=ZRO	SCX
	AH2=ZRO	SCX
	AH3=ZRO	SCX
	DO 4322 I2=1,2	SCX
	I6=I2+2	SCX
	I4=3-I2	SCX
	DO 4322 I3=1,2	SCX
	I5=3-I3	SCX
	I7=I3+2	SCX
	AH2=AH2+(CB(N1,I6) *CB(N2,I7) *XP1(I4,I5))	SCX 33
	AH3=AH3+(CB(N1,I2) *CB(N2,I7) *XP2(I2,I3))	SCX
	AH4=AH4+(CB(N1,I6) *CB(N2,I3) *XP2(I4,I5))	SCX
4322	AH1=AH1+(CB(N1,I2) *CB(N2,I3) *XP1(I2,I3))	SCX
	IF (ALAM(N2,2)-2.*ALAM(N1,2)) 438,4341,438	
4341	IF((0.,1.)*(ALAM(N2,2)-2.*ALAM(N1,2))) 438,4342,438	
4342	ISET=1	SCX
	GO TO 4381	SCX
438	ISET=2	SCX
4381	DO 45 I=1,2	SCX
	IP=3-I	SCX
	DO 45 J=1,2	SCX
	JP=3-J	SCX
	DO 45 L=1,2	SCX
	LP=3-L	SCX
	L2=2*(L-1)+I	SCX
	Y2=(-ONE)**I+(-ONE)**J	SCX
	ALM2(I,J,L)=- (ALAM(N1,2)*Y2+ALAM(N2,2)*(-ONE)**L)	SCX 40
	ALM3(L,I,J)=- (ALAM(N2,2)*Y2+ALAM(N1,2)*(-ONE)**L)	SCX 41
	DO 45 K=1,2	SCX
	K2=2*(K-1)+J	SCX
	KP=3-K	SCX
	DIV= (ALM(I,J)+ALM(L,K))	SCX
	IF(CABS(DIV) .GT. 1.E-38) GO TO 449 \$ DIV=0.	
448	XP5(L2,K2)=A23(N1,LP)*A23(N2,KP)*(-Z23)	SCX 46
	GO TO 45	SCX
449	XP5(L2,K2)=(A23(N1,LP)*A23(N2,KP)-A23(N1,I)*A23(N2,J))/DIV	SCX
45	CONTINUE	SCX
	GO TO 65	SCX
433	N2=N	SCX
	ALM(1,2)=-ALM(1,2)	SCX
	ALM(2,1)=-ALM(2,1)	SCX
	N1=M	SCX
	XA=XP1(1,2)	SCX
	XP1(1,2)=XP1(2,1)	SCX
	XP1(2,1)=XA	SCX
	XA=XP2(1,1)	SCX
	XP2(1,1)= XP2(2,2)	SCX
	XP2(2,2)= XA	
	XA=XP2(1,2)	SCX
	XP2(1,2)= XP2(2,1)	SCX
	XP2(2,1)= XA	SCX
	AH=AH3	SCX

	AH3=AH4	SCX
	AH4=AH	SCX
	DO 62 I=1,3	SCX
	IA=4-I	SCX
	DO 62 L=1,IA	SCX
	J=I+L	SCX
	XPP=XP5(I,J)	SCX
	XP5(I,J)=XP5(J,I)	SCX
	XP5(J,I)=XPP	SCX
62	CONTINUE	SCX
	DO 64 J=1,2	SCX
	DO 64 I=1,2	SCX
	DO 64 K=1,2	SCX
	XAL=ALM2(I,J,K)	SCX
	ALM2(I,J,K)=ALM3(K,I,J)	SCX
64	ALM3(K,I,J)=XAL	SCX
	ISET=2	SCX
65	DO 70 I=1,2	SCX
	IP=3-I	SCX
	XP4(I)=(ONE-A23(N1,I))/(ALAM(N1,2)*((-ONE)**(I+1))	SCX
	DO 70 J=1,2	SCX
	JP=3-J	SCX
	XP7(I,J)=A23(N2,JP)*XP4(I)	SCX 37
	DO 70 K=1,2	SCX
	GO TO (69,692),ISET	SCX
69	IF((-1.0)**I+(-1.0)**J+2.0*(-1.0)**K) 692,691,692	SCX
691	XP3(I,J,K)=-A23(N1,JP)*Z23	SCX 42
	GO TO 70	SCX
692	XP3(I,J,K)=(A23(N1,JP)-A23(N1,I)*A23(N2,K))/ALM2(I,J,K)	SCX
70	CONTINUE	SCX
	AL2=ZRO	SCX
	AL3=ZRO	SCX
	AL4=ZRO	SCX
	AL5=ZRO	SCX
	AK8=ZRO	SCX
	DO 48 I=1,2	SCX
	IP=3-I	SCX
	DO 48 J=1,2	SCX
	J1=J+2	SCX
	JP=3-J	SCX
	XT =((-ONE)**I-((-ONE)**J)	SCX
	DO 48 L=1,2	SCX
	LP=3-L	SCX
	L1=L+2	SCX
	L2=2*(L-1)+1	SCX
	AL2=AL2+(CB(N2,I)) *CB(N1,L)*(XT*XP7(L,J)+((-ONE)**J*XP1(L,SCX	
	1I)) -((-ONE)**I*XP2(L,J)) *CB(N2,J1))	SCX
	AL3=AL3+(CB(N2,I)) *CB(N1,L1)*(XT*XP7(LP,J)+((-ONE)**J*XP2(SCX	
	1LP,IP)) -((-ONE)**I*XP1(LP,JP)) *CB(N2,J1))	SCX
5	AL4=AL4+CB(N1,I)*CB(N1,J1)*CB(N2,L)*((XP3(I,J,L)-XP7(I ,LP))	SCX
	1/ALM(J,L)+(XP4(JP)-XP3(I,J,L))/ALM(I,L))	
5	AL5=AL5+CB(N1,I)*CB(N1,J1)*CB(N2,L1)*((XP3(JP,IP,LP)-XP4(I))/ALM(SCX	
4	1J,L)+(XP7(JP,L)) -XP3(JP,IP,LP))/ALM(I,L))	SCX
	DO 48 K=1,2	SCX

	KP=3-K	SCX
	K1=K+2	SCX
	K2=2*(K-1)+J	SCX
	AK8=AK8-(CB(N1,I)*CB(N2,J)) *CB(N2,K1)*((-ONE)**K*(((XP5(L2SCX 45	
	1,K2)-XP7(I,JP))/ALM3(L,J,K))+(-XP3(I,L,J)+XP7(I,JP))/ALM(L,J)-((XPSCX	
	25(L2,K2)-XP7(LP,K))/ALM3(I,J,K))+XP3(I,L,J)-XP4(LP))/ALM(I,J)) SCX	
	3+((-ONE)**J*(-((XP5(L2,K2)-XP7(I,JP))/ALM3(L,J,K))+XP3(LP,IP,KP) SCX	
	4-XP4(I))/ALM(L,K)+((XP5(L2,K2)-XP7(LP,K))/ALM3(I,J,K))+((-XP3(SCX	
	5LP,IP,KP)+XP7(LP,K))/ALM(I,K))))) *CB(N1,L1) SCX	
48	CONTINUE	SCX
	AK8=AK8/ALAM(N2,2)	SCX
	AL2=AL2/ALAM(N2,2)	SCX
	AL3=AL3/ALAM(N2,2)	SCX
	AK6= F12(N1)*AL2+G34(N1)*AL3	SCX 27
	AK5= AK6+F12(N1)*(AH1*F12(N2)+AH3*G34(N2))+G34(N1)*(AH2*G34(N2)+	SCX 25
	1 AH4*F12(N2))	SCX
	AK7=AK8+F12(N2)*AL4+G34(N2)*AL5	SCX 26
	DUM(4)=DUM(4)+AK5	SCX
	DUM(5)=DUM(5)+AK6	SCX
	DUM(6)=DUM(6)+AK7	SCX
499	DUM(7)=DUM(7)+AK8	SCX
	DUM(5)=DUM(5)+DUM(7)	SCX
	DUM(4)=DUM(4)+DUM(6)	SCX
500	RETURN	SCX
5	N=-N	
	WRITE OUTPUT TAPE 6,51	SCX
51	FORMAT(54H OMEGA IS SUCH THAT STRUCTURE IS RESONANT-NOT ALLOWED)SCX	
	RETURN	SCX
	END	
	SCOPE	

INPUT DATA DEFINED

TOROID



APPENDIX II

IMPEDANCE RESULTS FOR TOROID GEOMETRY PAGE 1
 THE FREQUENCY IS 0.1000E 07HZ. THE TURNS COUNT IS 1.
 THE RADIAL DIMENSIONS ARE 0.4130E-00 0.6450E 00 0.1275E 01 0.2540E 01
 THE Z AXIS DIMENSIONS ARE -0.1590E 01 -0.4470E-00 0.4930E-00 0.1590E 01

THE LENGTH MEASUREMENTS ARE IN CM. IMPEDANCES ARE IN OHMS. THE REST ARE DIMENSIONLESS.

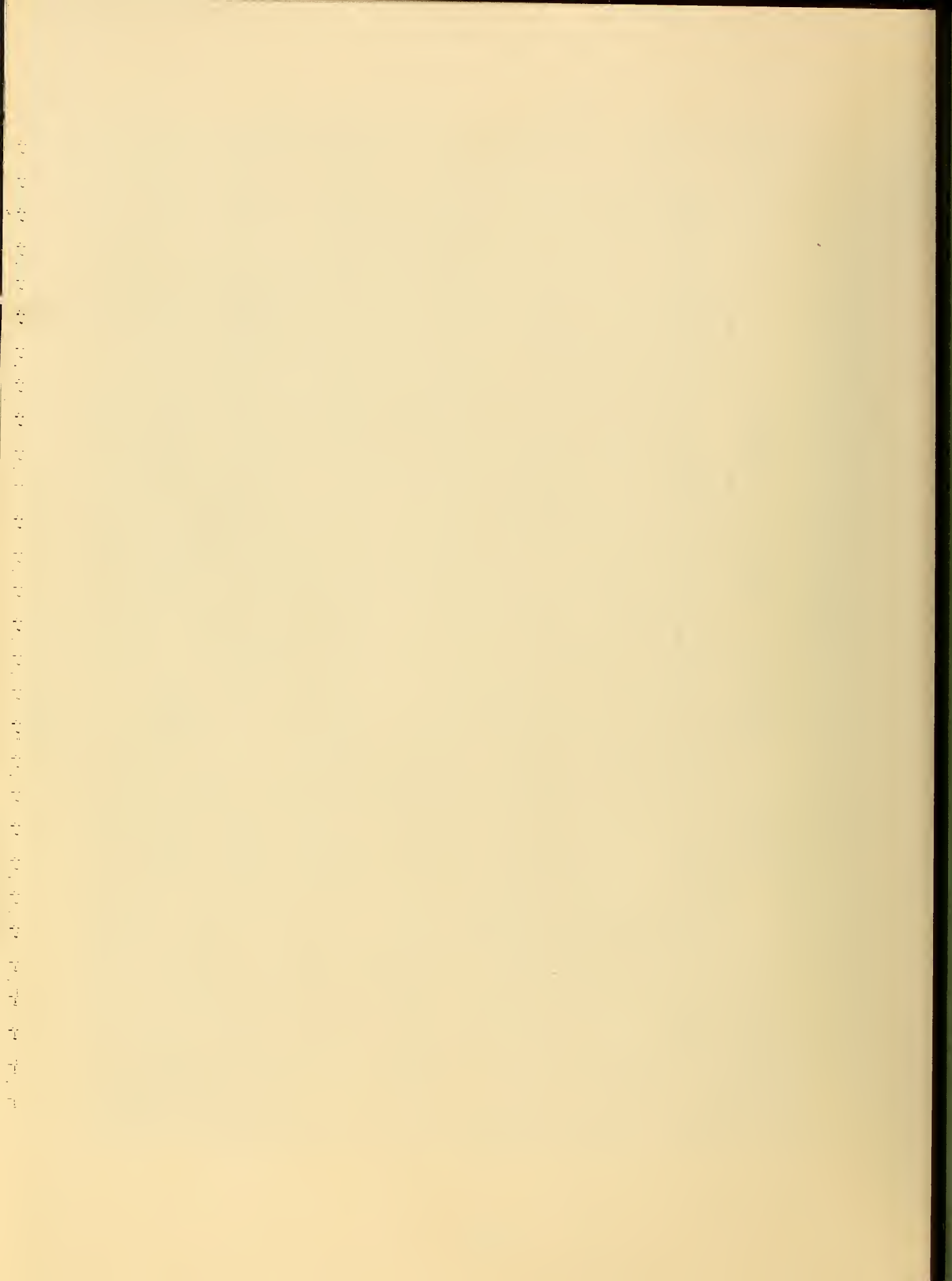
UI	MU	EPSILON	QUASI STATIC IMP	CORRECT IMP	DBL SUM IMP
0.2024E 04	+1 U2	E1	Z1	+1 ZC2	+1 DZ2
-0.1030E 04	-0.7540E 05	-0.2950E 05	+1 Z2	ZC1	DZ1
			0.8290E 01	0.1595E 02	0.1608E 02
			0.1635E 02	0.1608E 02	0.9130E 00
					-0.1048E 01
					UNT
					11

IMPEAOCE RESULTS FOR TOROID GEOMETRY PAGE 1

THE FREQUENCY IS 0.1500E 07HZ. THE TURNS COUNT IS 1.
 THE RADIAL DIMENSIONS ARE 0.4130E-00 0.6450E 00 0.1275E 01 0.2540E 01
 THE Z AXIS DIMENSIONS ARE -0.1590E 01 -0.4470E-00 0.4930E-00 0.1590E 01

THE LENGTH MEASUREMENTS ARE IN CM. IMPEAOCE\$ ARE IN OHMS. THE REST ARE DIMENSIONLESS.

#U	E1	E2	EPSILON	Z1	Z2	QUASI STATIC IMP	ZC1	ZC2	CORRECT IMP	DZ1	DZ2	DBL SUM IMP	CNT
01	0.1670E 04	0.7390E 05	-0.2880E 05	0.1280E 02	0.2026E 02	0.2026E 02	0.3090E 02	0.8581E 01	0.7944E 00	0.7310E 01	0.11		





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