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# EVALUATION OF CONVOLUTION INTEGRALS OCCURING IN THE THEORY OF MIXED PATH PROPAGATION

BY J.R. JOHLER AND C.M. LILLEY SCIENTIFIC REPORT NO. 14

Contract Title: Antenna Patterns in the Presence of Finite Ground Surfaces, CSO and A 58-40, dated March 31, 1958; Pro-61-568, dated February 2, 1961

Prepared for Electronics Research Directorate Air Force Cambridge Research Laboratories Office of Aerospace Research United States Air Force Bedford, Massachusetts

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November 8, 1961

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#### CONTENTS

ABS	STRACT	1
1.	Introduction	1
2.	Theory	2
3.	Computations	13
4.	Conclusions	13
REI	FERENCES	14

#### EVALUATION OF CONVOLUTION INTEGRALS OCCURRING IN THE THEORY OF MIXED PATH PROPAGATION

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#### ABSTRACT

The theory of propagation of electromagnetic waves around a sphere treats the smooth homogeneous case, i.e., the case in which the surface impedance of the sphere is uninterrupted by an abrupt change in conductivity such as a land/sea boundary. It is known, however, that such a theory can be extended to treat inhomogeneous, irregular terrain by formulating certain convolution integrals which utilize the smooth homogeneous formulas. The evaluation of these integrals can be accomplished with dispatch on a large-scale electronic computer with the aid of numerical analysis techniques.

The particular case of a land/sea boundary in a smooth, spherical surface is illustrated for a variety of cases by evaluating the convolution integrals on a large-scale computer.

#### l. Introduction.

The theory of propagation of electromagnetic waves around a sphere usually assumes the earth to be smooth and homogeneous. However, some extensions of the theory have been developed by Monteath [1951] and Hufford [1952]. More recently, Bremmer [1954] developed the plane earth theory for abrupt conductivity changes and the procedure was established in detail for plane and spherical earth theory by Wait [1956, 1957, 1961]. The formulation of Wait [1961] also developed procedures for the cases of transmitting and receiving terminals at very great heights. Furutsu [1957] has also developed methods for evaluating the field over sections of inhomogeneous ground and in particular has developed a method especially convenient for electronic computer computations in which the length of the sections of inhomogeneous ground are large. The formulation of this problem by the above mentioned authors, especially Wait [1961], comprised convolution integrals involving an integrand which can be calculated with the Watson [1918] type series of residues. This paper demonstrates a method of evaluating these integrals by an application of numerical analysis technique. Detailed numerical calculations for the particular case of the land/ sea boundary are presented.

#### 2. Theory.

The analysis of the propagation over a finite surface can be more conveniently formulated from the smooth, homogeneous earth theory by utilizing the attenuation factor, F, directly instead of the field (vertical electric field,  $E_r$ , for example), where

$$E_{r} = 2 E_{pr} F$$
 (1)

and the attenuation factor, F, is defined as the ratio of the total Hertz vector,  $\Pi$  to the primary Hertz vector  $\Pi$ ,

$$F = F(\omega, d) = \Pi / 2\Pi pr$$
(2)

where the radiation field,

$$E_{pr} = I_{o} \ell 10^{-7} \omega \exp(i\omega t - ik_{l} d)/d$$
(3)

$$\Pi_{\rm pr} = \exp\left(i\omega t - ik_1 d\right) / (-ik_1 d) \qquad (4)$$

The factor, F, can be evaluated with the aid of the series of residues, [Bremmer, 1949], and in particular can be written, [Johler, Kellar and Walters, 1956],

$$F = \left[2\pi \alpha^{\frac{2}{3}}(k_1 a)^{\frac{1}{3}} d/a\right] \sum_{s=0}^{\frac{1}{2}} \sum_{s=0}^{\infty} \frac{f_s(h_1) f_s(h_2)}{[2\tau_s - 1/\delta^2]}$$

$$\times \exp\left\{-i\left[\left(k_{1}a\right)^{\frac{1}{3}}\tau_{s}\alpha^{\frac{2}{3}}\frac{d}{a}+\frac{\alpha d}{2a}+\frac{\pi}{4}\right]\right\},\qquad(5)$$

s = 0, 1, 2, 3 ...,

where t = time, seconds, a is the radius of the spherical earth along which the vertical electric field,  $E_r$ , is directed,  $k_1$  is the wave number of air at the surface of the earth,  $k_1 = \omega \eta_1/c$  and  $\eta_1$  is the index of refraction of air,  $\eta_1 \sim 1$  (c = 2.997925<sup>(108)</sup> m/s,  $\eta_1 \sim 1.000338$  at the surface).  $\alpha$  is a factor which accounts for the vertical lapse of the index of refraction of air  $\alpha \sim 0.75$  to 0.85.  $h_1$  and  $h_2$  are the heights of the transmitter and receiver respectively and  $f_s(h)$ ,  $h = h_1$ ,  $h_2$ , are the height gain factor (identical for either receiver or transmitter as a result of reciprocity). The conductivity dielectric - constant parameter,  $\delta = \delta_e$  or  $\delta = \delta_m$  for vertical electric or vertical magnetic polarization of the Hertz dipole, respectively, can be written [Bremmer, 1949],

$$\delta_{e} = \frac{-i \frac{k_{2}^{2}}{k_{1}^{2}} \alpha^{\frac{1}{3}}}{(k_{1}a)^{\frac{1}{3}} \left[\frac{k_{2}^{2}}{k_{1}^{2}} - 1\right]^{\frac{1}{2}}},$$
(6)

$$\delta_{\rm m} = \frac{k_1^2}{k_2^2} \quad \delta_{\rm e}.$$
(7)

The wave number of the ground can be written,

$$k_{2} = \frac{\omega}{c} \sqrt{\epsilon_{2} - i \frac{\sigma \mu_{o} c^{2}}{\omega}}$$
(8)

The residues of the series (5) are defined by the special roots,  $\tau_s$ , of the differential equation of Riccati which has been formulated [Johler, Walters and Lilley, 1959]

$$\frac{\mathrm{d}\delta}{\mathrm{d}\tau} - 2 \,\delta^2 \tau + 1 = 0. \tag{9}$$

The height gain factors,  $f_s(h) = f_s(h_1)$  or  $f_s(h_2)$  are determined as the ratio of two Hankel functions,  $H_{\nu}^{(2)}(z)$ , [Bremmer, 1949],

$$f_{s}(h) = \sqrt{\frac{a}{r}} \frac{H_{\nu_{s}^{+} + \frac{1}{2}}^{(2)}(k_{1}r)}{H_{\nu_{s}^{+} + \frac{1}{2}}^{(2)}(k_{1}a)}$$
(10)

where

$$r = a + h \tag{11}$$

and

$$v_{\rm s} \sim k_{\rm l}^{\rm a} + (k_{\rm l}^{\rm a}) \tau_{\rm s}$$
, (12)

which can be approximately written [Johler, Kellar and Walters, 1956] in terms of the tabulated Furry [1945] modified Hankel functions,  $h_{12}(z)$ 

$$f_{s}(h) = \frac{h_{2} \left\{ \left(\frac{1}{2}\right)^{\frac{2}{3}} \left[ \left(k_{1}a\right)^{\frac{2}{3}} - \frac{2ha^{\frac{1}{3}}}{a} - 2\tau_{s} \right] \right\}}{h_{2} \left\{ -\left(2\right)^{\frac{1}{3}} \tau_{s} \right\}}$$
(13)

provided  $h \ll a$ .

The special case of a boundary (such as a land/sea boundary) located a distance,  $d_1$ , from the transmitter located on land with conductivity,  $\sigma_1$  and dielectric constant  $\epsilon_{2,1}$ , or wave number of the ground,  $k_2$ , is considered. A second finite ground surface located around the receiver from the distance  $d_1$  to d or the same boundary located a distance,  $d_2$ , from the receiver such that  $d = d_1 + d_2$  with a conductivity  $\sigma_2$ , a dielectric constant  $\epsilon_{2,2}$ . The field of the ground wave at the receiver is then, for a plane earth [Bremmer, 1949], and for a spherical earth [Wait, 1958, 1961],

$$E_r = 2 E_{pr} F^{(1,2)},$$
 (14)

where

$$F^{(1, 2)} = F^{(1, 2)}(\omega, d),$$

$$F^{(1, 2)} = F^{(1)}(d) - \sqrt{\frac{ik_1d}{2\pi}} \left[ \frac{Z^{(2)} - Z^{(1)}}{Z_0} \right]_0^{d_2} \frac{F^{(1)}(d - \rho) F^{(2)}(\rho)_{d\rho}}{\sqrt{\rho(d - \rho)}} (15)$$

$$(d - d_2 > 0)$$

where,

$$Z_{0} = \sqrt{\frac{\mu_{0}}{\epsilon_{0}}} = c\mu_{0} \sim 376.7$$
(16)

ohms, the impedance of space,

$$k_{1} = \frac{\omega}{c} \eta_{1}, \qquad (17)$$

$$Z^{(1)} = Z_{0} \frac{k_{1}}{k_{2,1}} \sqrt{1 - \frac{k_{1}^{2}}{k_{2,1}^{2}}}, \qquad (18)$$

$$Z^{(2)} = Z_{0} \frac{k_{1}}{k_{2,2}} \sqrt{1 - \frac{k_{1}^{2}}{k_{2,2}^{2}}}, \qquad (19)$$

$$k_{2,1} = \frac{\omega}{c} \left[ \epsilon_{2,1} - i \frac{\sigma_{1} \mu_{0} c^{2}}{\omega} \right]^{\frac{1}{2}}, \qquad (19)$$

$$k_{2,2} = \frac{\omega}{c} \left[ \epsilon_{2,2} - i \frac{\sigma_{2} \mu_{0} c^{2}}{\omega} \right]^{\frac{1}{2}}, \qquad (20)$$

where the integer p, identifies the particular finite surface under consideration. Thus, the smooth, homogeneous, spherical earthattenuation functions,  $F^{(1)}(d) = F^{(1)}(\omega, d)$ ,  $F^{(1)}(d-\rho) = F^{(1)}(\omega, d-\rho)$ ,  $F^{(2)}(\rho) = F^{(2)}(\omega, \rho)$  are employed in the calculation of the attenuation functions  $F^{(1, 2)}(\omega, d)$ , for mixed propagation i.e., propagation over finite surfaces. The integral of the combination of the two attenuation functions  $F^{(1)}(d-\rho)/\sqrt{d-\rho}$  and  $F^{(2)}(\rho)/\sqrt{\rho}$  appearing inside the integral (15) is by convention called the convolution of these functions, i.e.,  $F_1^*F_2 = \int_{0}^{t} F_1(t-\tau) F_2(\tau) d\tau$ . An alternate form (of equation (15)), can be written,

$$F^{(1, 2)} = F^{(2)}(d) - \sqrt{\frac{ik_{1}d}{2\pi}} \left[ \frac{Z^{(1)} - Z^{(2)}}{Z_{0}} \right]_{0}^{d_{1}} \frac{F^{(1)}(\rho)F^{(2)}(d-\rho)}{\sqrt{\rho(d-\rho)}} d\rho$$

$$(21)$$

$$(d > d_{1})$$

In either case, equations (15, 21), the problem is clearly resolved as the summation of the particular series of residues associated with the Watson [1918] transformation for the F, with various arguments function,  $F^{(1)}(\rho)$ ,  $F^{(2)}(d)$ ,  $F^{(2)}(d-\rho)$ ; or  $F^{(2)}(\rho)$ ,  $F^{(1)}(d)$ ,  $F^{(1)}(d-\rho)$ , whichever is appropriate; and the evaluation of a finite integral between the limits 0 and d<sub>1</sub> or 0 and d<sub>2</sub>, whichever is appropriate. This is accomplished quite readily with a Gaussian quadrature, [Kopal, 1955]

$$\int_{0}^{d} f(\rho) d(\rho) = \sum_{m=1}^{M} W_{m} f(\rho_{m}) + \epsilon (M), \qquad (22)$$

where  $\epsilon$  (M) is an error term which can be made arbitrarily small by increasing M, and,

$$W_{m} = \frac{1}{2} d_{1,2} H_{m}, m = 1, 2, 3 \dots M,$$
 (23)

$$\rho_{\rm m} = \frac{1}{2} \, d_{1,2} \, x_{\rm m} + \frac{1}{2} \, d_{1,2}. \tag{24}$$

The  $x_m$ 's are the Gaussian abscissas, and M determines the number of values of the integrand to be used in the quadrature. The integral can be split into any number of finite integrals to get convergence,

$$\int_{0}^{d_{1,2}} f(\rho) d\rho = \int_{0}^{a_{1}} f(\rho) d\rho + \int_{0}^{a_{2}} f(\rho) d\rho + \dots + \int_{n}^{d_{1,2}} f(\rho) d\rho.$$
(25)

The quadrature problem is therefore the evaluation of a finite integral,

$$\int_{a}^{b} F(x) dx = \sum_{m=1}^{M} W_{m} F(X_{m}) + \epsilon (M).$$
 (26)

where,

$$X_{m} = \frac{1}{2} [(b-a) x_{m} + (b+a)].$$
 (27)

The Gaussian weights and abscissas can be determined from the following:

$$\int_{-1}^{1} f(x) dx = \sum_{m=1}^{M} H_{m} f(x_{m}),$$

$$W_{m} = \frac{1}{2} (b-a) H_{m}.$$
(28)
(29)

The x 's are the roots of the Legendre polynomials defined by,

$$\frac{d^{m}}{d^{x^{m}}} (x^{2} - 1)^{m} = 2^{m} m! P_{m}(x), \qquad (30)$$

$$P_{0}(x) = 1$$

$$P_{1}(x) = x$$

$$P_{2}(x) = \frac{3}{2} x^{2} - \frac{1}{2}$$

$$\dots n^{and},$$

$$(m + 1) P_{m+1}(x) + m P_{m-1}(x) = (2 m + 1) \times P_{m}(x). \qquad (31)$$

Upon determination of the roots, the weight coefficients,  $H_m$ , of the corresponding quadrature formulas are evaluated as follows:

$$H_{m} = \frac{2}{(1-x_{m}^{2}) \left[P'_{m}(x_{m})\right]^{2}}$$
(32)

Forty-eight Gaussian weights and abscissas were used in the functions graphed in this paper, convergence being obtained by increasing the number of intervals of integration. These weights and abscissas have been published by Kopal [1955], Davis and Rabinowitz [1956].

The singularity, equation (21) for example,

$$[\rho (d-\rho)]^{-\frac{1}{2}} = [d (d-d_1)]^{-\frac{1}{2}}, (d = d_1),$$

would apparently present some difficulties in numerical integration.

However, it should be noted that the Gaussian quadrature, equations (24, 25, 26) does not include the end points of the interval of integration. A loss of precision, however, could possibly occur, but this could be sidestepped by the use of multiple precision. In any case, this did not become a problem for the examples presented in this paper.

The Gaussian quadrature techniques do not utilize the values of the integrand (15, 21) at the end points of integration. However, close to zero the series of residues, equation (5), converges slowly and, hence, the calculation of the attenuation function, F, becomes more difficult with a considerable reduction in computation efficiency. Since the Gaussian points,  $x_m$ , tend to cluster, equation (24) tends to cluster in the vicinity of  $d_{1,2} = 0$ , it is desirable to utilize the plane earth approximation in this area. Indeed, quite high precision can be obtained if the computation is preferred both by the series of residues and the plane-earth theory such that a comparison is made in the transition region to ascertain the computation precision. The attenuation for the plane earth theory [Norton, 1937] is,

$$F = \sqrt{\frac{2}{d^2 + h^2}} \left\{ \sin^2 \tau \left( 1 + R_e \right) + \left( 1 - R_e \right) \left[ 1 - \frac{k_1^2}{k_1^2} + \frac{k_1^4}{k_1^4} \sin \tau \right] f(\rho) \right\}$$

$$-2\left[\frac{k_1}{k_2}\sqrt{1-\frac{k_1^2}{k_2^2}}\sin\tau\right]\cos\tau \right], \qquad (33)$$

where,

$$f(\rho) = 1 - i \sqrt{\pi \rho} \exp(-\rho) \operatorname{erfc}(i \sqrt{\rho}) . \qquad (34)$$

if,

$$\rho_{1} = \frac{ik_{1}D}{2} \frac{k_{1}^{2}}{k_{2}^{2}} \left[ 1 - \frac{k_{1}^{2}}{k_{1}^{2}} \sin \tau \right] , \qquad (35)$$

where

$$D = \sqrt{d^2 + h^2}$$
, (36)

then,

$$\rho = \rho_1 \left\{ 1 + \frac{h}{\frac{k_1}{k_2} \sqrt{1 - \frac{k_1^2}{k_2^2}}} \right\}$$
(37)

where  $h_1$  is the height of either the transmitter,  $h = h_1$  or the receiver,  $h = h_2$  and  $k_1$  and  $k_2$  are the wave numbers of the medium previously defined, equation (18, 19, 20), which can be specified as some finite surface,  $k_{2,p}$ . The factor  $R_e$  is the Fresnel reflection coefficient for vertical polarization,

$$R_{e} = R_{e}(\tau) = \left\{ k_{2}^{2} \cos \tau / k_{1}^{2} - \left[ k_{2}^{2} - \sin \tau \right]^{\frac{1}{2}} \right\} / \left\{ k_{2}^{2} \cos \tau / k_{1}^{2} + \left[ k_{2}^{2} / k_{1}^{2} - \sin^{2} \tau \right]^{\frac{1}{2}} \right\}$$
(38)

where  $\tau = \tan^{-1} (d/h)$ . Thus, the height gain previously discussed for the residue series, equation (10, 12), is built into the plane-earth theory for short distances. It is now possible, with the aid of the convolution integrals (15, 21), to elevate transmitting and/or receiving dipoles and evaluate the field over mixed ground transmissions. Indeed, in the limit whereupon the receiver or transmitter is elevated to several earth radii; equation (10), the antenna pattern over mixed ground, is described. Furthermore, the convolution integrals can be developed for three or more finite ground surfaces [Wait, 1961]. Thus, discontinuities located at distances,  $d_1$ ,  $d_1 + d_2$ from the transmitter or  $d = d_2 + d_3$  and  $d - d_3$  respectively from the receiver, where  $d = d_1 + d_2 + d_3$  can be written,

$$F^{(1, 2, 3)} = F^{(1)} - \sqrt{\frac{ik_1d}{2\pi}} \left[ \frac{z^{(2)} - z^{(3)}}{z_0} \right] \int_{d_1+d_2}^{d} \frac{F^{(1)}(\rho) F^{(3)}d-\rho}{\sqrt{\rho (d-\rho)}} d\rho$$

$$-\sqrt{\frac{ikd}{2\pi}} \left[\frac{z^{(2)} - z^{(1)}}{z_{o}}\right] \int_{d_{1}}^{d_{1}+d_{2}} \frac{F^{(1)}(\rho)}{\sqrt{\rho}} \left[\frac{F^{(2)}(d-\rho)}{\sqrt{d-\rho}}\right]$$

$$-\sqrt{\frac{ik_{1}d}{2\pi}} \int_{d_{1}+d_{2}}^{d} \frac{F^{(2)}(\eta-\rho)F^{(3)}(d-\eta)}{\sqrt{(\eta-\rho)(d-\eta)}} d_{\eta} d_{\rho}.$$
(39)

#### 3. Computations.

The results of computations to illustrate technique are presented in figures 1, 2, 3, 4. The amplitude, figures 1, 2 illustrate the recovery effect of a land/sea boundary noted by Millington [1949]. There is, of course, a corresponding phase recovery, figures 3, 4. These computations were carried out on the CDC-1604 computer.

#### 4. Conclusions.

The results of this investigation indicate that the problem of wave propagation over inhomogeneous or finite ground surfaces can be treated with comparative ease on a large scale computer with the aid of numerical analysis techniques applied to the convolution integrals and an adequate technique for summing the series of residues.

#### 5. Acknowledgments.

The series of residues and the ground wave computer program concept was developed by John D. Harper, Jr. and J. R. Johler of Division 85.10 (Applied Electromagnetic Theory Section) and Harper programed same for IBM-704 and CDC-1604 computers. The convolution integrals were programed for CDC-1604 by C. M. Lilley of the same group.

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### LIST OF FIGURES

		PAGE
Fig. l	Amplitude of the ground wave as a function of distance for a variety of land/sea boundary locations, $d_1$ , illustrating the recovery effect, $f = 1000 \text{ kc/s}$ .	15
Fig. 2	Amplitude of the ground wave as a function of distance for a variety of land/sea boundary locations, d <sub>1</sub> , illustrating the recovery effect, f = 3000 kc/s, noted by Millington.	16
Fig. 3	Phase of the ground wave as a function of distance for a variety of land/sea boundary locations, $d_1$ , illustrating the recovery effect, f = 1000 kc/s.	17
Fig. 4	Phase of the ground wave as a function of distance for a variety of land/sea boundary locations, $d_1$ , illustrating the recovery effect, f = 3000 kc/s.	18

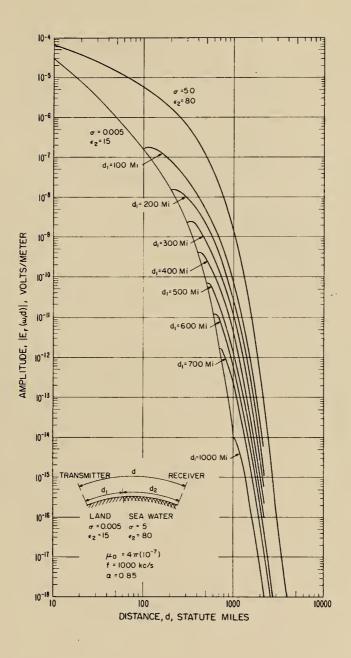


Figure 1. Amplitude of the ground wave as a function of distance for a variety of land/sea boundary locations, d<sub>1</sub>, illustrating the recovery effect, f = 1000 kc/s.

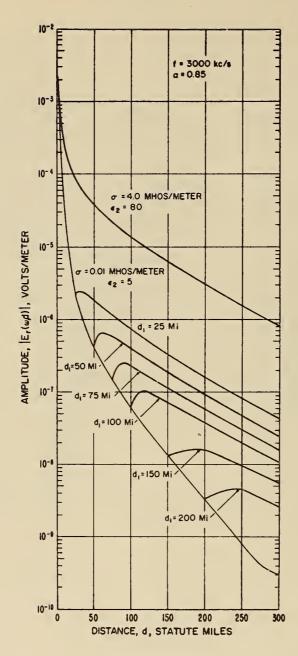


Figure 2. Amplitude of the ground wave as a function of distance for a variety of land/sea boundary locations,  $d_1$ , illustrating the recovery effect, f = 3000 kc/s, noted by Millington.

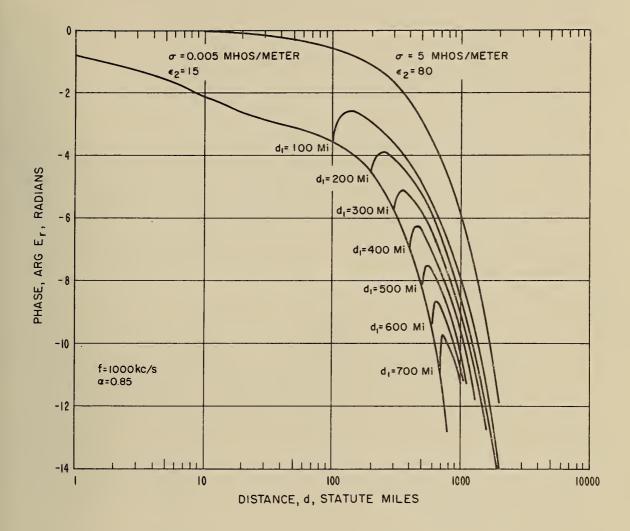


Figure 3. Phase of the ground wave as a function of distance for a variety of land/sea boundary locations, d<sub>1</sub>, illustrating the recovery effect, f = 1000 kc/s.

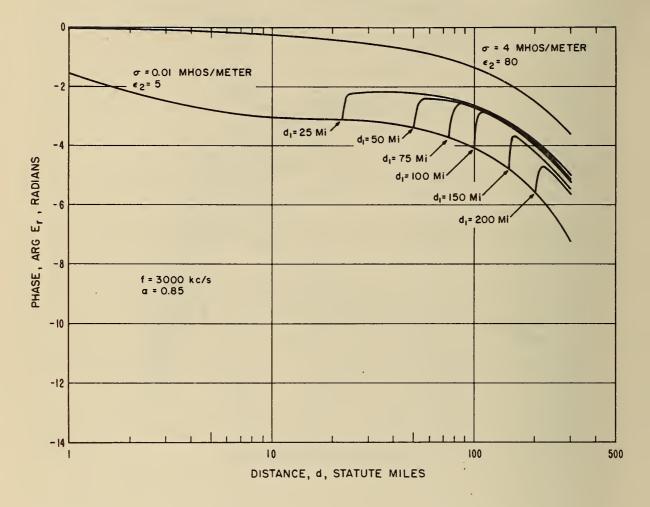


Figure 4. Phase of the ground wave as a function of distance for a variety of land/sea boundary locations,  $d_1$ , illustrating the recovery effect, f = 3000 kc/s.

#### 4. S. DEPARTMENT OF COMMERCE Luther H. Hodges, Secretary

NATIONAL BUREAU OF STANDARDS A. V. Astin, Director



#### THE NATIONAL BUREAU OF STANDARDS

The scope of activities of the National Bureau of Standards at its major laboratories in Washington, D.C., and Boulder, Colorado, is suggested in the following listing of the divisions and sections engaged in technical work. In general, each section carries out specialized research, development, and engineering in the field indicated by its title. A brief description of the activities, and of the resultant publications, appears on the inside of the

#### **WASHINGTON, D.C.**

Flectricity, Resistance and Reactance. Electrochemistry. Electrical Instruments. Magnetic Measurements. Dielectrics.

Metrology. Photometry and Colorimetry. Refractometry. Photographic Research. Length. Engineering Metrology. Mass and Scale. Volumetry and Densimetry.

Heat. Temperature Physics. Heat Measurements. Cryogenic Physics. Equation of State. Statistical Physics. Radiation Physics, X-ray, Radioactivity, Radiation Theory, High Energy Radiation, Radiological Equipment, Nucleonic Instrumentation, Neutron Physics.

Analytical and Inorganic Chemistry. Pure Substances. Spectrochemistry. Solution Chemistry. Standard Refer-ence Materials. Applied Analytical Research.

Mechanics. Sound. Pressure and Vacuum. Fluid Mechanics. Engineering Mechanics. Rheology. Combustion

Organic and Fibrous Materials. Rubber. Textiles. Paper. Leather. Testing and Specifications. Polymer Structure. Plastics. Dental Research.

Metallurgy. Thermal Metallurgy. Chemical Metallurgy. Mechanical Metallurgy. Corrosion. Metal Physics. Electrolysis and Metal Deposition.

Mineral Products. Engineering Ceramics. Glass. Refractories. Enameled Metals. Crystal Growth. Physical Properties. Constitution and Microstructure.

Building Research, Structural Engineering, Fire Research, Mechanical Systems, Organic Building Materials, Codes and Safety Standards, Heat Transfer, Inorganic Building Materials.

Applied Mathematics, Numerical Analysis, Computation, Statistical Engineering, Mathematical Physics, Op-

Data Processing Systems. Components and Techniques. Digital Circuitry. Digital Systems. Analog Systems. Applications Engineering.

Atomic Physics. Spectroscopy. Infrared Spectroscopy. Solid State Physics. Electron Physics. Atomic Physics. Instrumentation. Engineering Electronics. Electron Devices. Electronic Instrumentation. Mechanical Instruments. Basic Instrumentation.

Physical Chemistry. Thermochemistry. Surface Chemistry. Organic Chemistry. Molecular Spectroscopy. Molecular kinetics. Mass Spectrometry.

Office of Weights and Measures.

#### **BOULDER, COLO.**

Cryogenic Engineering. Cryogenic Equipment. Cryogenic Processes. Properties of Materials. Cryogenic Technical Services.

lonosphere Research and Propagation. Low Frequency and Very Low Frequency Research. Ionosphere Research. Prediction Services. Sun-Earth Relationships. Field Engineering. Radio Warning Services.

Radio Propagation Engineering. Data Reduction Instrumentation. Radio Noise. Tropospheric Measurements. Tropospheric Analysis. Propagation-Terrain Effects. Radio-Meteorology. Lower Atmosphere Physics. Radio Standards, High Frequency Electrical Standards. Radio Broadcast Service. Radio and Microwave Materi-als. Atomic Frequency and Time Interval Standards. Electronic Calibration Center. Millimeter-Wave Research.

Vicrowave Circuit Standards.

Radio Systems. High Frequency and Very High Frequency Research. Modulation Research. Antenna Research. Navigation Systems.

Upper Atmosphere and Space Physics. Upper Atmosphere and Plasma Physics. Ionosphere and Exosphere Scatter. Airglow and Aurora. Ionospheric Radio Astronomy.

