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# Computer Software for the Computation of the Scattered Field and the Optical Microscope Image of Line Objects Patterned in Thick Layers

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December 1987

Prepared for:

**U.S. DEPARTMENT OF COMMERCE**  
**National Bureau of Standards**  
**National Engineering Laboratory**  
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COMPUTATION OF THE SCATTERED FIELD  
AND THE OPTICAL MICROSCOPE IMAGE OF  
LINE OBJECTS PATTERNED IN THICK  
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U.S. DEPARTMENT OF COMMERCE, C. William Verity, *Secretary*  
NATIONAL BUREAU OF STANDARDS, Ernest Ambler, *Director*



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the Scattered Field and the Optical Microscope Image of  
Line Objects Patterned in Thick Layers

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**ABSTRACT:** This report contains computer software for calculating optical microscope images of line objects patterned in thick layers ( $>\lambda/4$  thick). The algorithms used are based on a monochromatic, waveguide model which can predict the images of line objects with arbitrary edge geometry including multilayer structures with sloped, curved, asymmetric, and undercut edges. Along with the computer software listing, the mathematics of the model, a short description of its structure and use, and test cases for help in implementation are given.

**KEY WORDS:** computer software; diffraction; dimensional metrology; linewidth; microscopy; optical imaging; optical metrology.

#### INTRODUCTION

The computer software described in this report was written in conjunction with the NBS project to develop fundamentally accurate optical measurement techniques for the width of micrometer and submicrometer lines patterned on integrated circuit wafers. Accurate and precise measurement techniques for linewidth are needed to improve yield, to ensure that lithographic and critical dimension (CD) measurement systems meet specification, to establish control of fabrication processes, and as input to device modeling and simulation programs.

In the course of research and development of the NBS laser linewidth measurement system [1], it was found that the microscope

image profiles for lines patterned in thick layers ( $>\lambda/4$  thick) were not properly predicted by the scalar theory conventionally used to calculate images of line objects [2, 3]. Scalar imaging theory assumes that the object is planar (thin compared to the wavelength of the illumination) and can be represented by a complex transmittance or reflectance function. The scalar model does not take into account the multiple reflections or standing waves that may occur within the patterned layer or the edge effects that occur in thick layers.

A new model was developed by Nyysönen [4, 5] based on a waveguide approach which characterizes the patterned layer by its complex dielectric constant and calculates the waveguide modes supported by the line structure. First, the mode expansion of the electromagnetic fields within the patterned layer is found and then the appropriate boundary conditions are used to calculate the scattered field.

Nyysönen's model assumes that the line structure is patterned in a nonmagnetic layer which can be characterized by its complex index of refraction which is taken to be constant with depth within the layer. Thus, this model can be used to represent homogeneous line structures with vertical edge walls. The spatial function representing the variation in the dielectric constant (square of the complex index of refraction) in the layer is expanded in a Fourier series. This series is substituted into the wave equation and the eigenvalue solutions to this equation which represent the waveguide modes are found. Assuming a single incident plane wave normal to the surface, the boundary value problem at the layer interfaces is solved to determine the Fourier coefficients in the expansions for the transmitted and reflected fields. This method allows the use of conventional scalar imaging theory to compute the image when no polarization effects are present. In such a case, the E- and H-field components are equivalent and either may be used in the scalar imaging equations.



To include the effects of partial coherence in the imaging system, it is possible to integrate over a finite illumination cone angle as long as the variation of the scattered field coefficients is negligible for the cone angle used. Hence, for layers approximately a few micrometers or less thick, finite coherence may be included.

This report covers the extension of this model by Kirk and Nyyssonen [6] to line structures of arbitrary edge geometry whose index of refraction may vary with depth in the layer. This line structure is approximated by subdividing it into a set of sublayers, each consisting of a line object with vertical edges and having a constant index of refraction over its small interval of depth. Each sublayer is treated in the same way as the single layer of Nyyssonen's earlier model. The waveguide modes and solution for the electromagnetic fields are found for each sublayer. There are now  $n+1$  boundary value equations where  $n$  is the number of sublayers. These equations when solved allow for the substitution of a single "equivalent" planar scattering layer for the multilayer structure. Thereafter, the solution for the transmitted and reflected field components and image are found in the same manner as for the single layer case. This model therefore allows for the modeling of line structures which contain different materials, have curved edges, asymmetry, and unpatterned (or patterned) sublayers. The mathematical details of the model are given in the expanded version of ref. [6] reproduced in Appendix III. A list of definitions of symbols used in the software and Appendix III is given in the attached table.

### Software Structure

As input to the software, the line structure may be characterized either by giving the index of refraction, thickness, and edge locations for the individual sublayers, or by a single refractive index, total thickness, and the polynomial coefficients used to

describe the edge geometry (See Appendix III.) In the latter case, the number of layers used to approximate this structure in the calculations must be designated. A good rule of thumb to use in determining the number of layers is to divide the change in linewidth (between the top and bottom interfaces) by  $0.1 \mu\text{m}$  and use a value larger than this. For most cases of interest in integrated circuit processing, 9 layers are sufficient. The program arrays, as written, allow up to 20 layers. The program calculates both the scattered field and the microscope image for TE-mode illumination. That is, the direction of polarization (i.e. E-field) is assumed to be parallel to the length of the line structure.

The scattered field is given in terms of the Fourier coefficients of the resulting far field diffraction pattern. In the test cases shown, 45 coefficients are given corresponding to  $\pm 22$  diffracted orders at diffraction angles corresponding to  $n\lambda/P$  where  $n$  is the order number,  $\lambda$  the wavelength, and  $P$  the period. The line structure is assumed to be repeated at a period  $P$  for convenience in calculating the Fourier series coefficients. An isolated line image is calculated by taking  $P$  large enough so that there is no influence from adjacent edges. Near resonances, this is not always possible due to time or storage limits on array sizes. The largest period used has been  $12 \mu\text{m}$  corresponding to a  $90 \times 90$  complex matrix for the boundary condition equations.

The microscopic image is calculated using conventional scalar imaging theory with the complex Fourier coefficients of the scattered field substituted for those of the planar object traditionally used. The software, as written, assumes a 1-D imaging system with a finite illumination cone angle. The Fourier coefficients of the scattered field are assumed to be constant over this angle. The calculations are accurate for a 2-D system as long as the condenser numerical aperture (N.A.) is much smaller

than that of the imaging objective. Hence, the calculations will accurately predict the image waveforms for the NBS laser system which has an objective N.A. of 0.85 or 0.95 and an illumination aperture 1/5 that of the objective at a wavelength of 514 nm.

The software also allows the calculation of the image for varying focus positions. Best focus is the top surface of the patterned layer. For all other focus positions, a quadratic phase factor is introduced (just as in scalar image theory).

The input for focus position is given as the number of wavelengths of defocus  $m$  where  $m$  is a dimensionless quantity. The relationship between the displacement of the measurement plane from the Gaussian focal plane,  $z$ , is given by

$$z = \frac{2}{\tan^2\theta} m$$

where  $\lambda$  is the wavelength and  $\tan\theta$  is substituted for N.A. ( $=\sin\theta$ ) for the high numerical apertures conventionally used for linewidth measurement.

### Implementation and Testing

This software (See Appendix I) is written in ANSI FORTRAN 77 consisting of a main program and 9 modules. The software is portable except for the complex matrix subroutines used for calculation of the eigenvectors and eigenvalues and matrix inversions. The software requires highly accurate routines because of the cascading of matrix multiplications and inversions. Those used at NBS on the Cyber 855\* are from the NAG library\* [7]. If these

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\*Certain commercial equipment, instruments, or materials are identified in this paper in order to adequately specify experimental or computational procedure. Such identification does not imply recommendation or endorsement by the National Bureau of Standards, nor does it imply that the materials or equipment identified are necessarily the best available for the purpose.

routines are unavailable, the user may substitute others as noted in the software. However, several test conditions should be met. A symmetric line object should produce symmetric Fourier coefficients to at least six significant figures. Also, a single-layer, thick line object with complex index of refraction should produce the same image structure as for a single layer when divided into 20 sublayers (or the maximum number of sublayers to be used). Repeatability to three or four significant figures in the image is sufficient to ensure this. If the matrix routines are inaccurate, the errors will increase as the number of layers is increased. Test cases are given in Appendix II. These six test cases are designed to test different portions of the software. Successful execution of one test case does not guarantee accuracy for the others.

All of the data calculated are output to files for use with available graphics. Plotting routines are not included here because they tend to be system dependent and user requirements may vary.

#### Accuracy

Testing of this software poses a challenging problem in that the only cases for which answers are known do not fully test the algorithms. For example, in comparing this model with scalar theory, agreement can be expected only in the limit as the thickness of the patterned layer approaches zero. Test case #2 for a patterned 0.09  $\mu\text{m}$  thick chromium layer is one such case which has been shown to agree with scalar theory. (See ref. 5.) Comparison has also been done for patterned thin layers of silicon dioxide showing excellent agreement. (See ref. 4.) In addition, it is required that the ratio of reflectivities of the patterned layer and substrate (or sublayers) away from the edges agrees with the values calculated from the Fresnel equations. (See ref. 8.) Otherwise, calculations can at this time only be compared with experimental measurements. This has been done for some cases

[6, 9]. However, although the agreement is good in most cases, the accuracy is suspect near resonances (where the line dimensions of width and/or thickness are equal to an integer multiple of the wavelength) and for small line dimensions. In both cases, the energy in the high spatial frequencies increases and errors in the eigenvalues and eigenvectors are expected to increase due to truncation errors.

#### Acknowledgments

A number of people have contributed to this software development. The original software was written for the Univac 1108 computer\* by Chris Kirk while a research associate at NBS. The author wishes to especially thank Ruth Varner, NBS, for her assistance in converting this software for use on the Cyber 855.\* The present software has been modified by the author to work for patterned metal layers. Other modifications to simplify its use and provide additional documentation have also been made.

#### References

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\*See previous footnote.

5. D. Nyysönen, "Optical Linewidth Measurement on Patterned Metal Layers," Proc. SPIE Vol. 480, Integrated Circuit Metrology II, 65-70 (1984).
6. D. Nyysönen and C. Kirk, "Modeling the Optical Microscope Images of Thick Layers for the Purpose of Linewidth Measurement," Proc. SPIE Vol. 538, Optical Microlithography IV, 179-187 (1985).
7. These routines are available from: Numerical Algorithms Group (NAG) Ltd., 7 Banbury Road, Oxford, Oxfordshire OX2 6NN, England.
8. D. Nyysönen, "Narrow-Angle Laser Scanning Microscope System for Linewidth Measurement on Wafers," NBSIR (to be published).
9. C. Kirk, "A Study of the Instrumental Errors in Linewidth and Registration Measurements Made with an Optical Microscope," Presented at the SPIE 1987 Santa Clara Symposium on Microlithography, Vol. 775, Integrated Circuit Metrology, Inspection and Process Control.

Table: Definition of Symbols Used in Software

	Math <u>Symbol</u>	Software <u>Label</u>
Free space wavelength in microns	$\lambda$	WAVE
Total layer thickness	T	TL
The number of layers	N	NS
Period of the grating	P	PER
RI of the air layer	$\hat{n}_0$	CRI(0)
RI of the substrate	$\hat{n}_s$	CRI(21)
RI of the n-th layer	$\hat{n}_n$	CRI(N)
Width polynomial coefficients	$X_j$	W(J)
Starting points for 1st, 2nd, and 3rd order	$Z_j$	PO(J)
The X position offset for each layer	--	XP(J)
The Z locations of the layer interfaces	$Z_n$	ZP (N)
Spatial frequency of the grating	1/P	SPAFRE
Free space wavenumber (in units of $\mu\text{m}^{-1}$ )	$k_0$	RK0
Fourier coefficients of dielectric constant	$E_{q,n}$	EQ(I)
Eigenvalue matrix	$D_{i,j}$	D(I,J)
Eigenvalues of the layers	$\alpha_{m,n}$	VAL(M)
Eigenvectors of the layers	$B_{j,m,n}$	VEC
Boundary equation matrices	$D_{j,m}^{11}, D_{j,m}^{12}, D_{j,m}^{21}, D_{j,m}^{22}$	RM(L,M)
(A) coefficients	$A_{m,1}, A'_{m,1}$	AV(M)
Fourier coefficients of the pseudo- object	$E_j^R$	FC(J)
1-D image of the line	--	ELEC(I)

Appendix I  
Computer Software Listing (THKIMAG)



1  
2 C  
3 C  
4 C  
5 C  
6 C  
7 C  
8 C  
9 C  
10 C  
11 C  
12 C  
13 C  
14 C  
15 C  
16 C  
17 C  
18 C  
19 C  
20 C  
21 C  
22 C  
23 C  
24 C  
25 C  
26 C  
27 C  
28 C  
29 C  
30 C  
31 C  
32 C  
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34 C  
35 C  
36 C  
37 C  
38 C  
39 C  
40 C  
41 C  
42 C  
43 C  
44 C  
45 C  
46 C  
47 C  
48 C  
49 C  
50 C

PROGRAM THKIMAG

THE ORIGINAL VERSION OF THIS PROGRAM WAS WRITTEN BY CHRIS KIRK, WHILE AN EMPLOYEE OF VICKERS INSTRUMENTS, (HAXBY ROAD, YORK, NORTH YORKSHIRE, ENGLAND) WHILE A RESEARCH ASSOCIATE AT THE NATIONAL BUREAU OF STANDARDS, GAITHERSBURG, MARYLAND, USA. SEPTEMBER 1984.

THE CURRENT VERSION HAS BEEN MODIFIED BY D. NYSSONEN TO WORK FOR METAL LAYERS. JUNE 1986

\*\*\*\*\*

THIS PROGRAM COMPUTES THE OPTICAL IMAGE OF LINE OBJECTS WITH ARBITRARY EDGE GEOMETRY, PATTERNED IN THICK LAYERS INCLUDING MULTILAYER STRUCTURES WITH SLOPING, CURVED, AND UNDERCUT EDGES, AS WELL AS ASYMMETRIC OBJECTS.

SEE REFERENCE: "MODELING OF THE OPTICAL IMAGING OF LINES PATTERNED IN THICK LAYERS WITH VARIABLE EDGE GEOMETRY," BY D. NYSSONEN AND C. P. KIRK

ALL EQUATION NUMBERS GIVEN REFER TO THIS MANUSCRIPT FOR QUESTIONS CONCERNING THIS PROGRAM CONTACT: R. D. LARRABEE, PRECISION ENGINEERING DIVISION, NATIONAL BUREAU OF STANDARDS OR D. NYSSONEN, CD METROLOGY, INC.

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SYSTEM DEPENDENT FEATURES.

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ALTHOUGH THIS SOFTWARE HAS BEEN WRITTEN IN FORTRAN 77 IN ORDER TO ALLOW FOR EASY PORTABILITY, THERE ARE A NUMBER OF FEATURES WHICH MAY BE SYSTEM DEPENDENT.

1. A FORTRAN 77 COMPILER IS REQUIRED.
2. THE SYSTEM MUST SUPPORT CERTAIN NAG LIBRARY ROUTINES.

A NOTE CONCERNING THE NAG LIBRARY.  
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THIS PROGRAM USES THE FOLLOWING NAG ROUTINES.

F02AKF            F04ADF

THESE ARE SINGLE PRECISION VERSIONS OF LIBRARY ROUTINES FOR THE CYBER WHICH MAY NOT BE SUPPORTED ON ALL MACHINES.

THESE ROUTINES ARE AVAILABLE FROM:

51 C NUMERICAL ALGORITHMS GROUP (NAG) LTD, 7 BANBURY ROAD,  
52 C OXFORD, OXFORDSHIRE. OX2 6NN. ENGLAND.

53 C  
54 C \*\*\*\*\*  
55 C PARAMETER (NLAY=20, NLAY1=21)  
56 C PARAMETER(NLIM=-22, LIM=22, KLIM=45, MLIM=90)  
57 C REAL PI, RK0, WAVE, SPAFRE, WI(NLAY), ZP(0:NLAY), NS, XP(NLAY), TRA(2)  
58 C COMPLEX D(NLIM:LIM, NLIM:LIM), RM(MLIM, MLIM), T(MLIM, MLIM)  
59 C \* , RN(KLIM, MLIM), AV(MLIM), FC(NLIM:LIM), CRI(0:NLAY1), Q0, VAL(KLIM)  
60 C PI = 4.0\*ATAN(1.0)  
61 C OPEN (UNIT=9, FILE='PARFIL')  
62 C OPEN (UNIT=10, FILE='FCOFDC')  
63 C OPEN (UNIT=11, FILE='DMATRX')  
64 C OPEN (UNIT=12, FILE='EIGVAL')  
65 C OPEN (UNIT=13, FILE='EIGVEC')  
66 C OPEN (UNIT=14, FILE='BOUNEQ')  
67 C OPEN (UNIT=15, FILE='BEMATX')  
68 C OPEN (UNIT=16, FILE='FCOFPO')  
69 C OPEN (UNIT=17, FILE='IMAGE')

70 C  
71 C \*\*\*\*\*  
72 C  
73 C SUMMARY OF THE INPUT/OUTPUT STRUCTURE.

74 C -----  
75 C  
76 C INPUT INPUT DATA FILE.  
77 C PARFIL UNIT9 PARAMETER FILE CREATED.  
78 C FCOFDC UNIT10 FOURIER COEFFICIENTS OF  
79 C DIELECTRIC CONSTANT CALCULATED  
80 C FOR FIRST LAYER.  
81 C DMATRX UNIT11 EIGENVALUE MATRIX D OF THE 1ST  
82 C LAYER.  
83 C EIGVAL UNIT12 EIGENVALUES OF THE LAYERS.  
84 C EIGVEC UNIT13 EIGENVECTORS OF THE LAYERS.  
85 C BOUNEQ UNIT14 BOUNDARY EQUATION MATRICES.  
86 C BEMATX UNIT15 (A) COEFFICIENTS.  
87 C FCOFPO UNIT16 FOURIER COEFFICIENTS OF  
88 C THE PSEUDO-OBJECT.  
89 C IMAGE UNIT17 1-D IMAGE OF THE LINE.

90 C  
91 C \*\*\*\*\*  
92 C  
93 C SET UP THE PARAMETERS FOR THE SECTIONED LINE.

94 C -----  
95 C  
96 C INPUT  
97 C WAVE = FREE SPACE WAVELENGTH IN MICRONS.  
98 C NS = THE NUMBER OF LAYERS.  
99 C PER = THE PERIOD.  
100 C CRI(0) = RI OF THE AIR LAYER.

```

101 C      CRI(21)= RI OF THE SUBSTRATE.
102 C      ID      = DEFINE LINE:  PARAMETRICALLY(0)
103 C      =      :  LAYER BY LAYER(1).
104 C      W(I)    = WIDTH POLYNOMIAL COEFFICIENTS.  I=ORDER
105 C      PO(I)    = STARTING POINTS FOR 1ST, 2ND AND 3RD ORDER.
106 C      CRI(1)   = REFRACTIVE INDEX OF PATTERNED LAYER.
107 C      TL      = LAYER THICKNESS.
108 C      TRA(1)  = PHOTOMULTIPLIER SLIT WIDTH.
109 C      TRA(2)  = TELEVISION CAMERA GAUSSIAN WIDTH PARAMETER.
110 C      LSKIP   = 0:  GENERATE A NEW FOURIER SERIES FOR
111 C      THE PSEUDO OBJECT.
112 C      1:  USE PREVIOUSLY GENERATED FOURIER SERIES.
113 C      -1:  GENERATE THE IMAGE.
114 C      DF = DEFOCUS IN WAVES WHERE DF IS A DIMENSIONLESS QUANTITY.
115 C      IF DEFOCUS IS KNOWN IN TERMS OF DISPLACEMENT OF THE
116 C      MEASUREMENT PLANE FROM THE FOCAL PLANE DZ, IT MUST BE
117 C      CONVERTED TO WAVES USING THE RELATIONSHIP:
118 C      DZ=(2*WAVE/(TAN(ANGLE))**2)*DF
119 C      WHERE ANGLE IS THE SAME AS DEFINED BY N.A. = SIN(ANGLE)
120 C      INPUT IF ID = 1.
121 C      WI(N)    = THE WIDTH OF THE LINES IN THE NTH LAYER.
122 C      ZP(N)    = THE Z LOCATIONS OF THE LAYER INTERFACES.
123 C      CRI(N)   = RI OF EACH LAYER.
124 C      XP(N)    = THE X POSITION OFFSET FOR EACH LAYER.
125 C
126 C      COMPUTATION
127 C      SPAFRE   = SPATIAL FREQUENCY OF THE GRATING (1/PERIOD).
128 C
129 C      COMPUTATIONS IF ID = 0.
130 C      CRI(N)   = RI OF EACH LAYER.
131 C      WI(N)    = THE WIDTHS OF THE LINES IN EACH LAYER.
132 C      ZP(N)    = THE Z LOCATIONS OF THE LAYER INTERFACES.
133 C      THE TOP SURFACE IS ASSUMED TO BE Z=0.
134 C      XP(N)    = THE X POSITION OFFSET FOR EACH LAYER (FOR
135 C      ASYMMETRIC CROSS-SECTIONS).
136 C
137 C      COMPUTATION
138 C      RK0 = FREE SPACE WAVENUMBER (IN UNITS OF 1/UM).
139 C
140 C      CALL LINE(PI,WAVE,RK0,CRI,ZP,WI,NS,SPAFRE,
141 C      *      TRA,XP,DF,LSKIP)
142 C
143 C      NOTE: CALCULATIONS ARE FOR LINE OBJECT, SPACES ARE
144 C      CALCULATED BY TAKING LINEWIDTH = PERIOD MINUS
145 C      DESIRED SPACE WIDTH.
146 C
147 C      *****
148 C
149 C      IF LSKIP =1 THEN GO STRAIGHT TO THE IMAGE SUBROUTINE.
150 C

```

```

151      IF (LSKIP.EQ.1) GO TO 20
152 C
153 C      *****
154 C
155 C      FIND THE EIGENVALUES AND EIGENVECTORS FOR EACH LAYER.
156 C
157 C      D IS THE EIGENVALUE MATRIX FOR EACH LAYER.
158 C      THE EIGENVALUES AND EIGENVECTORS ARE FOUND FOR EACH LAYER
159 C      AND THEN STORED ON DISK, STARTING WITH THE TOP LAYER.
160 C
161 C      THE DIFFRACTION SERIES IS TRUNCATED FOR DIFFRACTION
162 C      ANGLES WHICH EXCEED 90 DEGREES IN AIR. THIS IS SET
163 C      BY THE VARIABLE 'LIM'. LIM AND ITS RELATED VARIABLES
164 C      ARE SET IN THE PARAMETER STATEMENTS.
165 C
166 C      L = (REFRACTIVE INDEX OF AIR)*(PERIOD)/(WAVELENGTH)
167 C      NLIM = -L, LIM = L, KLIM = 2L+1, MLIM = 2(2L+1)
168 C
169 C      DO 10 N=1,NS
170 C          CALL SETUPD(PI,WI(N),CRI,D,SPAFRE,N,XP(N),WAVE,Q0,EQ2)
171 C          CALL EIGD(D,N,Q0,RK0,WAVE,SPAFRE,EQ2)
172 C      10 CONTINUE
173 C
174 C      *****
175 C
176 C      SET UP THE MATRIX ELEMENTS FOR THE FIRST INTERFACE
177 C      BOUNDARY CONDITIONS.
178 C      -----
179 C
180 C      THE ARRAY RM CONTAINS THE UPPER MATRIX ELEMENTS (SEE EQ. 21).
181 C
182 C      CALL SETRM(RM,RK0,WAVE,SPAFRE,CRI(0))
183 C
184 C      *****
185 C
186 C      SET UP THE PRODUCT MATRICES FOR THE N-LAYERS INTERFACES
187 C      -----
188 C
189 C      SET THE MATRIX T TO THE IDENTITY MATRIX.
190 C      THE T MATRIX IS MODIFIED BY ALL THE INTERMEDIATE INTERFACES.
191 C      STARTING WITH THE IDENTITY MATRIX, THE T MATRIX IS PASSED
192 C      TO THE SUBROUTINE FOR EACH INTERFACE AND RETURNED AFTER
193 C      MODIFICATION.
194 C
195 C      DO 40 L=1,MLIM
196 C          DO 30 M=1,MLIM
197 C              IF(L.EQ.M) THEN
198 C                  T(L,M) = CMPLX(1.0,0.0)
199 C              ELSE
200 C                  T(L,M) = CMPLX(0.0,0.0)

```

```

201           END IF
202 30        CONTINUE
203 40        CONTINUE
204           IF(NS.GT.1) THEN
205             DO 50 N=1,NS-1
206               CALL SETPQN(N,ZP(N),T)
207 50        CONTINUE
208           END IF
209 C
210 C *****
211 C
212 C SET UP THE MATRIX ELEMENTS FOR THE LAST INTERFACE BOUNDARY
213 C CONDITIONS.
214 C -----
215 C
216 C THE ARRAY RN CONTAINS THE LAST INTERFACE MATRIX
217 C ELEMENTS.
218 C
219 C CALL SETRN(RN,NS,ZP(NS),VAL)
220 C
221 C SOLVE THE MATRIX EQUATION FOR THE (A) COEFFICIENTS. (EQ. 21)
222 C
223 C THE VECTOR AV IS RETURNED CONTAINING THE (A) COEFFICIENTS.
224 C
225 C CALL SETAM(RM,T,RN,RK0,AV,VAL,NS,WAVE,SPAFRE,CRI(21))
226 C
227 C SET UP THE FOURIER COEFFICIENTS.
228 C -----
229 C
230 C THE FOURIER COEFFICIENTS ARE RETURNED IN VECTOR FC.
231 C T IS USED AS A WORKSPACE.
232 C
233 C CALL FCOEFF(AV,FC,T)
234 20        CONTINUE
235 C
236 C IF LSKIP = -1 THEN SKIP THE IMAGE ROUTINE.
237 C
238 C IF (LSKIP.EQ.-1) GOTO 60
239 C
240 C COMPUTE THE IMAGE FROM THE FOURIER COEFFICIENTS.
241 C -----
242 C
243 C CALL IMAGE(FC,PI,SPAFRE,TRA,DF,LSKIP,WAVE)
244 C
245 60        CONTINUE
246          STOP
247        END
248 C
249 C *****
250 C *****

```

```

251 C
252 C      SUBROUTINES TO SUPPORT THE MAIN PROGRAM.
253 C
254 C      *****
255 C
256 C      SUBROUTINE LINE(PI,WAVE,RK0,CRI,ZP,WI,NS,SPAFRE,
257 *      TRA,XP,DF,LSKIP)
258 C      COMPLEX CRI(0:21)
259 C      REAL PI,RK0,WAVE,SPAFRE,W(0:3),WI(20),PER,PO(3),
260 *      Z,ZP(0:20),NS,TL,XP(20),TRA(2)
261 C
262 C      SET THE ARRAYS TO ZERO.
263 C
264 C      DO 210 I=0,20
265 C          ZP(I) = 0.0
266 C          CRI(I) = CMPLX(0.0,0.0)
267 C          IF(I.GT.0) THEN
268 C              WI(I) = 0.0
269 C              XP(I) = 0.0
270 C          END IF
271 C      210 CONTINUE
272 C
273 C      READ IN THE WAVELENGTH, THE PERIOD, THE NUMBER OF LAYERS
274 C
275 C      READ*,WAVE,PER,NS
276 C      SPAFRE = 1.0/PER
277 C
278 C      READ IN THE RI'S OF THE AIR LAYER AND THE SUBSTRATE.
279 C      ID DETERMINES WHETHER THE LINE IS DEFINED BY PARAMETERS
280 C      OR LAYER BY LAYER.
281 C
282 C      READ*,CRI(0),CRI(21),ID
283 C
284 C      READ IN THE WIDTH POLYNOMIAL COEFFICIENTS.
285 C
286 C      READ*,W(0),W(1),W(2),W(3)
287 C
288 C      READ IN THE Z POSITION STARTING POINTS FOR THE 1ST,2ND
289 C      AND 3RD ORDERS.
290 C
291 C      READ*,PO(1),PO(2),PO(3)
292 C
293 C      READ IN THE REFRACTIVE INDEX OF THE PATTERNED LAYER
294 C      AND THE LAYER THICKNESS.
295 C
296 C      READ*,CRI(1),TL
297 C
298 C      READ IN THE PHOTOMULTIPLIER SLIT WIDTH AND THE VIDEO
299 C      CAMERA WIDTH PARAMETER.
300 C

```

```

301      READ*,TRA(1),TRA(2)
302      IF(TRA(1).GT.0.01) THEN
303          TRA(2) = 0.0
304      END IF
305 C
306 C      THE CAMERA WIDTH PARAMETER TRA(2) IS SET TO ZERO WHENEVER
307 C      THERE IS A FINITE SLIT WIDTH TRA(1).
308 C
309 C
310 C      READ IN THE SKIP PARAMETER (TO SKIP IMAGE CALCULATIONS)
311 C      AND THE AMOUNT OF DEFOCUS IN NUMBER OF WAVES.
312 C
313 C
314      READ*,LSKIP,DF
315 C
316      IF(ID.EQ.0) THEN
317 C
318 C          SET UP THE LINE DEFINED PARAMETRICALLY.
319 C
320          DO 220 I=1,NS
321              CRI(I) = CRI(1)
322              Z = TL*(I-0.5)/NS
323              WI(I) = W(3)*(Z-PO(3))**3+W(2)*(Z-PO(2))**2
324              WI(I) = WI(I)+W(1)*(Z-PO(1))+W(0)
325              ZP(I) = TL*I/NS
326 220      CONTINUE
327      ELSE IF(ID.EQ.1) THEN
328 C
329 C          SET UP THE LINE LAYER BY LAYER.
330 C
331 C          READ IN THE WIDTH, INTERFACE POSITION RI, AND X OFFSET
332 C          OF EACH LAYER
333 C
334          DO 230 I=1,NS
335              READ*,WI(I),ZP(I),CRI(I),XP(I)
336 230      CONTINUE
337      END IF
338 C
339 C      CALCULATE THE WAVE NUMBER.
340 C
341      RK0 = 2.0*PI/WAVE
342
343 C      SEND PARAMETERS TO A FILE.
344 C
345      WRITE(9,*) 'RUN PARAMETERS.'
346      WRITE(9,*) 'WAVELENGTH = ',WAVE
347      WRITE(9,*) 'NUMBER OF LAYERS=',NS
348      WRITE(9,*) 'WAVE NUMBER = ',RK0
349      WRITE(9,*) 'AIR LAYER = ',CRI(0)
350      WRITE(9,*) 'SUBSTRATE = ',CRI(21)

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351 WRITE(9,*) 'DEFOCUS =',DF
352 WRITE(9,*) 'SLIT WIDTH IN MICRONS =',TRA(1)
353 WRITE(9,*) 'CAMERA WIDTH PARAMETER =',TRA(2)
354 IF(LSKIP.LE.0) THEN
355 WRITE(9,*) 'LAYER      WIDTH      POSITION      RI      OFFSET'
356     DO 299 I=1,NS
357         WRITE(9,240) I,WI(I),ZP(I),CRI(I),XP(I)
358 240     FORMAT (1X,I2,6X,F6.3,6X,F5.3,4X,F5.3,1X,F5.3,
359 *           6X,F5.3)
360 299     CONTINUE
361     END IF
362     RETURN
363     END
364 C
365 C *****
366 C
367 SUBROUTINE SETUPD(PI,W,CRI,D,SPAFRE,N,XP,WAVE,Q0)
368 PARAMETER(NLIM=-22, LIM=22, KLIM=45, MLIM=90)
369 REAL PI,G05CAF,SPAFRE,RP,W,XP,RIP,WAVE
370 COMPLEX D(NLIM:LIM,NLIM:LIM),EQ(NLIM:LIM),E1,CI,Q0,
371 * CRI(0:21)
372 C
373 C W = LINEWIDTH OF LAYER N.
374 C XP = X OFFSET OF THE LINE IN LAYER N.
375 C
376 C THE LAYER FOURIER COEFFICIENTS FOR THE EXPANSION OF THE
377 C DIELECTRIC CONSTANT ARE STORED IN ARRAY EQ.
378 C CI = CMLPX(0.0,1.0)
379 C
380 C
381 C SET THE ARRAYS TO ZERO.
382 C
383 DO 320 I=NLIM,LIM
384     EQ(I) = CMLPX(0.0,0.0)
385     DO 310 J=NLIM,LIM
386         D(I,J) = CMLPX(0.0,0.0)
387 310     CONTINUE
388 320     CONTINUE
389 C
390 C ROUTINE TO GENERATE THE FOURIER COEFFICIENTS.
391 C THE OFFSET XP ADDS A PHASE TERM TO THE FOURIER
392 C COEFFICIENTS. SEE EQUATION 8 OF APPENDIX III.
393 C
394 E1 = CRI(N)**2 - CRI(0)**2
395 EQ(0) = CRI(0)**2+E1*W*SPAFRE
396 DO 330 I=NLIM,LIM
397     IF(I.EQ.0) GO TO 330
398     EQ(I) = E1*SIN(PI*I*SPAFRE*W)/PI/I
399     EQ(I) = EQ(I)*CEXP(2.0*PI*CI*SPAFRE*XP*I)
400     IF(N.EQ.1) THEN

```



```

401          WRITE(10,*) EQ(I),I
402          END IF
403 330      CONTINUE
404          Q0 = EQ(0)
405 C
406 C      ROUTINE TO CONSTRUCT THE D MATRIX.  SEE APPENDIX OF
407 C      APPENDIX III.
408          DO 350 L=NLIM,LIM
409              DO 340 M=NLIM,LIM
410                  IF(L.EQ.M) THEN
411                      D(L,M) = -(WAVE*L*SPAFRE)**2
412                  ELSE IF((IABS(L-M)).GT.22)THEN
413                      D(L,M) = CMPLX(0.0,0.0)
414                  ELSE
415                      D(L,M) = EQ(M-L)
416                  END IF
417 340      CONTINUE
418 350      CONTINUE
419          RETURN
420          END
421 C
422 C      *****
423 C
424          SUBROUTINE EIGD(D,N,Q0,RK0)
425          PARAMETER(NLIM=-22, LIM=22, KLIM=45, MLIM=90 )
426          COMPLEX D(KLIM,KLIM),VAL(KLIM),Q0
427          REAL AR(KLIM,KLIM),AI(KLIM,KLIM),RR(KLIM),
428 *          RI(KLIM),VR(KLIM,KLIM),VI(KLIM,KLIM),IWORK(KLIM),RK0
429 C
430 C      THE D MATRIX FOR THE FIRST LAYER IS WRITTEN ON TAPE11.
431 C      THE COMPLEX D MATRIX IS SPLIT INTO TWO REAL MATRICES
432 C      AND STORED IN ARRAYS AR() (REAL) AND AI() (IMAGINARY).
433 C
434          DO 420 I=1,KLIM
435              RR(I) = 0.0
436              RI(I) = 0.0
437              IWORK(I) = 0.0
438              DO 410 J=1,KLIM
439                  VR(I,J) = 0.0
440                  VI(I,J) = 0.0
441                  AR(I,J) = REAL(D(I,J))
442                  IF(ABS(AR(I,J)).LT.1.E-9) AR(I,J)=+0.0
443                  AI(I,J) =AIMAG(D(I,J))
444                  IF(ABS(AI(I,J)).LT.1.E-9) AI(I,J)=+0.0
445                  IF(N.EQ.1) THEN
446                      WRITE(11,*) D(I,J),I,J
447                  END IF
448 410      CONTINUE
449 420      CONTINUE
450 C

```

```

451 C      F02AKF SOLVES FOR THE EIGENVALUES AND EIGENVECTORS.
452 C      THE VECTORS ARE RETURNED IN VR AND VI AND THE EIGENVALUES
453 C      ARE RETURNED IN RR AND RI.
454 C      THE SQUARE ROOTS OF THE EIGENVALUES ARE WRITTEN ON
455 C      TAPE12 AND THE EIGENVECTORS ON TAPE13.
456 C
457       IFAIL = 0
458       CALL F02AKF(AR,KLIM,AI,KLIM,KLIM,RR,RI,VR,KLIM,VI,KLIM,
459 *          IWORK,IFAIL)
460 C
461 C      THE CYBER RESTRICTS PHASE ANGLES ON SQRT OF A COMPLEX
462 C      NO. TO  $-\pi/2 < \theta < \pi/2$  WHICH RESULTS IN INCORRECT
463 C      PHASE ANGLES. ANGLES ARE CORRECTED FOLLOWING
464 C      CALCULATION OF CSQRT.
465 C
466       DO 440 I=1,KLIM
467         VAL(I) = CMPLX(RR(I),RI(I))
468         VAL(I) = -VAL(I)-Q0
469         VAL(I) = CSQRT(VAL(I))*RK0
470         IF(AIMAG(VAL(I)).LT.0.0) THEN
471           VAL(I) = -VAL(I)
472         END IF
473         WRITE(12,*) VAL(I)
474         PRINT*,VAL(I),I
475         DO 430 J=1,KLIM
476           WRITE(13,*) VR(I,J),VI(I,J)
477 430      CONTINUE
478 440      CONTINUE
479       IF(IFAIL.EQ.0) THEN
480         PRINT*,'EIGENVALUE SUCCESS AT LAYER',N
481       ELSE
482         PRINT*,'EIGENVALUE FAILURE AT LAYER',N
483       END IF
484       RETURN
485       END
486 C
487 C      *****
488 C
489       SUBROUTINE SETRM(RM,RK0,WAVE,SPAFRE,CRO)
490       PARAMETER(NLIM=-22, LIM=22, KLIM=45, MLIM=90)
491       COMPLEX RM(MLIM,MLIM),VEC,VECCN,VAL(KLIM),CI,RKL,V,CRO
492       REAL VR,VI,RK0,WAVE,SPAFRE
493 C
494 C      CRO = REFRACTIVE INDEX OF THE AIR LAYER.
495 C
496 C      THE EIGENVALUE AND EIGENVECTORS ARE RECOVERED FOR THE
497 C      FIRST LAYER FROM UNIT12 AND UNIT13.
498 C
499       REWIND 12
500       REWIND 13

```

```

501      CI = CMPLX(0.0,-1.0)
502      DO 520 L=1,KLIM
503          V = CMPLX((WAVE*SPAFRE*(L-LIM-1))**2,0.0)
504          RKL = CSQRT(CRO**2-V)
505          DO 510 M=1,KLIM
506              IF(L.EQ.1) THEN
507                  READ(12,*) VAL(M)
508              END IF
509              READ(13,*) VR,VI
510              VEC = CMPLX(VR,VI)
511              RM(L,M) = (RKL-VAL(M)*CI/RK0)*VEC
512              RM(L,M+KLIM) = (RKL+VAL(M)*CI/RK0)*VEC
513              RM(L+KLIM,M) = CMPLX(0.0,0.0)
514              RM(L+KLIM,M+KLIM) = CMPLX(0.0,0.0)
515      510      CONTINUE
516      520      CONTINUE
517      RETURN
518      END
519 C
520 C      *****
521 C
522      SUBROUTINE SETPQN(N,ZP1,T)
523      PARAMETER(NLIM=-22, LIM=22, KLIM=45, MLIM=90)
524      COMPLEX PQ(MLIM,MLIM),T(MLIM,MLIM),RIPQ(MLIM,MLIM),
525 *      VEC,VECCN,VAL(KLIM)
526      REAL WKSPACE(MLIM),ZP1,VR,VI
527 C
528 C      ZP1 = THE Z LOCATION OF THE INTERFACE.
529 C
530      REWIND 12
531      REWIND 13
532 C
533 C      READ PAST THE EIGENVALUES AND EIGENVECTORS RELATING TO
534 C      THE PREVIOUS LAYERS.
535 C
536      IF(N.GT.1) THEN
537          DO 630 I=1,N-1
538              DO 620 J=1,KLIM
539                  READ(12,*) VAL(J)
540                  DO 610 K=1,KLIM
541                      READ(13,*) VR,VI
542      610          CONTINUE
543      620          CONTINUE
544      630          CONTINUE
545      END IF
546 C
547 C      SET UP THE UPPER LAYER MATRIX.
548 C
549      DO 650 L=1,KLIM
550          DO 640 M=1,KLIM

```

```

551         IF(L.EQ.1) THEN
552             READ(12,*) VAL(M)
553         END IF
554         READ(13,*) VR,VI
555         VEC = CMPLX(VR,VI)
556         PQ(L,M) = VEC*CEXP(-VAL(M)*ZP1)
557         PQ(L,M+KLIM) = VEC*CEXP(VAL(M)*ZP1)
558         PQ(L+KLIM,M) = VAL(M)*PQ(L,M)
559         PQ(L+KLIM,M+KLIM) = -VAL(M)*PQ(L,M+KLIM)
560 640     CONTINUE
561 650     CONTINUE
562 C
563 C     MULTIPLY THE PREVIOUS T MATRIX BY THE UPPER LAYER MATRIX.
564 C
565     DO 685 L=1,MLIM
566         DO 680 M=1,MLIM
567             RIPQ(L,M) = CMPLX(0.0,0.0)
568             DO 675 I=1,MLIM
569                 RIPQ(L,M) = RIPQ(L,M)+PQ(L,I)*T(I,M)
570 675         CONTINUE
571 680     CONTINUE
572 685     CONTINUE
573 C
574 C     WRITE THE MATRIX PRODUCT ONTO UNIT14.
575 C     CLEAR THE MATRICES RIPQ AND T.
576 C
577     REWIND 14
578     DO 658 L=1,MLIM
579         DO 654 M=1,MLIM
580             WRITE(14,*) RIPQ(L,M)
581             RIPQ(L,M) = CMPLX(0.0,0.0)
582             IF(L.EQ.M) THEN
583                 T(L,M) = CMPLX(1.0,0.0)
584             ELSE
585                 T(L,M) = CMPLX(0.0,0.0)
586             END IF
587 654     CONTINUE
588 658     CONTINUE
589 C
590 C     SET UP THE LOWER LAYER MATRIX.
591 C
592     DO 670 L=1,KLIM
593         DO 660 M=1,KLIM
594             IF(L.EQ.1) THEN
595                 READ(12,*) VAL(M)
596             END IF
597             READ(13,*) VR,VI
598             VEC = CMPLX(VR,VI)
599             PQ(L,M) = VEC*CEXP(-VAL(M)*ZP1)
600             PQ(L,M+KLIM) = VEC*CEXP(VAL(M)*ZP1)

```

```

601          PQ(L+KLIM,M) = VAL(M)*PQ(L,M)
602          PQ(L+KLIM,M+KLIM) = -VAL(M)*PQ(L,M+KLIM)
603 660      CONTINUE
604 670      CONTINUE
605 C
606 C      F04ADF INVERTS THE MATRIX PQ AND RETURNS THE INVERSE IN
607 C      THE MATRIX RIPQ.
608 C
609          IFAIL = 0
610          CALL F04ADF(PQ,MLIM,T,MLIM,MLIM,MLIM,RIPQ,MLIM,
611 *          WKSPACE,IFAIL)
612          IF(IFAIL.EQ.0) THEN
613              PRINT*,'SUCCESSFUL MATRIX INVERSION AT INTERFACE',N
614          ELSE
615              PRINT*,'SINGULAR MATRIX AT INTERFACE',N
616          STOP
617          END IF
618 C
619 C      RECOVER THE PREVIOUS MATRIX FROM UNIT14.
620 C
621          REWIND 14
622          DO 608 L=1,MLIM
623              DO 604 M=1,MLIM
624                  READ(14,*) PQ(L,M)
625 604          CONTINUE
626 608          CONTINUE
627 C
628 C      MULTIPLY THE TWO MATRICES AND RETURN THE PRODUCT IN T.
629 C
630          DO 694 L=1,MLIM
631              DO 694 M=1,MLIM
632                  T(L,M) = CMPLX(0.0,0.0)
633                  DO 690 I=1,MLIM
634                      T(L,M) = T(L,M)+RIPQ(L,I)*PQ(I,M)
635 690          CONTINUE
636 694          CONTINUE
637 698          CONTINUE
638          RETURN
639          END
640 C
641 C      *****
642 C
643          SUBROUTINE SETRN(RN,NS,ZS,VAL)
644          PARAMETER (NLIM=-22, LIM=22, KLIM=45, MLIM=90)
645          COMPLEX RN(KLIM,MLIM),VEC,VECCN,VAL(KLIM)
646          REAL VR,VI,ZS,NS
647 C
648 C      ZS = THE Z POSITION OF THE SUBSTRATE INTERFACE.
649 C      CRS = THE COMPLEX REFRACTIVE INDEX OF THE SUBSTRATE.
650 C

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```

651          REWIND 12
652          REWIND 13
653 C
654 C          SKIP OVER THE ALPHA(M) AND B(L,M) TERMS OF THE UPPER
655 C          LAYERS.
656 C
657          IF(NS.GT.1) THEN
658              DO 730 I=1,NS-1
659                  DO 720 J=1,KLIM
660                      READ(12,*) VAL(J)
661                      DO 710 K=1,KLIM
662                          READ(13,*) VR,VI
663                          710          CONTINUE
664                          720          CONTINUE
665                          730          CONTINUE
666          END IF
667 C
668 C          SET UP THE RN MATRIX.
669          DO 750 L=1,KLIM
670              DO 740 M=1,KLIM
671                  IF(L.EQ.1) THEN
672                      READ(12,*) VAL(M)
673                  END IF
674                  READ(13,*) VR,VI
675                  VEC = CMPLX(VR,VI)
676                  RN(L,M)= CEXP(-VAL(M)*ZS)*VEC
677                  RN(L,M+KLIM) = CEXP(VAL(M)*ZS)*VEC
678                  740          CONTINUE
679                  750          CONTINUE
680          RETURN
681          END
682 C
683 C          *****
684 C
685          SUBROUTINE SETAM(RM,T,RN,RK0,AV,VAL,NS,WAVE,SPAFRE,CRS)
686          PARAMETER(NLIM=-22, LIM=22, KLIM=45, MLIM=90)
687          COMPLEX RM(MLIM,MLIM),T(MLIM,MLIM),RN(KLIM,MLIM),RKL,
688          *          RV(MLIM),AV(MLIM),CI,VAL(KLIM),CRS,CRSUB,S2,S3,A1,ASUB
689          REAL RK0,WKSPCE(MLIM),WAVE,SPAFRE,NS
690 C
691 C          ASSEMBLE THE RM,RN AND T MATRICES INTO ONE MATRIX (RM).
692 C
693 C          THIS VERSION ALLOWS FOR THE INCLUSION OF AN UNPATTERNED
694 C          SUBLAYER (OXIDE). IF NOT WANTED, SET TSUB=0.0
695 C          TSUB IS THE THICKNESS IN UM OF THIS LAYER AND CRSUB IS
696 C          THE COMPLEX INDEX OF REFRACTION.
697 C
698          TSUB = 0.0
699          CRSUB = CMPLX(1.46,0.0)
700          CI = CMPLX(0.0,-1.0)

```

```

701      IF(NS.EQ.1) THEN
702      DO 820 L=1,KLIM
703          V = CMPLX((WAVE*SPAFRE*(L-LIM-1))**2,0.0)
704          S2 = CSQRT(CRSUB**2-V)
705          S3 = CSQRT(CRS**2-V)
706          A1 = CEXP(CMPLX(0.0,2.0)*RK0*TSUB*S2)
707          ASUB = ((S2+S3)-(S2-S3)*A1)/((S2+S3)+(S2-S3)*A1)
708          RKL = ASUB*S2
709          DO 810 M=1,KLIM
710              RM(L+KLIM,M) = RN(L,M)*(RKL+VAL(M)*CI/RK0)
711              RM(L+KLIM,M+KLIM) = RN(L,M+KLIM)*(RKL-VAL(M)*CI/RK0)
712      810      CONTINUE
713      820      CONTINUE
714      ELSE
715      DO 836 L=1,KLIM
716          V = CMPLX((WAVE*SPAFRE*(L-LIM-1))**2,0.0)
717          S2 = CSQRT(CRSUB**2-V)
718          S3 = CSQRT(CRS**2-V)
719          A1 = CEXP(CMPLX(0.0,2.0)*RK0*TSUB*S2)
720          ASUB = ((S2+S3)-(S2-S3)*A1)/((S2+S3)+(S2-S3)*A1)
721          RKL = ASUB*S2
722          DO 834 M=1,KLIM
723              DO 832 I=1,KLIM
724                  RM(L+KLIM,M) = RM(L+KLIM,M) + RKL*(RN(L,I)*T(I,M) +
725      *          RN(L,I+KLIM)*T(I+KLIM,M)) + (RN(L,I)*VAL(I)*CI/RK0*
726      *          T(I,M) - RN(L,I+KLIM)*VAL(I)*CI/RK0*T(I+KLIM,M))
727                  RM(L+KLIM,M+KLIM) = RM(L+KLIM,M+KLIM) + RKL*(RN(L,I)*
728      *          T(I,M+KLIM) + RN(L,I+KLIM)*T(I+KLIM,M+KLIM)) +
729      *          (RN(L,I)*VAL(I)*CI/RK0*T(I,M+KLIM) - RN(L,I+KLIM)*
730      *          VAL(I)*CI/RK0*T(I+KLIM,M+KLIM))
731      832      CONTINUE
732      834      CONTINUE
733      836      CONTINUE
734      END IF
735      C
736      C      SET UP THE RIGHT HAND SIDE VECTOR RV.
737      C
738      DO 840 L=1,MLIM
739          WKSPACE(L) = 0.0
740          RV(L) = CMPLX(0.0,0.0)
741      840      CONTINUE
742      RV(LIM+1)=CMPLX(2.0,0.0)
743      C
744      C      F04ADF SOLVES THE SET OF SIMULTANEOUS EQUATIONS DEFINED
745      C      BY RM AND RV.  SEE EQ. 21 OF APPENDIX III.
746      C      THE SOLUTIONS ARE RETURNED IN THE VECTOR AV.
747      C
748      IFAIL = 0
749      C
750      PRINT *, '4 SECTIONS OF RM MATRIX'

```

```

751 C      WRITE(1,*)'FIRST SECTION RM(1,M),M=1,45'
752      PRINT *, (RM(1,M),M=1,45)
753 C      WRITE(1,*)(RM(1,M),M=1,45)
754      PRINT *, (RM(1,M+45),M=1,45)
755 C      WRITE(1,*)'SECOND SECTION RM(1,M+45),M=1,45'
756      WRITE(1,*)(RM(1,M+45),M=1,45)
757      PRINT *, (RM(46,M),M=1,45)
758 C      WRITE(1,*)'THIRD SECTION RM(46,M),M=1,45'
759      WRITE(1,*)(RM(46,M),M=1,45)
760      PRINT *, (RM(46,M+45),M=1,45)
761 C      WRITE(1,*)'FOURTH SECTION RM(46,M+45),M=1,45'
762      WRITE(1,*)(RM(46,M+45),M=1,45)
763
764 C
765 C
766      IFAIL=0
767      CALL F04ADF(RM,MLIM,RV,MLIM,MLIM,1,AV,MLIM,
768 *      WKSPACE,IFAIL)
769      IF(IFAIL.EQ.0) THEN
770          PRINT*,'SUCCESSFUL SOLUTION FOR (A) COEFFICIENTS.'
771      ELSE
772          PRINT*,'FAILURE TO SOLVE FOR (A) COEFFICIENTS.'
773      END IF
774 C      PRINT *, (AV(L),L,'AV',L=1,MLIM)
775      RETURN
776      END
777 C
778 C      *****
779 C
780      SUBROUTINE FCOEFF(AV,FC,T)
781      PARAMETER (NLIM=-22, LIM=22, KLIM=45, MLIM=90 )
782      COMPLEX AV(MLIM),FC(KLIM),T(MLIM,MLIM)
783      REAL VR,VI
784      REWIND 13
785 C
786 C      CALCULATE THE FOURIER COEFFICIENTS OF THE PSEUDO-OBJECT
787 C      AND WRITE THE RESULTS ON UNIT15.
788 C      STORE THE FIRST LAYER EIGENVECTORS IN MATRIX T.
789 C
790      DO 920 L=1,KLIM
791          WRITE(15,*) AV(L),AV(L+KLIM),L
792          DO 910 M=1,MLIM
793              READ(13,*) VR,VI
794              T(L,M) = CMPLX(VR,VI)
795 910      CONTINUE
796 920      CONTINUE
797 C
798 C      CLEAR THE FC VECTOR AND SET THE FC(0) VALUE TO -1.0.
799 C      CALCULATE FOURIER SERIES COEFFICIENTS FC USING THE A-
800 C      COEFFICIENTS AND THE EIGENVECTORS. SEE EQ. 22 OF APPENDIX III.

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```

801 C
802 DO 940 L=1,KLIM
803 IF(L.EQ.LIM+1) THEN
804 FC(L) = CMPLX(-1.0,0.0)
805 ELSE
806 FC(L) = CMPLX(0.0,0.0)
807 END IF
808 DO 930 M=1,KLIM
809 FC(L) = FC(L)+(AV(M)+AV(M+KLIM))*T(L,M)
810 930 CONTINUE
811 940 CONTINUE
812 C
813 C THE FOURIER SERIES FOR THE PSEUDO-OBJECT IS WRITTEN ON
814 C UNIT16, IN A FORMAT WHICH ALLOWS IT TO INTERFACE TO A
815 C PLANAR IMAGING OR OTHER PROGRAM.
816 C
817 DO 950 K=1,KLIM
818 WRITE(16,999) FC(K)
819 PRINT*,FC(K)
820 950 CONTINUE
821 999 FORMAT(1X,2F20.8)
822 RETURN
823 END
824 C
825 C *****
826 C
827 SUBROUTINE IMAGE(FC,PI,SPAFRE,TRA,DF,LSKIP,WAVE)
828 PARAMETER(NLIM=-22, LIM=22, KLIM=45, MLIM=90)
829 REAL PI,WAVE,SPAFRE,RR,RI,XPOS,TRA(2),
830 2 RIM(-600:600),IR,CAM(-100:100),ELEC(-500:500),X,YN,XI
831 COMPLEX FC(NLIM:LIM),CI,CCI,RT,FFC,FOC(-KLIM:KLIM)
832 C
833 C THIS SUBROUTINE COMPUTES THE IMAGE OF THE LINE ASSUMING
834 C ONE DIMENSIONAL OPTICS.
835 C
836 C THE IMAGE IS COMPUTED FOR 1000 POINTS OVER A RANGE
837 C OF +/- 0.5*THE PERIOD.
838 C
839 CI = CMPLX(0.0,1.0)
840 C
841 C IF LSKIP = 1 THEN READ IN THE PREVIOUS COEFFICIENTS.
842 C
843 IF(LSKIP.EQ.1) THEN
844 DO 105 I=NLIM,LIM
845 READ(16,*) RR, RI
846 FC(I) = CMPLX(RR,RI)
847 105 CONTINUE
848 END IF
849 C
850 C OBJ = NUMERICAL APERTURE OF THE OBJECTIVE.

```

```

851 C      CON = NUMERICAL APERTURE OF THE "CONDENSER"
852 C      (ILLUMINATION APERTURE).
853 C      THESE PARAMETERS ARE PRESENTLY SET FOR THE NBS LASER SYSTEM
854 C
855      OBJ = 0.85
856      CON = 0.17
857      FM = OBJ*LIM
858      FR = CON*LIM
859 C
860 C      THE DEFOCUS TERM DF ADDS A PHASE MODULATION TO THE
861 C      FOURIER SERIES.
862 C
863      CCI = CI*2.0*PI*SPAFRE
864      DO 115 I=-KLIM,KLIM
865          FOC(I) = CEXP(2.0*PI*CI*DF*(WAVE*I*SPAFRE)**2)
866 115      CONTINUE
867 C
868 C      THE IMAGE IS COMPUTED FOR AN ADDITIONAL 100 POINTS ON
869 C      EITHER SIDE IN ORDER TO ALLOW FOR CONVOLUTION WITH A
870 C      DETECTOR APERTURE OR IMPULSE RESPONSE.
871 C
872      DO 130 IX=-600,600
873          XPOS = IX/(SPAFRE*1000.0)
874          RIM(IX) = 0.0
875          DO 120 J=-FR,FR
876              RT = CMPLX(0.0,0.0)
877              DO 110 K=J-FM,J+FM
878                  IF(ABS(K).GT.LIM) THEN
879                      FFC = CMPLX(0.0,0.0)
880                  ELSE
881                      FFC = FC(K)*FOC(K-J)
882                  END IF
883                  RT = RT+FFC*CEXP(CCI*XPOS*K)
884 110          CONTINUE
885              RIM(IX) = RIM(IX)+REAL(RT)**2+AIMAG(RT)**2
886 120          CONTINUE
887 130          CONTINUE
888 C
889 C      THE ARRAY RIM() CONTAINS THE OPTICAL IMAGE PROFILE.
890 C
891 C      CONVOLVE THE IMAGE WITH THE DETECTOR LINE SPREAD
892 C      FUNCTION.
893 C      IF TRA(1) IS GREATER THAN 0.01 THEN A SLIT APERTURE IS
894 C      ASSUMED.
895 C      OTHERWISE IF TRA(2) IS GREATER THAN 0.05 THEN A VIDEO
896 C      CAMERA WITH A GAUSSIAN LINE SPREAD FUNCTION IS ASSUMED.
897 C      IF BOTH ARE ZERO THEN NO CONVOLUTION TAKES PLACE.
898 C      THE ARRAY CAM() CONTAINS THE VIDEO CAMERA RESPONSE.
899 C      THE ARRAY ELEC() CONTAINS THE ELECTRICAL OUTPUT SIGNAL
900 C      OF THE TRANSDUER.

```

```

901 C
902     IF(TRA(1).GT.0.01) THEN
903         IR = TRA(1)*SPAFRE*500.0
904         IF(IR.GT.100) THEN
905             IR = 100
906         ELSE IF(IR.LT.1) THEN
907             IR = 1
908         END IF
909         DO 150 I=-500,500
910             ELEC(I) = 0.0
911             DO 140 J=-IR,IR
912                 ELEC(I) = ELEC(I)+RIM(J+I)
913     140         CONTINUE
914     150         CONTINUE
915     ELSE IF(TRA(2).GT.0.05) THEN
916         DO 160 I=-100,100
917             X = I/(SPAFRE*1000)
918             CAM(I) = EXP(-0.5*(X**2)/(TRA(2)**2))
919     160         CONTINUE
920         DO 180 I=-500,500
921             ELEC(I) = 0.0
922             DO 170 J=-100,100
923                 ELEC(I) = ELEC(I)+CAM(J)*RIM(I+J)
924     170         CONTINUE
925     180         CONTINUE
926     ELSE
927         DO 190 I=-500,500
928             ELEC(I) = RIM(I)
929     190         CONTINUE
930     END IF
931 C
932 C     WRITE THE OUTPUT IMAGE DATA ON UNIT17.
933 C
934 C     THE IMAGE IS NORMALISED WITH RESPECT TO THE INTENSITY AT
935 C     A DISTANCE OF 1/2 THE PERIOD AWAY FROM THE ORIGIN.
936 C
937     YN = AMAX1(ELEC(0),ELEC(500))
938     DO 195 I=-500,500
939         XI = I/(SPAFRE*1000.0)
940         ELEC(I) = ELEC(I)/YN
941         WRITE(17,*) ELEC(I),XI
942     195     CONTINUE
943 C
944 C     THE FOLLOWING 'END OF DATA' STATEMENT IS FOR THE
945 C     'DATAPLOT' GRAPHICS PACKAGE.
946 C
947     WRITE(17,*) 'END OF DATA'
948     RETURN
949     END

```

Appendix II  
Test Cases for Thick Layer Imaging Program (THKIMAG)

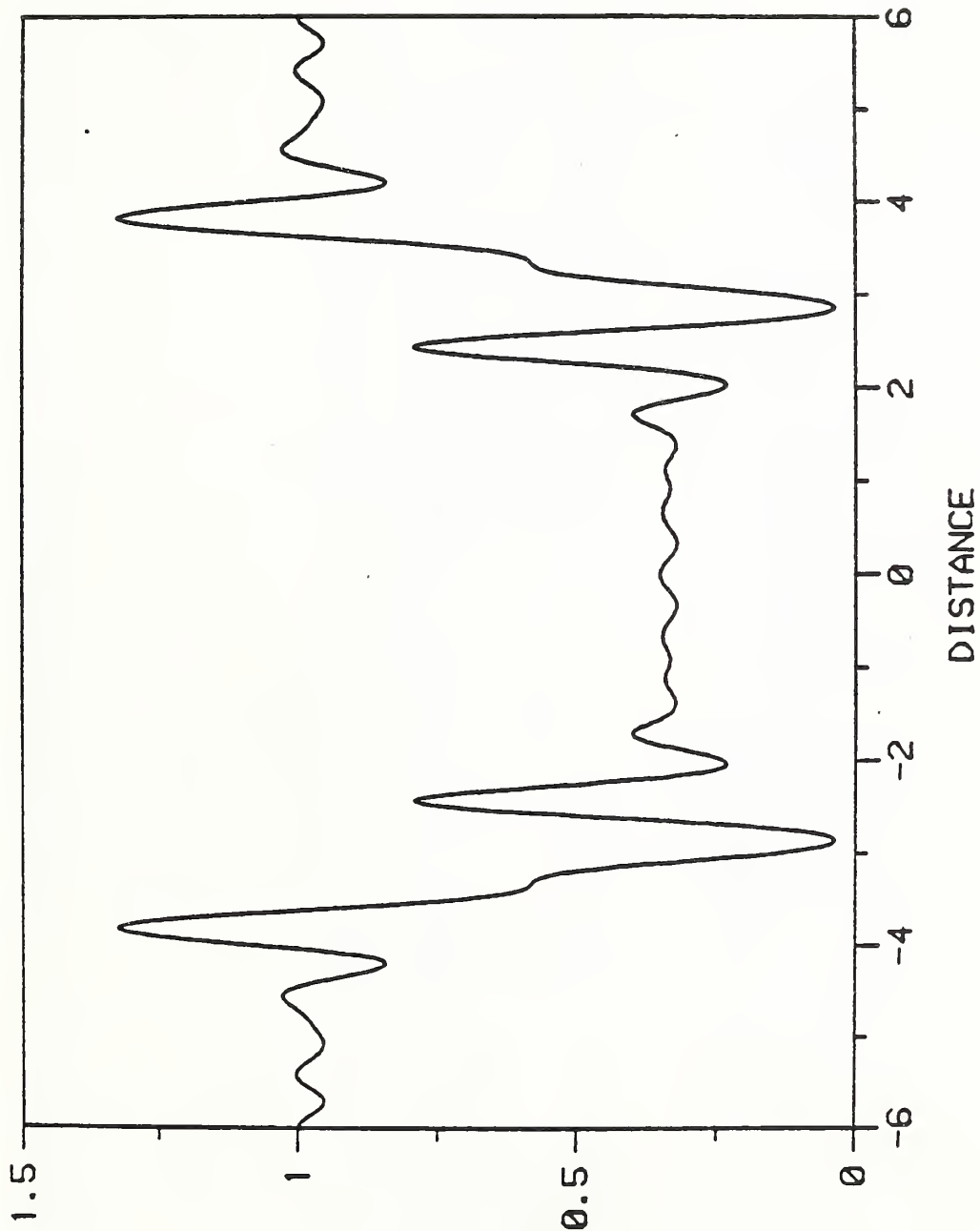
Test Case #1: THICK DIELECTRIC LAYER WITH VERTICAL EDGES  
 (6  $\mu\text{m}$  wide line with vertical edges patterned  
 in a 0.65  $\mu\text{m}$  thick silicon dioxide layer on a  
 silicon substrate).

1	0.53	12.0	1	:	WAVELENGTH, PERIOD, NUMBER OF LAYERS
2	(1.00,0.00)	(4.10,0.06)	0	:	RI(AIR), RI(SUBSTRATE), ID
3	6.00	0.00	0.00	:	WIDTH, 3 POLYNOMIAL COEFFICIENTS
4	0.00	0.00	0.00	:	DISPLACEMENTS IN Z
5	(1.46,0.0)	0.65		:	RI(PATTERNED LAYER), LAYER THICKNESS
6	0.20	0.00		:	SLIT WIDTH, CAMERA WIDTH
7	0	0.00		:	SKIP, DEFOCUS

Test Case #1: File FCOFPO: Fourier Coefficients of the  
Pseudo-Object

1	-.09108273	.00519622
2	-.06079249	.01308535
3	.00478018	.01658667
4	.04455374	.00100212
5	.00810522	-.00558479
6	-.03495132	-.01667204
7	-.00672853	.00298502
8	.01767570	.02462243
9	.00840075	-.00504291
10	-.00193394	-.02052349
11	-.01237166	.00493059
12	-.00605733	.00957640
13	.01508445	-.00253726
14	.00599687	.00264588
15	-.01573390	-.00030316
16	-.00008822	-.01386078
17	.01502888	.00264068
18	-.01048221	.02557344
19	-.01386187	-.00424937
20	.02992448	-.04556850
21	.01290886	.00491631
22	-.10926479	.13891567
23	.39447293	.04366360
24	-.10926479	.13891567
25	.01290886	.00491631
26	.02992448	-.04556850
27	-.01386187	-.00424937
28	-.01048221	.02557344
29	.01502888	.00264068
30	-.00008822	-.01386078
31	-.01573390	-.00030316
32	.00599687	.00264588
33	.01508445	-.00253726
34	-.00605733	.00957640
35	-.01237166	.00493059
36	-.00193394	-.02052349
37	.00840075	-.00504291
38	.01767570	.02462243
39	-.00672853	.00298502
40	-.03495132	-.01667204
41	.00810522	-.00558479
42	.04455374	.00100212
43	.00478018	.01658667
44	-.06079249	.01308535
45	-.09108273	.00519622

CASE1 PROFILE FOR DN THKIMAG



Test Case #2: THIN METAL LAYER WITH VERTICAL EDGES  
 (6  $\mu\text{m}$  wide line with vertical edges patterned  
 in a 0.09  $\mu\text{m}$  thick chromium layer on a silicon  
 substrate).

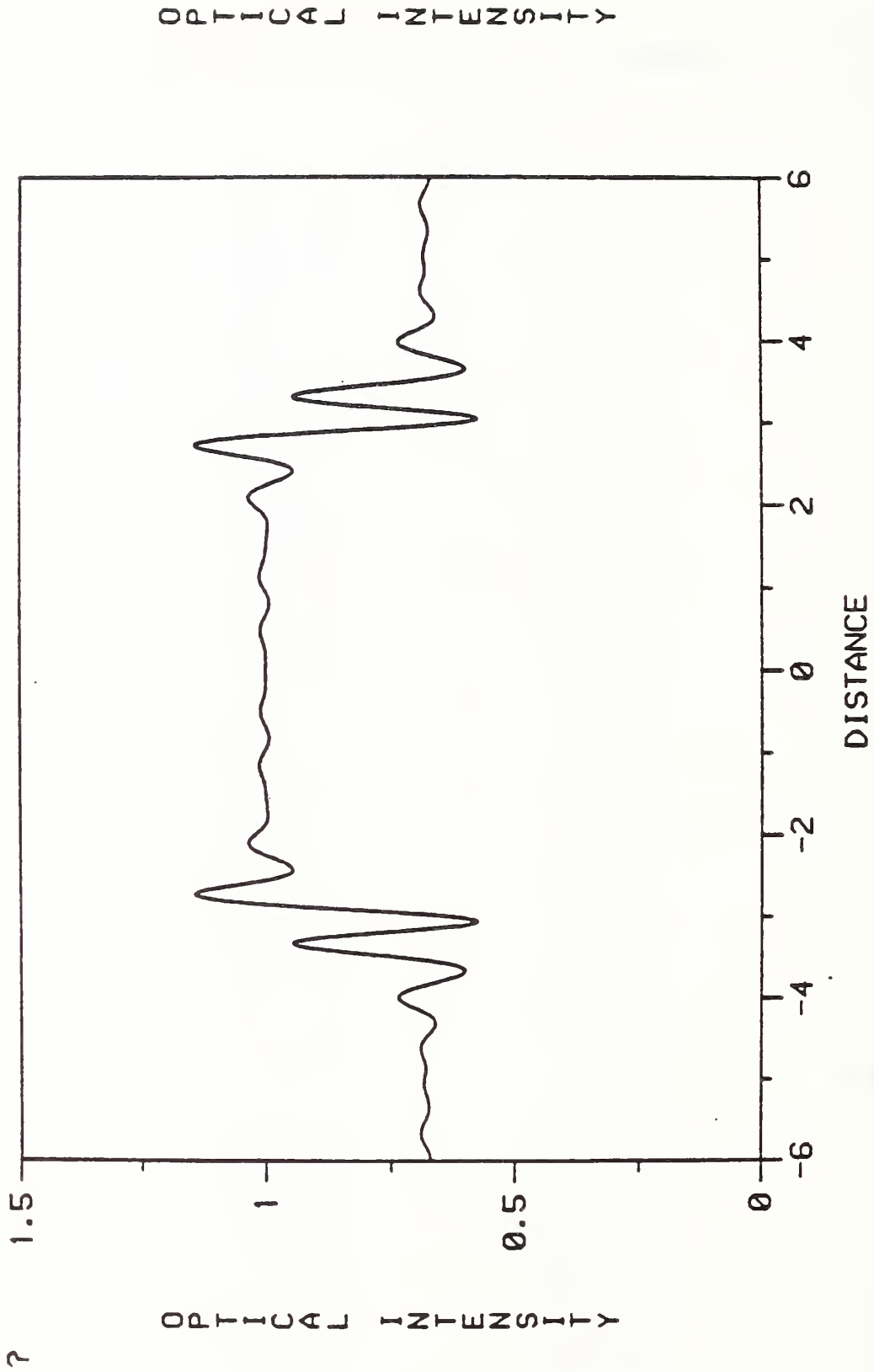
1	0.53	12.0	1	:	WAVELENGTH, PERIOD, NUMBER OF LAYERS
2	(1.00,0.00)	(4.10,0.06)	0	:	RI(AIR), RI(SUBSTRATE), ID
3	6.00	0.00	0.00	:	WIDTH, 3 POLYNOMIAL COEFFICIENTS
4	0.00	0.00	0.00	:	DISPLACEMENTS IN Z
5	(1.40,2.55)	0.09		:	RI(PATTERNED LAYER), LAYER THICKNESS
6	0.20	0.00		:	SLIT WIDTH, CAMERA WIDTH
7	0	0.00		:	SKIP, DEFOCUS



Test Case #2: File FCOFPO: Fourier Coefficients of the  
Pseudo-Object

1	.02598611	-.01092292
2	-.01353529	.00904450
3	-.01731358	.00311739
4	.01756685	-.00795909
5	.01356888	-.00057359
6	-.02047819	.00686085
7	-.01107900	-.00064297
8	.02327938	-.00598552
9	.00921602	.00126586
10	-.02648625	.00536360
11	-.00773381	-.00157148
12	.03061512	-.00501271
13	.00650267	.00169434
14	-.03646409	.00498365
15	-.00544231	-.00171108
16	.04564310	-.00541449
17	.00449766	.00167165
18	-.06227954	.00669869
19	-.00362787	-.00161367
20	.10149616	-.01032568
21	.00279663	.00156888
22	-.29965151	.02984478
23	-.14235833	-.46645156
24	-.29965151	.02984478
25	.00279663	.00156888
26	.10149616	-.01032568
27	-.00362787	-.00161367
28	-.06227954	.00669869
29	.00449766	.00167165
30	.04564310	-.00541449
31	-.00544231	-.00171108
32	-.03646409	.00498365
33	.00650267	.00169434
34	.03061512	-.00501271
35	-.00773381	-.00157148
36	-.02648625	.00536360
37	.00921602	.00126586
38	.02327938	-.00598552
39	-.01107900	-.00064297
40	-.02047819	.00686085
41	.01356888	-.00057359
42	.01756685	-.00795909
43	-.01731358	.00311739
44	-.01353529	.00904450
45	.02598611	-.01092292

CASE2, ONE LAYER, 3/20/87, THKBEST



Test Case #3: THICK DIELECTRIC LAYER WITH CURVED EDGES .  
 (6  $\mu\text{m}$  wide line with curved edges patterned in  
 a 0.65  $\mu\text{m}$  thick silicon dioxide layer on a  
 silicon substrate).

1	0.53	12.0	9	:	WAVELENGTH, PERIOD, NUMBER OF LAYERS	
2	(1.00,0.00)	(4.10,0.06)	0	:	RI(AIR), RI(SUBSTRATE), ID	
3	6.00	0.00	2.00	0.00	:	WIDTH, 3 POLYNOMIAL COEFFICIENTS
4	0.00	0.00	0.00	:	DISPLACEMENTS IN Z	
5	(1.46,0.00)	0.65	:	:	RI(PATTERNED LAYER), LAYER THICKNESS	
6	0.20	0.00	:	:	SLIT WIDTH, CAMERA WIDTH	
7	0	0.00	:	:	SKIP, DEFOCUS	

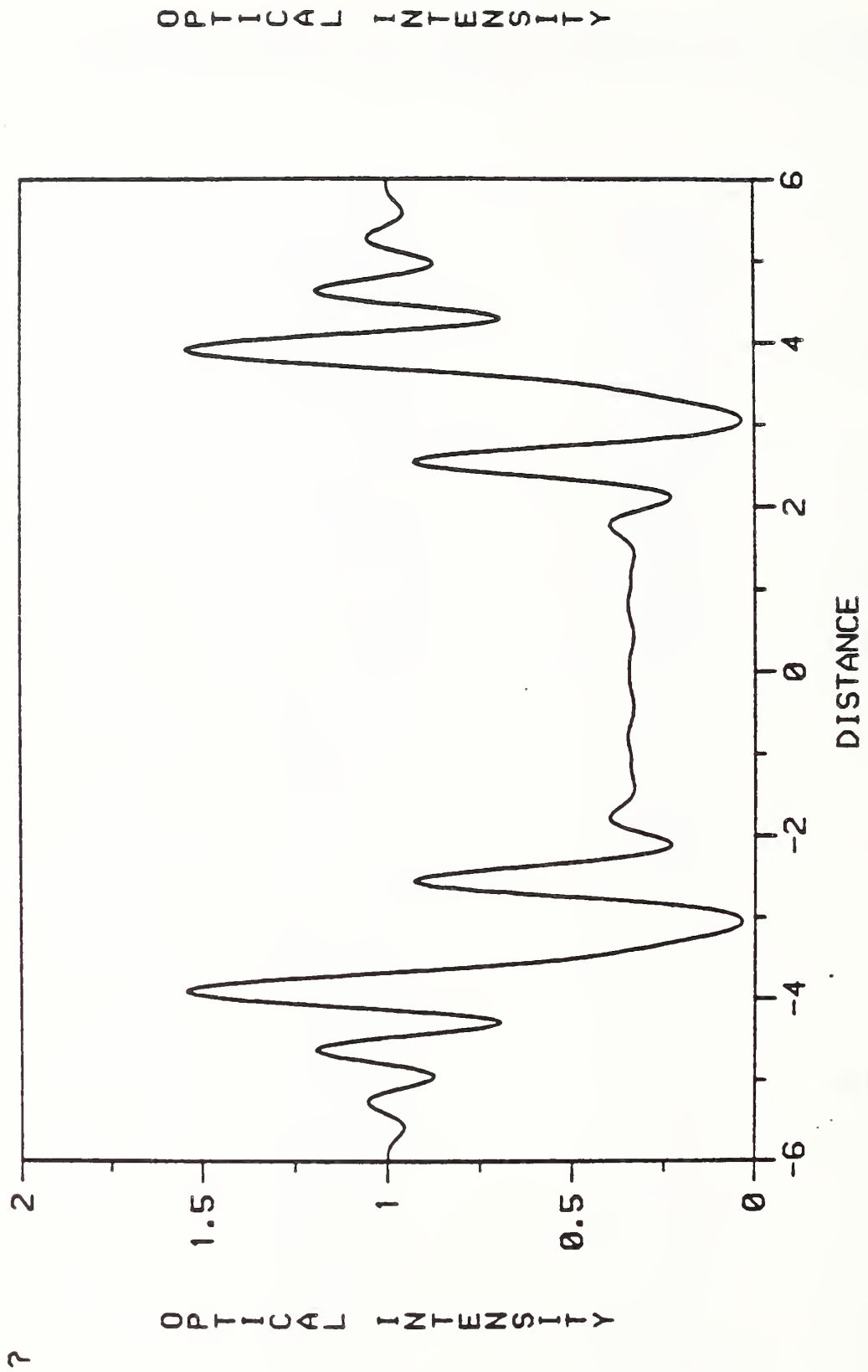
Test Case #3: File PARFIL: Parameter File Created

```
1  RUN PARAMETERS.
2  WAVELENGTH = .53
3  NUMBER OF LAYERS=9.
4  WAVE NUMBER =11.85506661732
5  AIR LAYER =(1.,0.)
6  SUBSTRATE =(4.1,.06)
7  DEFOCUS =0.
8  SLIT WIDTH IN MICRONS =.2
9  CAMERA WIDTH PARAMETER =0.
10 LAYER      WIDTH      POSITION      RI      OFFSET
11  1          6.003      .072        1.460   .000   .000
12  2          6.023      .144        1.460   .000   .000
13  3          6.065      .217        1.460   .000   .000
14  4          6.128      .289        1.460   .000   .000
15  5          6.211      .361        1.460   .000   .000
16  6          6.316      .433        1.460   .000   .000
17  7          6.441      .506        1.460   .000   .000
18  8          6.587      .578        1.460   .000   .000
19  9          6.754      .650        1.460   .000   .000
```

Test Case #3: File FCOFPO: Fourier Coefficients of the  
Pseudo-Object

1	.03128047	-.04799119
2	-.00117774	-.02619416
3	-.04538472	.02843279
4	.00285586	-.00315021
5	.04860053	.01296409
6	-.01663746	-.01247821
7	-.02455180	-.02653567
8	.00370391	.02884274
9	.01005389	.01942759
10	.01510389	-.02721254
11	-.00775574	-.00992680
12	-.02538226	.01395119
13	.01109421	.00413423
14	.02493682	.00192007
15	-.01569382	-.00211696
16	-.01660809	-.01617415
17	.01970374	.00272111
18	.00248225	.02957102
19	-.02276161	-.00472439
20	.02118083	-.04960428
21	.02524568	.00671538
22	-.10579441	.14143811
23	.38000129	.04024071
24	-.10579441	.14143811
25	.02524568	.00671538
26	.02118083	-.04960428
27	-.02276161	-.00472439
28	.00248225	.02957102
29	.01970374	.00272111
30	-.01660809	-.01617415
31	-.01569382	-.00211696
32	.02493682	.00192007
33	.01109421	.00413423
34	-.02538226	.01395119
35	-.00775574	-.00992680
36	.01510389	-.02721254
37	.01005389	.01942759
38	.00370391	.02884274
39	-.02455180	-.02653567
40	-.01663746	-.01247821
41	.04860053	.01296409
42	.00285586	-.00315021
43	-.04538472	.02843279
44	-.00117774	-.02619416
45	.03128047	-.04799119

CASE3N, 2/22/87, THKBEST



Test Case #4: THICK METAL LAYER WITH VERTICAL EDGES

(6  $\mu\text{m}$  wide line with vertical edges patterned  
in a 0.6  $\mu\text{m}$  thick chromium layer on a silicon  
substrate).

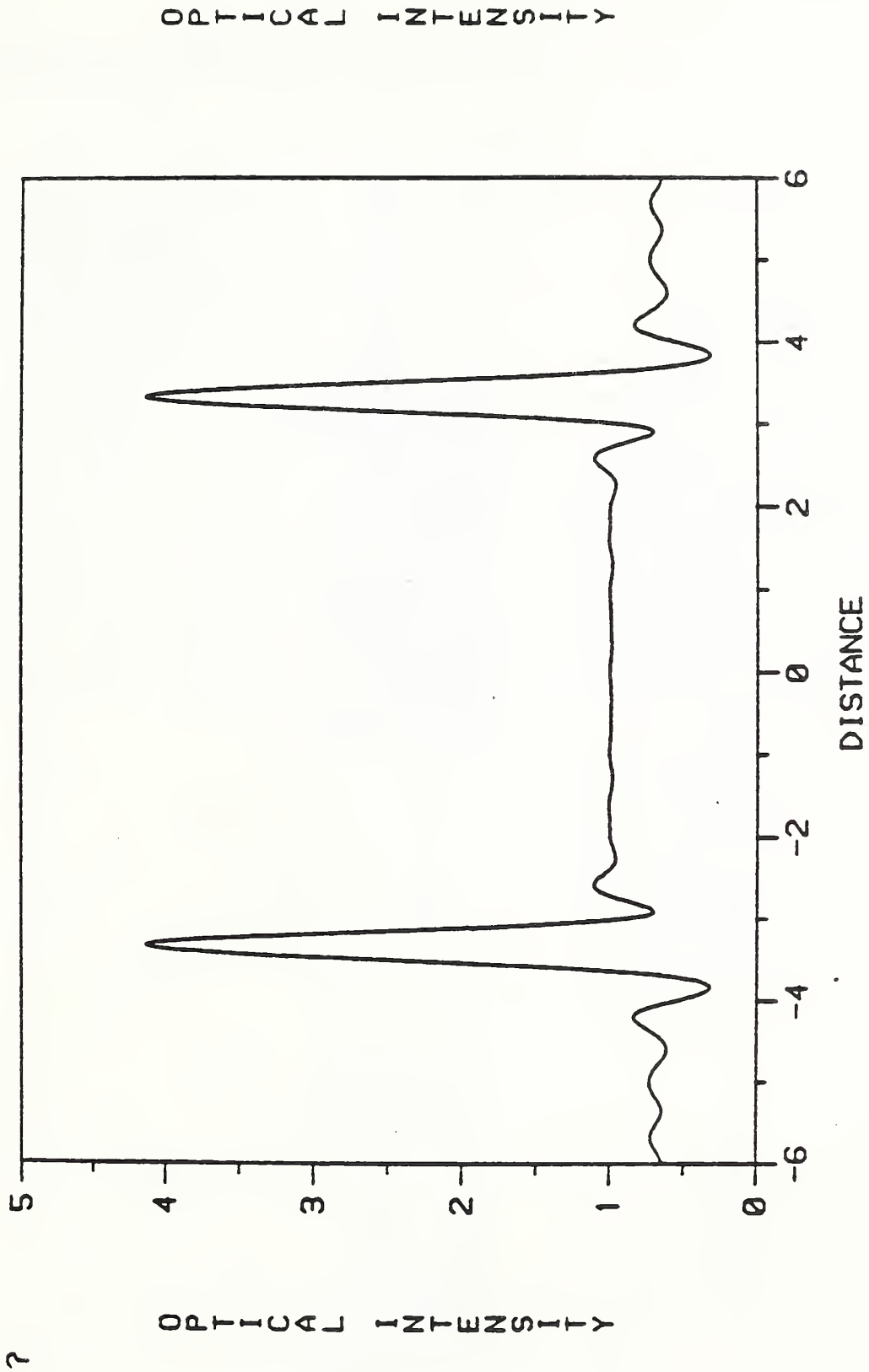
1	0.53	12.0	1	:	WAVELENGTH, PERIOD, NUMBER OF LAYERS
2	(1.00, 0.00)	(4.10, 0.06)	0	:	RI(AIR), RI(SUBSTRATE), ID
3	6.00	0.00	0.00	:	WIDTH, 3 POLYNOMIAL COEFFICIENTS
4	0.00	0.00	0.00	:	DISPLACEMENTS IN Z
5	(1.40, 2.55)	0.60		:	RI(PATTERNED LAYER), LAYER THICKNESS
6	0.20	0.00		:	SLIT WIDTH, CAMERA WIDTH
7	0	0.00		:	SKIP, DEFOCUS

Test Case #4: File FCOFPO: Fourier Coefficients of the  
Pseudo-Object

1	.02870492	-.00727836
2	-.03936288	-.01604885
3	.00858718	.02754465
4	.02527876	.00266374
5	-.01939032	-.03161583
6	-.00720966	.00412280
7	.01258914	.04300149
8	-.00154677	-.02651014
9	.00149106	-.03863546
10	-.00367560	.04706727
11	-.00975738	.02263821
12	.01454326	-.05862991
13	.01080683	-.00255415
14	-.02546317	.06130368
15	-.00753362	-.01701846
16	.03522425	-.05701120
17	.00344633	.03333270
18	-.04739825	.04842201
19	-.00070328	-.04416899
20	.07514600	-.04152960
21	-.00182714	.04869017
22	-.21542281	.06822844
23	-.26974701	-.56211569
24	-.21542281	.06822866
25	-.00182724	.04869002
26	.07514611	-.04152962
27	-.00070319	-.04416899
28	-.04739853	.04842213
29	.00344651	.03333271
30	.03522438	-.05701149
31	-.00753393	-.01701814
32	-.02546294	.06130364
33	.01080669	-.00255426
34	.01454345	-.05863000
35	-.00975762	.02263852
36	-.00367547	.04706706
37	.00149104	-.03863546
38	-.00154668	-.02651021
39	.01258891	.04300184
40	-.00720944	.00412244
41	-.01939040	-.03161581
42	.02527873	.00266401
43	.00858720	.02754443
44	-.03936286	-.01604884
45	.02870492	-.00727828



CASE4, 2/22/87, THKBEST



Test Case #5: THICK METAL LAYER WITH CURVED EDGES

(6  $\mu\text{m}$  wide line with curved edges patterned in  
 a 0.60  $\mu\text{m}$  thick silicon dioxide layer on a  
 silicon substrate).

1	0.53	12.0	9	:	WAVELENGTH, PERIOD, NUMBER OF LAYERS	
2	(1.00,0.00)	(4.10,0.06)	0	:	RI(AIR), RI(SUBSTRATE), ID	
3	6.00	0.00	2.00	0.00	:	WIDTH, 3 POLYNOMIAL COEFFICIENTS
4	0.00	0.00	0.00	:	DISPLACEMENTS IN Z	
5	(1.40,2.55)	0.60	:	:	RI(PATTERNED LAYER), LAYER THICKNESS	
6	0.20	0.00	:	:	SLIT WIDTH, CAMERA WIDTH	
7	0	0.00	:	:	SKIP, DEFOCUS	

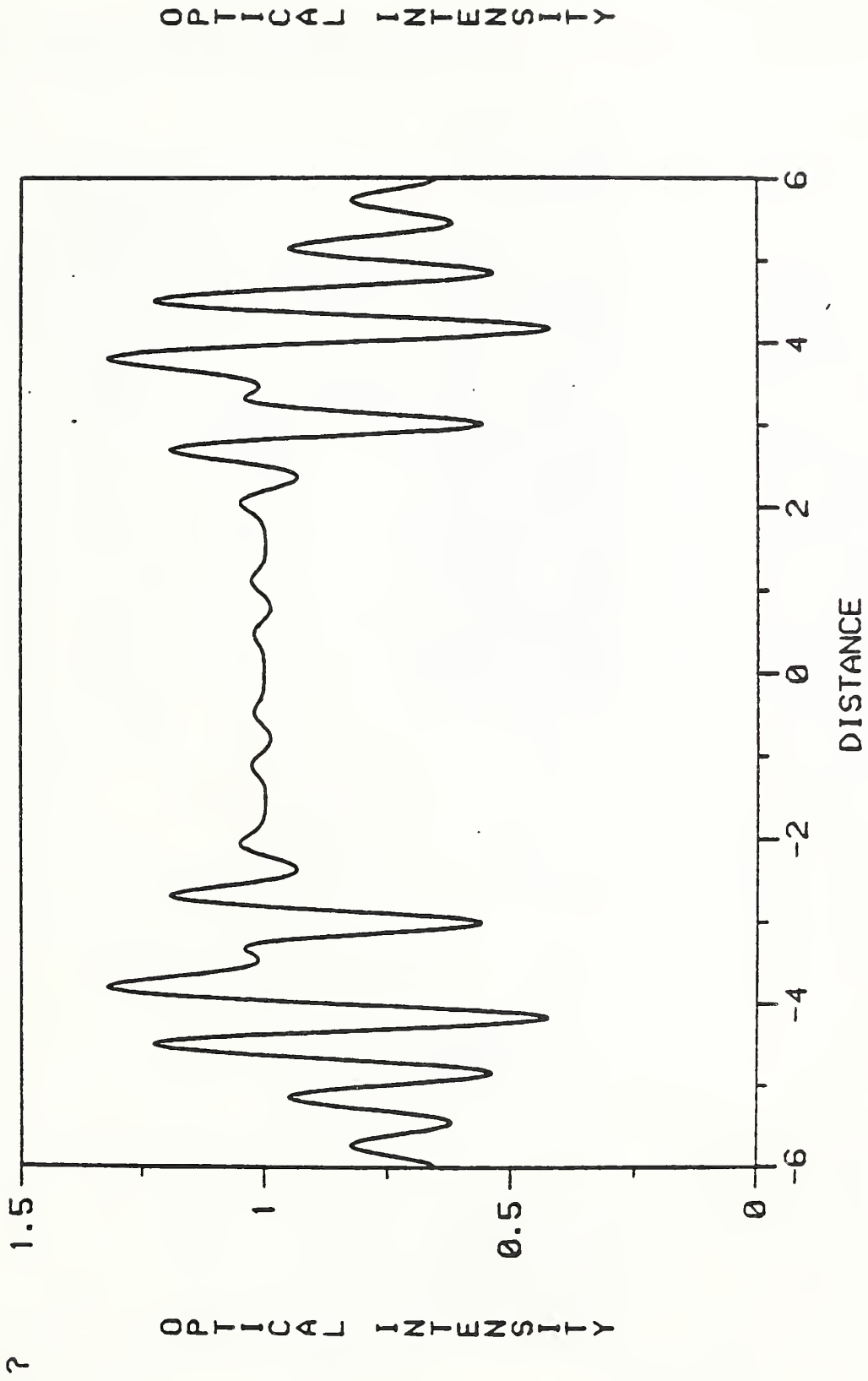
- Test Case #5: File PARFIL: Parameter File Created

```
1 RUN PARAMETERS.
2 WAVELENGTH = .53
3 NUMBER OF LAYERS=9.
4 WAVE NUMBER =11.85506661732
5 AIR LAYER =(1.,0.)
6 SUBSTRATE =(4.1,.06)
7 DEFOCUS =0.
8 SLIT WIDTH IN MICRONS =.2
9 CAMERA WIDTH PARAMETER =0.
10 LAYER      WIDTH      POSITION      RI          OFFSET
11  1         6.002      .067        1.400 2.550    .000
12  2         6.020      .133        1.400 2.550    .000
13  3         6.056      .200        1.400 2.550    .000
14  4         6.109      .267        1.400 2.550    .000
15  5         6.180      .333        1.400 2.550    .000
16  6         6.269      .400        1.400 2.550    .000
17  7         6.376      .467        1.400 2.550    .000
18  8         6.500      .533        1.400 2.550    .000
19  9         6.642      .600        1.400 2.550    .000
```

Test Case #5: File FCOFPO: Fourier Coefficients of the  
Pseudo-Object

1	.02859869	.01883469
2	-.01777494	-.05351227
3	-.01575921	.07268917
4	.01125702	-.04988998
5	.03861030	.00407903
6	-.05863673	.02278771
7	.00882784	-.01948456
8	.04261764	.00856627
9	-.02643598	-.00272235
10	-.01766390	-.00659771
11	.01759161	.01642617
12	.00861215	-.00775205
13	-.00318939	-.01494420
14	-.01427936	.01811821
15	-.00661590	.00686578
16	.02942126	-.02111680
17	.00859497	-.00085589
18	-.05174989	.02242236
19	-.00051795	-.00098279
20	.08337158	-.03235164
21	-.01442410	.00796601
22	-.20644118	.07115434
23	-.27539556	-.53212378
24	-.20644114	.07115399
25	-.01442399	.00796627
26	.08337133	-.03235159
27	-.00051788	-.00098292
28	-.05174964	.02242235
29	.00859462	-.00085587
30	.02942143	-.02111668
31	-.00661589	.00686571
32	-.01427944	.01811798
33	-.00318921	-.01494381
34	.00861184	-.00775220
35	.01759189	.01642605
36	-.01766396	-.00659770
37	-.02643602	-.00272207
38	.04261751	.00856600
39	.00882814	-.01948459
40	-.05863689	.02278783
41	.03861020	.00407919
42	.01125717	-.04989043
43	-.01575919	.07268954
44	-.01777511	-.05351234
45	.02859880	.01883462

CASE5, 2/22/87, THKBEST



Test Case #6: THICK POLYSILICON LAYER WITH CURVED EDGES

(6  $\mu\text{m}$  wide line with curved edges patterned in a 0.6  $\mu\text{m}$  thick polysilicon layer on a silicon substrate).

1	0.53	12.0	9	:	WAVELENGTH, PERIOD, NUMBER OF LAYERS	
2	(1.00,0.00)	(4.10,0.06)	0	:	RI(AIR), RI(SUBSTRATE), ID	
3	6.00	0.00	2.00	0.00	:	WIDTH, 3 POLYNOMIAL COEFFICIENTS
4	0.00	0.00	0.00	:	DISPLACEMENTS IN Z	
5	(3.80,0.10)	0.60	:	:	RI(PATTERNED LAYER), LAYER THICKNESS	
6	0.20	0.00	:	:	SLIT WIDTH, CAMERA WIDTH	
7	0	0.00	:	:	SKIP, DEFOCUS	

Test Case #6: File PARFIL: Parameter File Created

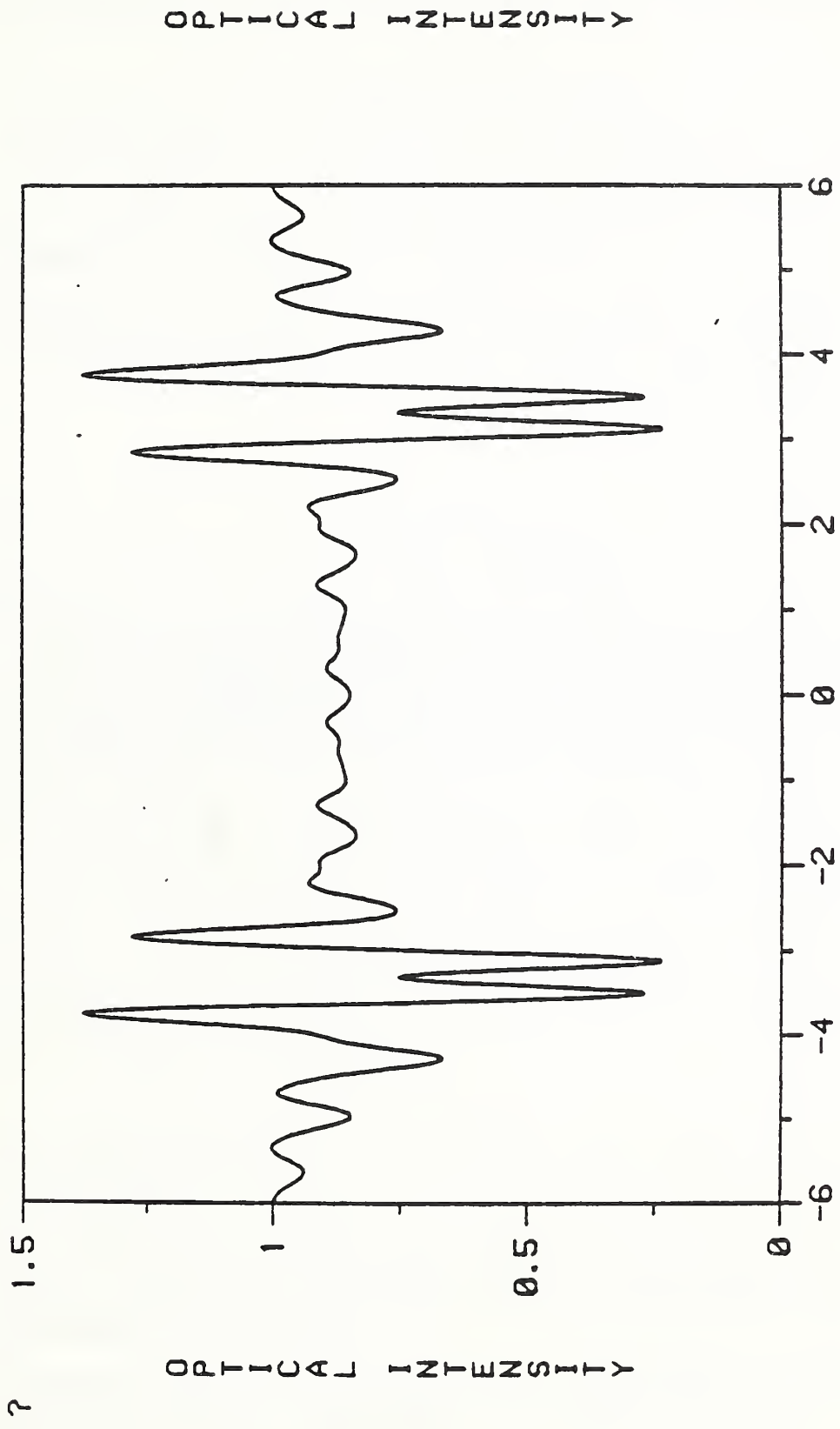
```
1  RUN PARAMETERS.
2  WAVELENGTH = .53
3  NUMBER OF LAYERS=9.
4  WAVE NUMBER =11.85506661732
5  AIR LAYER =(1.,0.)
6  SUBSTRATE =(4.1,.06)
7  DEFOCUS =0.
8  SLIT WIDTH IN MICRONS =.2
9  CAMERA WIDTH PARAMETER =0.
10 LAYER      WIDTH      POSITION      RI      OFFSET
11  1          6.002      .067        3.800   .100     .000
12  2          6.020      .133        3.800   .100     .000
13  3          6.056      .200        3.800   .100     .000
14  4          6.109      .267        3.800   .100     .000
15  5          6.180      .333        3.800   .100     .000
16  6          6.269      .400        3.800   .100     .000
17  7          6.376      .467        3.800   .100     .000
18  8          6.500      .533        3.800   .100     .000
19  9          6.642      .600        3.800   .100     .000
```

Test Case #6: File FCOFPO: Fourier Coefficients of the  
Pseudo-Object

1	-.06445069	.00175338
2	.10217661	-.01417309
3	-.06446184	.01696297
4	-.03134388	-.00120886
5	.07807896	-.01799169
6	-.01812634	.01902242
7	-.06195582	-.00139914
8	.04708281	-.01378403
9	.03549830	.01048018
10	-.05741779	.00419147
11	-.01086065	-.01166372
12	.05640716	.00465509
13	-.00874018	.00740007
14	-.04823824	-.01104616
15	.02114851	.00151754
16	.03683013	.01295479
17	-.02419279	-.01478962
18	-.02814956	-.00686661
19	.01508793	.03343811
20	.03801847	-.01927935
21	.00830037	-.06197198
22	-.16673161	.17059941
23	-.31875250	-.22803972
24	-.16673161	.17059941
25	.00830037	-.06197198
26	.03801847	-.01927935
27	.01508793	.03343811
28	-.02814956	-.00686661
29	-.02419279	-.01478962
30	.03683013	.01295479
31	.02114851	.00151754
32	-.04823824	-.01104616
33	-.00874018	.00740007
34	.05640716	.00465509
35	-.01086065	-.01166372
36	-.05741779	.00419147
37	.03549830	.01048018
38	.04708281	-.01378403
39	-.06195582	-.00139914
40	-.01812634	.01902242
41	.07807896	-.01799169
42	-.03134388	-.00120886
43	-.06446184	.01696297
44	.10217661	-.01417309
45	-.06445069	.00175338



CASE 6, 2/22/87, THKBEST



### Appendix III

Reprint - "Modeling of the optical microscope image of lines patterned in thick layers with variable edge geometry," by D. Nyysönen and C. P. Kirk (submitted to J. Opt. Soc. Am.).

Modeling of the optical microscope imaging of lines  
patterned in thick layers with variable edge geometry\*

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Abstract

A monochromatic, waveguide model is presented which can predict the optical microscope images of line objects with arbitrary edge geometry, patterned in thick-layers including multilayer structures with sloping, curved, and undercut edges, granular structures such as lines patterned in polysilicon, as well as asymmetric objects. The model is used to illustrate the effects of line edge structure on the optical image. Qualitative agreement with experimentally obtained optical image profiles is demonstrated. Application of the model to study the effects of variations in layer thickness and edge geometry on linewidth measurements made at different stages of manufacturing integrated-circuit devices is discussed.

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## INTRODUCTION

During the manufacture of integrated circuit wafers, critical dimensions of certain features must be measured at different stages of the production process. The control of the critical dimension of linewidth is usually monitored by using an optical or scanning-electron microscope. The measurements made by a correctly aligned microscope fitted with an appropriate measurement attachment are usually limited in accuracy and precision not by instrumental errors but by the lack of understanding of the characteristics of the image profile and the lack of accurate edge detection algorithms.

In the past, optical imaging in the microscope has been described by the scalar theory of partially coherent imaging [1-3] which characterizes the object by a planar complex transmittance or reflectance function. However, this approach does not accurately predict either the scattered field or the image structure for micrometer-sized line objects thicker than approximately one-quarter of the illuminating wavelength [3]. Unfortunately, most of the line features to be measured during the production of integrated circuits fall into this category of thick objects. Typical line objects are patterned in layers with thicknesses ranging from 0.3 to over 1.0  $\mu\text{m}$  (e.g. photoresist).

Integrated-circuit features encompass a wide variety of materials including dielectrics and absorbing materials of high and low conductivity (e.g., semiconductors in crystalline and amorphous states with varying amounts of dopant impurities, refractory metals, etc.) and line geometries with edge shapes varying from vertical to the complex shapes shown in Fig. 1. These lines may be etched into layers and be situated on top of metal or dielectric layers. Because of this wide variation, most of the approaches found in the literature are of only limited usefulness.

Scattering from micrometer-sized thick objects with complex indices of refraction falls in the domain of scattering by objects whose size is comparable to the illuminating wavelength such as Mie scattering [4] and requires an electromagnetic field treatment to accurately predict the scattered field. The various approaches that have been used to treat scattering by objects on the order of a wavelength can be found in the large volume of papers primarily in the areas of scattering by particles [5], scattering by surfaces with defects or protuberances [6], or diffraction by gratings [7]. Many of these approaches are limited either to objects of specific shapes [4,8] or weakly scattering objects which satisfy either the Rayleigh [9] or Born [10] approximations, or require infinite conductivity or specific grating geometries. Many are mathematically cumbersome when applied to the wide range of object shapes and materials of present interest.

Recently, the problem of imaging of line objects with vertical edge walls patterned in thick layers of dielectrics and metals was treated by Nyysönen [3,11] who developed a model based on the waveguide analysis for thick dielectric structures of Burckhardt [12]. Burckhardt's model was originally developed for diffraction from sinusoidal dielectric gratings and was applied to the reconstruction of bleached holograms. Kaspar [13,14] extended Burckhardt's approach to include absorbing photographic emulsions and non-sinusoidal grating structures and applied this approach to contact printing as well. Kaspar's work related the dielectric constant of the grating to the photographic density of the film.

The waveguide model previously used by Nyysönen [3,11] assumes that the line structure is patterned in a nonmagnetic layer which can be characterized by its complex index of refraction which is taken to be constant with depth within the layer. Thus, this model can be used to represent homogeneous line structures with

vertical edge walls. The spatial function representing the variation in the dielectric constant (square of the complex index of refraction) in the layer is expanded in a Fourier series. When this is substituted into the wave equation, Hill's equation [15] is obtained. The eigenvalue solutions to this equation represent waveguide modes. Assuming a single incident plane wave normal to the surface, the boundary value problem at the layer interfaces is solved to determine the Fourier coefficients in the expansions for the transmitted and reflected fields. This method allows the use of conventional scalar imaging theory to compute the image when no polarization effects are present. In such a case, the E- and H-field components are equivalent and either may be used in the scalar imaging equations.

In this paper, the waveguide model is extended to line objects whose index of refraction and geometry vary with depth in the layer, thus enabling line objects with nonvertical edges and multilayer structures to be considered. This extension of the model allows the images of virtually all line structures commonly encountered in IC fabrication to be modeled.

#### Thick-layer Model for Variable Edge Geometry

The method of solution for line structures of arbitrary edge geometry whose index of refraction may vary with depth in the layer is an extension of Nyysönen's earlier method used for vertical edge walls discussed above. A line object of the type to be treated in this paper is shown in Fig. 2a. This line structure has been chosen because it contains different materials, curved edges, asymmetry, and an unpatterned sub-layer. The line structure in the patterned layer varies in both width and dielectric constant as a function of depth  $z$  within the layer. Figure 2b shows how this line structure may be approximated by a multilayer set of line objects each of which has a constant index of refraction over a small interval of depth  $z$ . The line may thus

be represented by a set of layers with the complex dielectric constant  $\epsilon$  in each layer represented by the Fourier series expansion of:

$$\epsilon_n(x) = \begin{cases} \hat{\eta}_n^2, & \left(\frac{-W_n}{2} + \Delta_n\right) < x < \left(\frac{W_n}{2} + \Delta_n\right) \\ 1, & \text{between the lines} \end{cases} \quad (1)$$

where:

- $\hat{\eta}_n$  = refractive index of the nth layer of the line structure,
- $W_n$  = width of the nth layer, and
- $\Delta_n$  = offset of the nth layer in the x-direction. (See Fig. 2b)

Nonperiodic structures are taken as periodic with a very large repeating distance (period).

The eigenvalue solutions for the electromagnetic field components are found independently for each layer. The boundary value equations at each interface form a set of complex matrix equations which when solved allow the substitution of a single "equivalent" scattering layer for the multilayer structure. Thereafter, the solution for the transmitted and reflected field components and image are found in the same manner as for the single layer case [3,11].

Following Burckhardt's method [12] and assuming monochromatic illumination where  $\underline{E} = \underline{E}_0 e^{-i\omega t}$ , we start with the inhomogeneous wave equation for either the  $\underline{E}$ - or  $\underline{H}$ -field [16]:

$$\nabla^2 \underline{E} + \hat{\epsilon} \mu k_0^2 \underline{E} + \text{grad} \frac{\underline{E}}{\hat{\epsilon}} \cdot \text{grad} \hat{\epsilon} = 0 \quad (2)$$

$$\nabla^2 \underline{H} + \hat{\epsilon} \mu k_0^2 \underline{H} + \frac{\text{grad} \hat{\epsilon}}{\hat{\epsilon}} \times \text{curl} \underline{H} = 0 \quad (3)$$

with  $k_0^2 \hat{\epsilon} \mu = k_0^2 \hat{\eta}^2$  for nonmagnetic media.

If we choose the  $(x, y, z)$ -axes as we have done so that  $\hat{\epsilon}$  is constant along the length of the line in the  $y$ -direction, for TE-mode ( $\underline{E} = E_y \underline{y}$ ) and TM-mode ( $\underline{H} = H_y \underline{y}$ ) respectively, these equations simplify to

$$\nabla^2 E_y + k_0^2 \hat{\epsilon} E_y = 0 \quad (\text{TE-mode}) \quad (4)$$

and

$$\nabla^2 H_y - \frac{1}{\hat{\epsilon}} \frac{\partial \epsilon}{\partial x} \frac{H_y}{\partial x} + k_0^2 \hat{\epsilon} H_y = 0 \quad (\text{TM-mode}) \quad (5)$$

In the following analysis we consider only TE-mode because in the experimental NBS laser linewidth measuring system, the direction of polarization of the laser source is aligned parallel to the  $y$ -direction in the sample plane and the  $x$ -direction is the image scan direction. See Fig. 3. However, for any other system configuration and mode of illumination, the field can be considered as a superposition of TE- and TM-modes [17] and the corresponding analysis for TM-mode will be required. See ref. 12. Note that these modes are determined with respect to the orientation of the structure (not the incident wave) and are uncoupled (TE-mode:  $E_y$ ,  $H_x$ ,  $H_z$ , TM-mode  $H_y$ ,  $E_x$ ,  $E_z$ ).

The solution for the TE-mode is taken in the form  $E_y(x, z) = X(x) Z(z)$  and since  $\hat{\epsilon} = \hat{\epsilon}(x)$ , the method of separation of variables using a separation constant  $\alpha$  results in

$$\frac{\partial^2 X(x)}{\partial x^2} + k_0^2 \hat{\epsilon}(x) X(x) + \alpha^2 X(x) = 0 \quad (6)$$



$$\frac{\partial^2 Z}{\partial z^2} = \alpha^2 Z(z) \quad (7)$$

These equations hold for any functional form of  $\hat{\epsilon}(x)$ . However, following Burckhardt, we choose the Fourier approach where  $\hat{\epsilon}_n(x)$  from eq. (1) is represented by the Fourier series

$$\hat{\epsilon}_n(x) = \sum_q E_{q,n} \exp(2\pi i q x / P) \quad (8)$$

where

$$E_{q,n} = \begin{cases} (W_n/P) \hat{\eta}_n^2 + (1-W_n/P) \hat{\eta}_0^2 & q=0 \\ \frac{(\hat{\eta}_n^2 - \hat{\eta}_0^2)}{q\pi} \exp(2\pi i \Delta_n q / P) \sin(q\pi W_n / P) & q \neq 0 \end{cases}$$

where P is the period chosen arbitrarily large when the object consists of a single line and q is the summation index. Substituting eq. (8) into eq. (6) yields the differential equation known as Hill's equation

$$\frac{\partial^2 X(x)}{\partial x^2} + k_0^2 \left[ \sum_q E_{q,n} \exp(2\pi i q x / P) \right] X(x) + \alpha^2 X(x) = 0 \quad (9)$$

The solution for this general form of Hill's equation is given by Kaspar and others [13,15]. The form of the eigenvalue matrix is given in the Appendix. The solution of eq. (9) for the n-th layer is of the form

$$E_Y^n(x, z) = \sum_m \left[ A_{m,n} \exp(\alpha_{m,n} z) + A'_{m,n} \exp(-\alpha_{m,n} z) \right] \cdot \sum_j B_{j,m,n} \exp(2\pi i j x / P) \quad (10)$$

where the  $\alpha_{m,n}$ 's are the eigenvalues and the  $B_{j,m,n}$ 's are the eigenvector solutions to Hill's equation. The  $A_{m,n}$  and  $A'_{m,n}$  are weighting constants which must be determined from the boundary conditions. Each of these terms represents an inhomogeneous plane wave or waveguide mode which is supported by the line structure. Note that, when there is no absorption, ( $\hat{n}$  real) the eigenvalues are purely imaginary and this form reduces to that of Burckhardt.

The solution for the transmitted and reflected fields is found by equating the tangential components of the  $\underline{E}$ - and  $\underline{H}$ -fields at each layer interface. The x-component of the  $\underline{H}$ -field is found from  $ik_0 \underline{H} = \text{curl } \underline{E}$  or, here,  $ik_0 H_x = -\partial E_y / \partial z$ :

$$H_x^n(x, z) = \sum_m \left[ A_{m,n} \left( \frac{\alpha_{m,n}}{-ik_0} \right) \exp(\alpha_{m,n} z) - A'_{m,n} \left( \frac{\alpha_{m,n}}{-ik_0} \right) \exp(-\alpha_{m,n} z) \right] \cdot \sum_j B_{j,m,n} \exp(2\pi i j x / P). \quad (11)$$

With the tangential components of the incident field given by:

$$E_y^I = E_0^I \exp(ik_0 z) \quad (12a)$$

$$H_x^I = -E_0^I \exp(ik_0 z), \quad (12b)$$

the tangential components of the reflected field by:

$$E_y^R(x, z) = \sum_j E_j^R \exp \left\{ -ik_0 \left[ \left( \frac{\lambda_j}{P} \right) x + K_j^R z \right] \right\} \quad (13a)$$

$$H_x^R(x, z) = \sum_j K_j E_j^R \exp \left\{ -ik_0 \left[ \left( \frac{\lambda_j}{P} \right) x + K_j^R z \right] \right\}, \quad (13b)$$

$$\text{where } K_j^R = \sqrt{1 - \left( \frac{\lambda_j}{P} \right)^2}$$

and the tangential components of the transmitted field by:

$$E_y^T(x, z) = \sum_j E_j^T \exp \left\{ ik_0 \left[ \left( \frac{\lambda_j}{P} \right) x + K_j^T z \right] \right\} \quad (14a)$$

$$H_x^T(x, z) = \sum_j -K_j^T E_j^T \exp \left\{ ik_0 \left[ \left( \frac{\lambda_j}{P} \right) x + K_j^T z \right] \right\} \quad (14b)$$

where  $K_j^T = \sqrt{\hat{n}_s^2 - \left( \frac{\lambda_j}{P} \right)^2}$  and  $\hat{n}_s$  is the complex index of the substrate.

For each Fourier component  $j$ , the following boundary equations must be satisfied:

1. At  $z=0$  ( $x=y=0$ ), the solution must match the tangential components of the  $E$ - and  $H$ -fields at the top surface:

$$E_O^I \delta_{j0} + E_j^R = \sum_m \left[ A_{m,1} + A'_{m,1} \right] B_{j,m,1}; \quad (15a)$$

$$-E_O^I \delta_{j0} + K_j^R E_j^R = \sum_m \left[ A_{m,1} \left( \frac{\alpha_{m,1}}{-ik_0} \right) - A'_{m,1} \left( \frac{\alpha_{m,1}}{-ik_0} \right) \right] B_{j,m,1},$$

$$\text{where } \delta_{j0} = \begin{cases} 1 & j = 0 \\ 0 & j \neq 0. \end{cases}$$

2. At the interface between the n-th and (n+1)-th layer,  $z = -Z_n$  ( $x=y=0$ ), the  $\underline{E}$ - and  $\underline{H}$ -fields must satisfy the boundary conditions at each interface:

$$\sum_m \left[ A_{m,n} \exp(\alpha_{m,n} Z_n) + A'_{m,n} \exp(-\alpha_{m,n} Z_n) \right] B_{j,m,n} \quad (16a)$$

$$= \sum_m \left[ A_{m,n+1} \exp(\alpha_{m,n+1} Z_n) + A'_{m,n+1} \exp(-\alpha_{m,n+1} Z_n) \right] B_{j,m,n+1};$$

$$\sum_m \left[ A_{m,n} \left( \frac{\alpha_{m,n}}{-ik_0} \right) \exp(\alpha_{m,n} Z_n) - A'_{m,n} \left( \frac{\alpha_{m,n}}{-ik_0} \right) \exp(-\alpha_{m,n} Z_n) \right] B_{j,m,n}$$

$$= \sum_m \left[ A_{m,n+1} \left( \frac{\alpha_{m,n+1}}{-ik_0} \right) \exp(\alpha_{m,n+1} Z_n) - A'_{m,n+1} \left( \frac{\alpha_{m,n+1}}{-ik_0} \right) \exp(-\alpha_{m,n+1} Z_n) \right] B_{j,m,n+1} . \quad (16b)$$

3. At  $z = -Z_N = -T$  ( $x=y=0$ ), the solution must satisfy the boundary conditions imposed by the substrate (assumed infinite in extent - see Fig. 3):

$$\sum_m \left[ A_{m,N} \exp(\alpha_{m,N} T) + A'_{m,N} \exp(-\alpha_{m,N} T) \right] B_{j,m,N} \quad (17a)$$

$$= E_j^T \exp(-ik_0 K_j^T T) ;$$

$$\sum_m \left[ A_{m,N} \left( \frac{\alpha_{m,N}}{-ik_0} \right) \exp(-\alpha_{m,N} T) - A'_{m,N} \left( \frac{\alpha_{m,N}}{-ik_0} \right) \exp(\alpha_{m,N} T) \right] B_{j,m,N}$$

$$= -K_j^T E_j^T \exp(-ik_0 K_j^T T) . \quad (17b)$$

In matrix notation,

$$\begin{pmatrix} E_x^{IR} \\ H_y^{IR} \end{pmatrix} = \begin{pmatrix} B_1(0) \end{pmatrix} \cdot \begin{pmatrix} A_1 \\ A'_1 \end{pmatrix} \quad (18a)$$

$$\begin{pmatrix} B_n(z_n) \end{pmatrix} \cdot \begin{pmatrix} A_n \\ A'_n \end{pmatrix} = \begin{pmatrix} B_{n+1}(z_n) \end{pmatrix} \cdot \begin{pmatrix} A_{n+1} \\ A'_{n+1} \end{pmatrix} \quad (18b)$$

$$\begin{pmatrix} B_N(T) \end{pmatrix} \cdot \begin{pmatrix} A_N \\ A'_N \end{pmatrix} = \begin{pmatrix} E_x^T \\ H_y^T \end{pmatrix} \quad (18c)$$

where

$$B_n(z_n) = \begin{pmatrix} B_n^{11}(z_n) & B_n^{12}(z_n) \\ B_n^{21}(z_n) & B_n^{22}(z_n) \end{pmatrix}$$

and the matrix elements are given by

$$B_{j,m,n}^{11}(z_n) = \exp(-\alpha_{m,n} z_n) B_{j,m,n}$$

$$B_{j,m,n}^{12}(z_n) = \exp(\alpha_{m,n} z_n) B_{j,m,n}$$

$$B_{j,m,n}^{21}(z_n) = \left( \frac{\alpha_{m,n}}{-ik_0} \right) \exp(-\alpha_{m,n} z_n) B_{j,m,n}$$

$$B_{j,m,n}^{22}(z_n) = - \left( \frac{\alpha_{m,n}}{-ik_0} \right) \exp(\alpha_{m,n} z_n) B_{j,m,n}$$

Depending upon whether the solution for the transmitted or reflected field is desired, all of the  $A_n$  and  $A_n'$ s are eliminated from these equations except for either  $A_N$  and  $A_N'$  or  $A_1$  and  $A_1'$ , i.e., for the transmitted field:

$$\begin{pmatrix} A_1 \\ A_1' \end{pmatrix} = \left( B_1(z_1) \right)^{-1} \cdot \left( B_2(z_1) \right) \cdot \left( B_2(z_2) \right)^{-1} \dots \quad (19a)$$

$$\cdot \left( B_{N-1}(z_{N-1}) \right)^{-1} \cdot \left( B_N(z_{N-1}) \right) \cdot \begin{pmatrix} A_N \\ A_N' \end{pmatrix} = B \cdot \begin{pmatrix} A_N \\ A_N' \end{pmatrix}$$

or, for the reflected field:

$$\begin{pmatrix} A_N \\ A_N' \end{pmatrix} = \left( B_N(z_{N-1}) \right)^{-1} \cdot \left( B_{N-1}(z_{N-1}) \right) \cdot \left( B_{N-1}(z_{N-2}) \right)^{-1} \dots \quad (19b)$$

$$\cdot \left( B_2(z_1) \right)^{-1} \cdot \left( B_1(z_1) \right) \cdot \begin{pmatrix} A_1 \\ A_1' \end{pmatrix} = B^{-1} \cdot \begin{pmatrix} A_1 \\ A_1' \end{pmatrix}$$

This operation replaces the multilayer structure with an equivalent layer characterized by the matrix  $B$ .

The equations are now solved as for the single layer case by eliminating the unknown  $E_j^R$  and  $E_j^T$  to solve for the  $A_1$ s or  $A_N$ s again depending upon whether the Fourier expansion of the reflected or transmitted field is desired. In the present paper, we compute the reflected field coefficients,

$$\begin{pmatrix} E_x^{IR} \\ H_y^{IR} \end{pmatrix} = \left( B_1(0) \right) \cdot \begin{pmatrix} A_1 \\ A_1' \end{pmatrix} \quad (20a)$$

$$\begin{pmatrix} B^{11} & B^{12} \\ B^{21} & B^{22} \end{pmatrix} \cdot \begin{pmatrix} A_1 \\ A_1' \end{pmatrix} = \begin{pmatrix} E_x^T \\ H_y^T \end{pmatrix} \quad (20b)$$

where

$$\begin{pmatrix} B^{11} & B^{12} \\ B^{21} & B^{22} \end{pmatrix} = \left( B_N(T) \right) \cdot B^{-1} \quad (20c)$$

Using the format of eqs. (18) this reduces to the single matrix equation:

$$\begin{pmatrix} D_{11} & D_{12} \\ D_{21} & D_{22} \end{pmatrix} \cdot \begin{pmatrix} A_1 \\ A_1' \end{pmatrix} = \begin{pmatrix} R \\ 0 \end{pmatrix} \quad (21)$$

where

$$D_{j,m}^{11} = \left[ K_j^R - \left( \frac{\alpha_{m,1}}{-ik_0} \right) \right] B_{j,m,1}$$

$$D_{j,m}^{12} = \left[ K_j^R + \left( \frac{\alpha_{m,1}}{-ik_0} \right) \right] B_{j,m,1}$$

$$D_{j,m}^{21} = \left[ K_j^T B_{11} + B_{21} \right]$$

$$D_{j,m}^{22} = \left[ K_j^T B_{12} + B_{22} \right]$$

and

$$R_j = \left[ E_0^I (1 + K_j^R) \delta_{j0} \right]$$

where the desired coefficients  $E_j^R$  are found from eq. (15a)

$$E_j^R = \sum_m \left[ A_{m,1} + A'_{m,1} \right] B_{j,m,1} - E_0^I \delta_{j0} \quad (22)$$

These coefficients represent the Fourier coefficients of a pseudo or equivalent (planar) object. The microscope image of the reflected field is then computed using scalar imaging theory (See Ref. 3.) with the line object reflectance function given by

$$t(x) = \sum_j E_j^R \exp (ik_0 \lambda jx/P) \quad (23)$$

This method of computing both the reflected field and the corresponding microscope image requires no approximations of the type usually found, such as limits on the conductivity or slope of the surface, etc. Limitations may be imposed, however, by the computation capability available. First, the number of layers used to approximate the structure increases the computing time linearly. In most of the cases to be shown here, seven to nine layers were sufficient to produce significant results and required approximately one minute of CPU time on a Univac 1108.\*

The second limitation is in the truncation of the series, i.e., the matrix sizes used in the computations. In the present case, as for a single layer [3], all of the reflected plane waves are included which have diffraction angles less than  $\pm \pi/2$  in air. With  $P = 12\mu\text{m}$  and  $\lambda = 0.53\mu\text{m}$ , 22 diffracted orders are included which requires a 45 x 45 complex eigenvalue matrix and a 90 x 90

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\*Certain commercial equipment, instruments, or materials are identified in this paper in order to adequately specify experimental or computational procedure. Such identification does not imply recommendation or endorsement by the National Bureau of Standards, nor does it imply that the materials or equipment identified are necessarily the best available for the purpose.



complex matrix for inversion of the B matrices in eqs. (19). This choice necessarily truncates the series which represents the field in the layers with higher refractive index. In the cases considered here, this truncation does not appear to significantly affect the results.

Also, for grating objects with  $P \leq 12\mu\text{m}$ , the computations are exact. However, for isolated line objects near resonances (where the thickness times index of refraction of the patterned layer is approximately equal to the wavelength),  $P = 12\mu\text{m}$  is not large enough to eliminate the effect of the adjacent lines on the image. Larger matrix sizes would have to be used where  $P > 12\mu\text{m}$  is required.

### The Effect of Geometry on the Image Profile

The present method was developed primarily for the purpose of modeling the effect of nonvertical edges on the image profile. There is a very wide range of edge geometries which can be modeled, therefore, only a few key examples will be presented here in order to demonstrate trends. Figure 1 shows six basic shapes which are frequently encountered in integrated circuit processing. In order to simplify the definition of these shapes, they will be represented by a polynomial expansion which defines the width as a function of  $z$ .

$$W(z) = \sum_{j=0}^J X_j (z - Z_j)^j, \quad (24)$$

where  $J$  = the polynomial order,

$X_j$  = coefficient of the  $j$ -th order, and

$Z_j$  = offset of the  $j$ -th order.

The shapes in Fig. 1 have been restricted to fifth order polynomials ( $J \leq 5$ ). In these examples it will be assumed that the patterned layer is  $0.6 \mu\text{m}$  thick, and the nominal linewidth  $X_0$  is  $6.0 \mu\text{m}$ . The line consists of loss-free silicon dioxide with a real refractive index of 1.46. The substrate is taken to be silicon with a complex refractive index of  $4.1 + 0.06i$  at an illumination wavelength of  $0.53 \mu\text{m}$ . The diffraction-limited images have been computed for a 0.14 NA illumination aperture and 0.85 NA objective aperture. Note that in this polynomial representation of the line object, although  $X_0$  as shown in Fig. 1 is the nominal linewidth,  $X_0$  is: the width at the top ( $z=0$ ) in cases (b) and (g); the width at the bottom ( $z=d$ ) in cases (c), (d), and (f), and the mean width in cases (e) and (h). Figure 4 shows the theoretical image intensity profile of an ideal line with vertical edge walls. The oxide line extends from  $-3.0$  to  $+3.0 \mu\text{m}$  as shown by the dashed lines, and the rest of the object is bare silicon. In the calculation of the Fourier series coefficients, this structure is assumed to be repeated with a period of  $12.0 \mu\text{m}$ . The line edge image is characterized by a dark fringe wider than the interference fringe that would be calculated for a similar thin layer.

The set of image profiles in Fig. 5 demonstrates the effect of different edge geometries on the image. The line has been assumed to be symmetrical and so only half the image profile is shown. The line objects are defined in terms of Eq. (2) and the width of the line edge is getting progressively broader down each of the three columns. There are two effects which are common to all three geometries. As the edge becomes broader, the dark fringe associated with the edge also becomes broader with the bright fringes on either side becoming brighter than occurs with partially coherent imaging of planar objects. When the edge becomes very broad, the dark fringe itself begins to broaden out and small peaks form within it.

## Modeling Granular Materials

According to the method described so far, the object is assumed to consist of homogeneous slabs of materials. This assumption is implicit in Eq. (1). In practice many materials such as polycrystalline silicon cannot be regarded as homogeneous as they have a definite internal structure which will result in a variation in refractive index within the material. These structures can be modeled by representing the random refractive index disturbances within the material by functions.

$$\hat{\epsilon}'_n(x) = \hat{\epsilon}_n(x) + \sum_{j=1}^J C_j \delta(x-X_j) , \quad (25)$$

where:

- $\hat{\epsilon}_n(x)$  = refractive index profile from Eq. (1),
- $C_j$  = amplitude of j-th refractive index disturbance, and
- $X_j$  = position of j-th disturbance.

The function  $\hat{\epsilon}_n(x)$  now describes a noisy layer, and by restricting the range of values for  $X_j$ , the noise may be confined to different parts of the layer.

Figure 6 shows three image profiles for a noisy structure. The object is the oxide structure used to produce the profile in Fig. 4 and noise has been added to the entire layer. In each of the three cases shown in Fig. 6, different sets of noise data have been used corresponding to different sets of  $C_j$  selected from a random number table. The most striking feature of these curves is that, although the field is perturbed within the area of the object, the dark fringe at the edge of the line remains well defined, and relatively unperturbed.

## Comparison with Experimental Results

In order to test the model against experimental image profiles, a test specimen was prepared which consisted of lines patterned in photoresist on a silicon substrate. The same test pattern was put down repeatedly by projection printing but with a range of focus positions. This produced a series of patterns with a range of edge slopes. Two patterns were selected to illustrate different edge properties. One pattern corresponded to the in-focus exposure and had near vertical edge walls; the second corresponded to a considerable defocus on exposure and had significantly sloping edges.

The in-focus exposure was assumed to produce near vertical edge walls, and when the structures were examined in a scanning electron microscope, this was found to be the case. The edge walls had slopes of about  $80^\circ$ , which according to the theory may effectively be assumed to be vertical for films of submicrometer thickness. The width of the window in the resist was taken as  $2.5 \mu\text{m}$ .

The out-of-focus pattern had significant edge slopes, and the structure was modeled by the asymmetric third order window shown in Fig. 7, based on SEM pictures of the cross section. The photoresist thickness was not known exactly but was nominally  $1 \mu\text{m}$ . The refractive index of the resist is not known exactly, but this is not serious as small changes in the refractive index can be offset by small adjustments in the layer thickness without significantly changing the image profile. The refractive index of the photoresist was assumed to be 1.513. The theoretical layer thickness was adjusted until it produced the same contrast as the experimental data; this gave a thickness for the resist of  $0.94 \mu\text{m}$ .

It is very difficult to accurately determine the position corresponding to "best focus" when viewing these thick structures in

the optical linewidth measurement system. Therefore, it is necessary to generate a series of profiles as a function of focus position and compare these with the theoretical data. Figure 8 shows a comparison of the experimental and theoretical image profiles produced by these structures. The experimental data were generated using the coherent optical linewidth measurement system which has been described in detail in the literature by Nyyssonen.[3]

The profiles for the sloping edge structure are quite different from those for the vertical edge structure. However, in both cases they show good agreement with the models. The theoretical profiles display the same features as the experimental profiles, but the agreement is not perfect. One significant source of differences is the assumption that the resist is homogeneous. This is not strictly true. To get better agreement, it would have been advantageous to use a material such as silicon dioxide which is homogeneous and which would produce a stable and accurately known refractive index.

A second source of differences may arise from the assumption that the illumination is normally incident. As shown in the experimental data, the central peak of the sloping edge profile moves as the object moves through focus and the line with nearly vertical edges generates slightly asymmetric profiles (See Fig. 8). Both of these may be attributed to slight asymmetry in either the illumination system or imaging optics. The NBS profiling system used to generate these profiles was checked for illumination symmetry when profiling thin layers. However, it is expected that these thick layers are significantly more sensitive to asymmetric illumination than is a thin-layer object.

#### Case Study of the Production of an MOS Transistor

This section considers some typical structures which are encountered in microlithography and the effect of small variations in

these structures on the optical image profile. The two structures shown in Fig. 9 are considered. These structures represent the key stages of patterning the polysilicon layer when making an MOS device. The structure in Fig. 9(a) represents the stage between patterning the resist and etching the polysilicon. Controlling the resist dimensions at this stage will help control the final etched polysilicon linewidth. However, measuring the width of the resist line at this stage is difficult because the image profile is a function of the geometry, refractive index, and thickness of each of the three layers. The polysilicon was assumed to have a refractive index of  $3.8 + 0.1i$ .

The curves in Fig. 10 show how variations in the thickness of the polysilicon affect the image profile of the resist line. Again, only half the image profile is shown. The most striking feature of these curves is that the image structure in the immediate vicinity of the line edge is sensitive to the thickness of the polysilicon layer.

The structure in Fig. 11 represents a polysilicon line after etching and stripping of the resist. The image profile is a function of the thickness of the silicon-dioxide layer. Figure 12 shows the effect of changing the thickness of the oxide sublayer on the image profile of the structure shown in Fig. 9(b). The polysilicon is assumed in this case to have vertical edge walls and the oxide thickness is varied over a range of 40 nm. As the thickness changes, the image profile changes considerably. This makes linewidth measurement difficult as there appears to be no feature or threshold which locates the edge independently of oxide thickness.

The image profile is also a function of edge geometry and in Fig. 9(b) it has been assumed that the edge shape can be defined by a second order polynomial with the polynomial of Eq. (2) given by

$$W(z) = X_0 + X_2 \cdot z^2 , \quad (26)$$

where, again, the nominal linewidth  $X_0$  is  $6.0 \mu\text{m}$ . The curves in Fig. 12 show the effect of variations in the edge curvature of the polysilicon layer on the image profile. The edge image profile is clearly sensitive to the curvature of the physical edge of a curved feature. In this case, three possible definitions of linewidth are the minimum, maximum, or mean width of the layer. For the examples shown in Fig. 12, these definitions all give linewidths which do not correspond to that given by the position of the minimum of the dark fringe. More importantly, however, the offset between the true and measured linewidth varies with the degree of curvature over a range of a few tenths of a micrometer regardless of the definition used for the line edge.

### Summary

A waveguide model has been presented which enables the optical images of line structures patterned in thick layers to be computed. The model has been shown to be applicable to a wide range of structures, and qualitative agreement with experimental image profiles has been demonstrated. From the results presented here, it can be seen that the image profile of a thick line object, and therefore the measured linewidth, is affected by the physical shape of the edges.

The model enables the effects of process variations on the optical image profile to be determined and this, in turn, enables the accuracy and sensitivity of different measurement techniques to be investigated. Further work is in progress to improve existing measurement methods in order to minimize the effects of process variations on the accuracy and repeatability of linewidth measurements.

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## Appendix

In matrix notation, the eigenvalue problem of eq. (9) is:

$$D \cdot B_n = \alpha^2 I \cdot B_n$$

where the matrix elements  $D_{i,j}$  for the n-th layer are given by

$$D = \begin{pmatrix} -(E_{0,n} + K_j^2) & E_{1,n} & \cdot & \cdot & \cdot & E_{i-j,n} & 0 & \cdot & \cdot & \cdot & 0 \\ E_{1,n} & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & E_{1,n} & \cdot & \cdot & \cdot & \cdot & 0 \\ E_{i-j,n} & \cdot & \cdot & \cdot & \cdot & E_{1,n} & -(E_{0,n} + K_0^2) & E_{1,n} & \cdot & \cdot & E_{i-j,n} \\ 0 & \cdot & \cdot & \cdot & \cdot & E_{1,n} & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & E_{1,n} \\ 0 & \cdot & \cdot & \cdot & \cdot & \cdot & 0 & E_{i-j,n} & \cdot & \cdot & \cdot & E_{1,n} & -(E_{0,n} + K_j^2) \end{pmatrix}$$

and  $K_j = \frac{\lambda_j}{P}$

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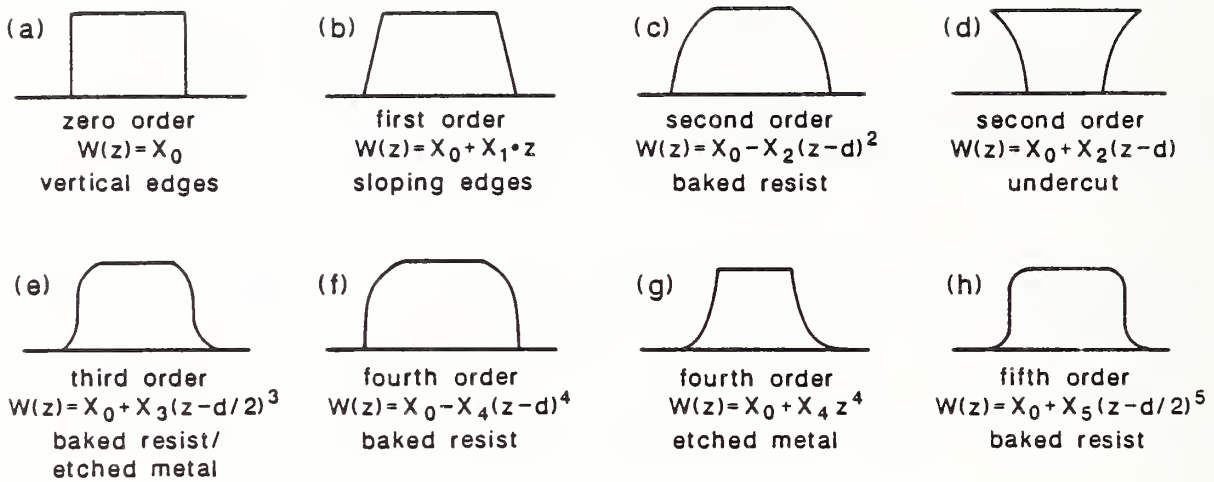


Fig. 1. The geometry of most line objects encountered in integrated circuits may be approximated by low order polynomials.

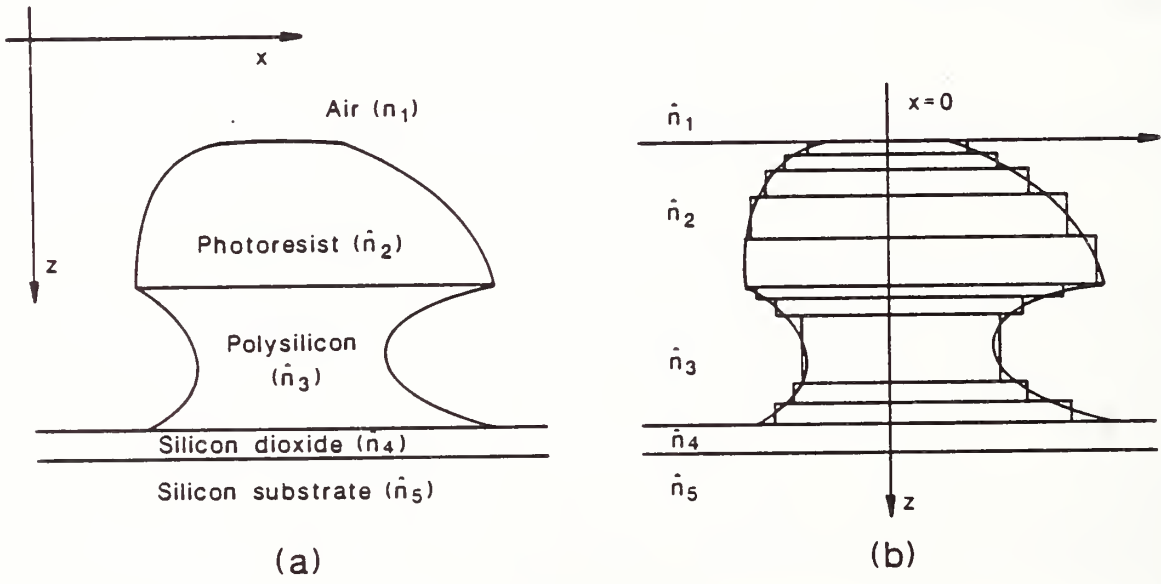


Fig. 2. Cross section of a hypothetical thick line object (a) and the corresponding multi-layer representation (b).

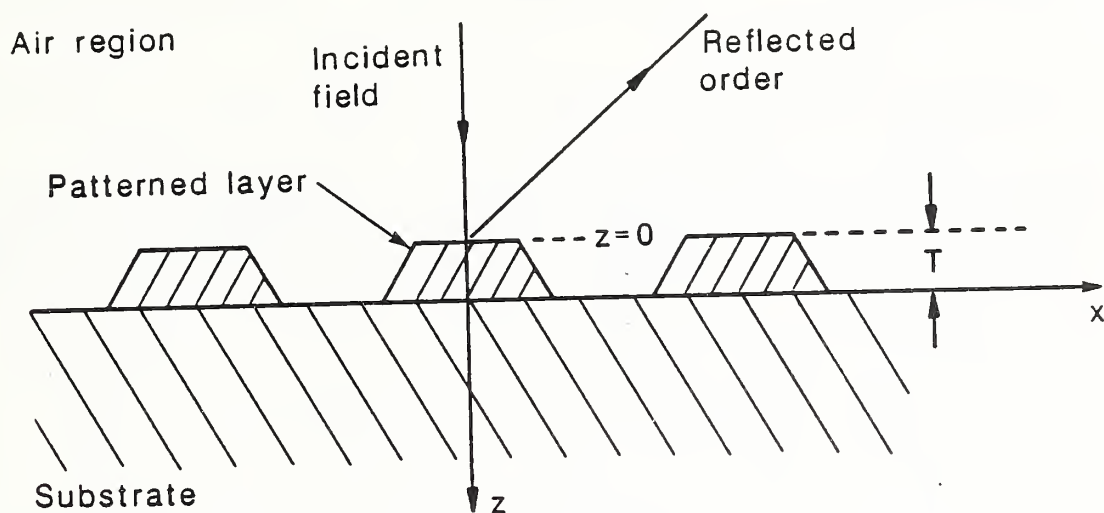


Fig. 3. Orientation of the line structure and incident fields in the NBS linewidth measuring system. The y-direction is out of the page and parallel to the length of the line patterns.

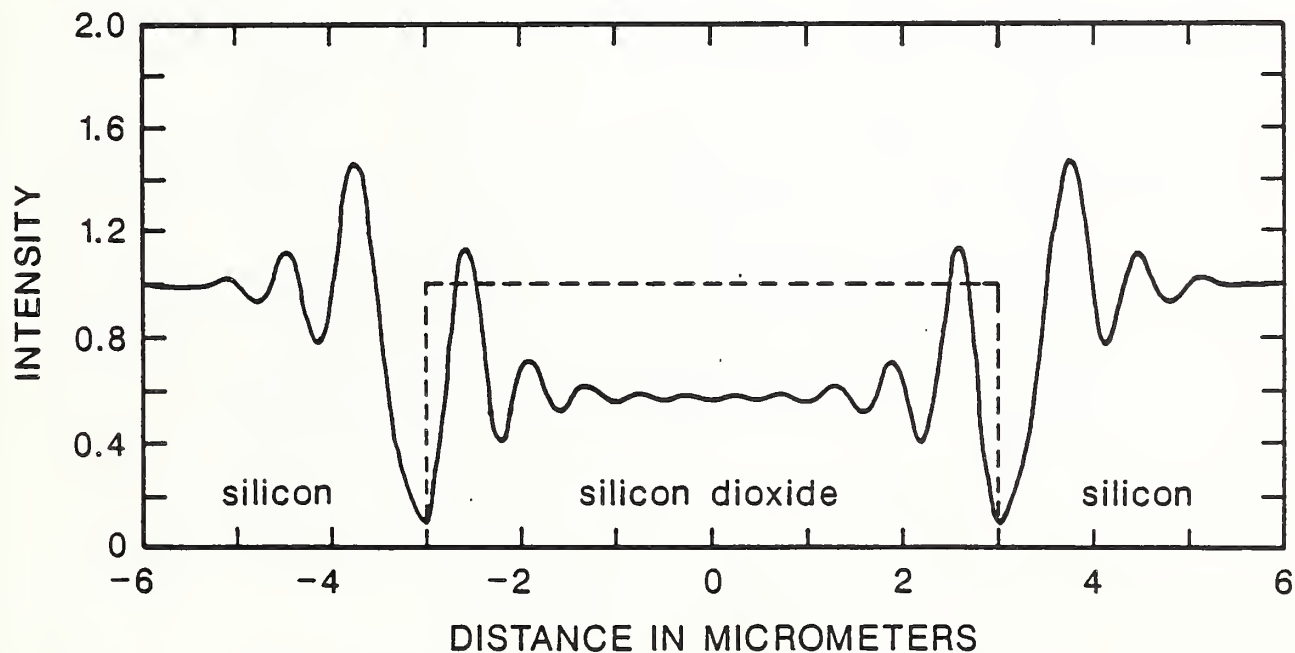


Fig. 4. Theoretical image intensity profile of a  $6.0 \mu\text{m}$  wide line centered at zero, patterned in a  $0.6 \mu\text{m}$  thick  $\text{SiO}_2$  layer on silicon.  $\lambda = 530 \text{ nm}$ .

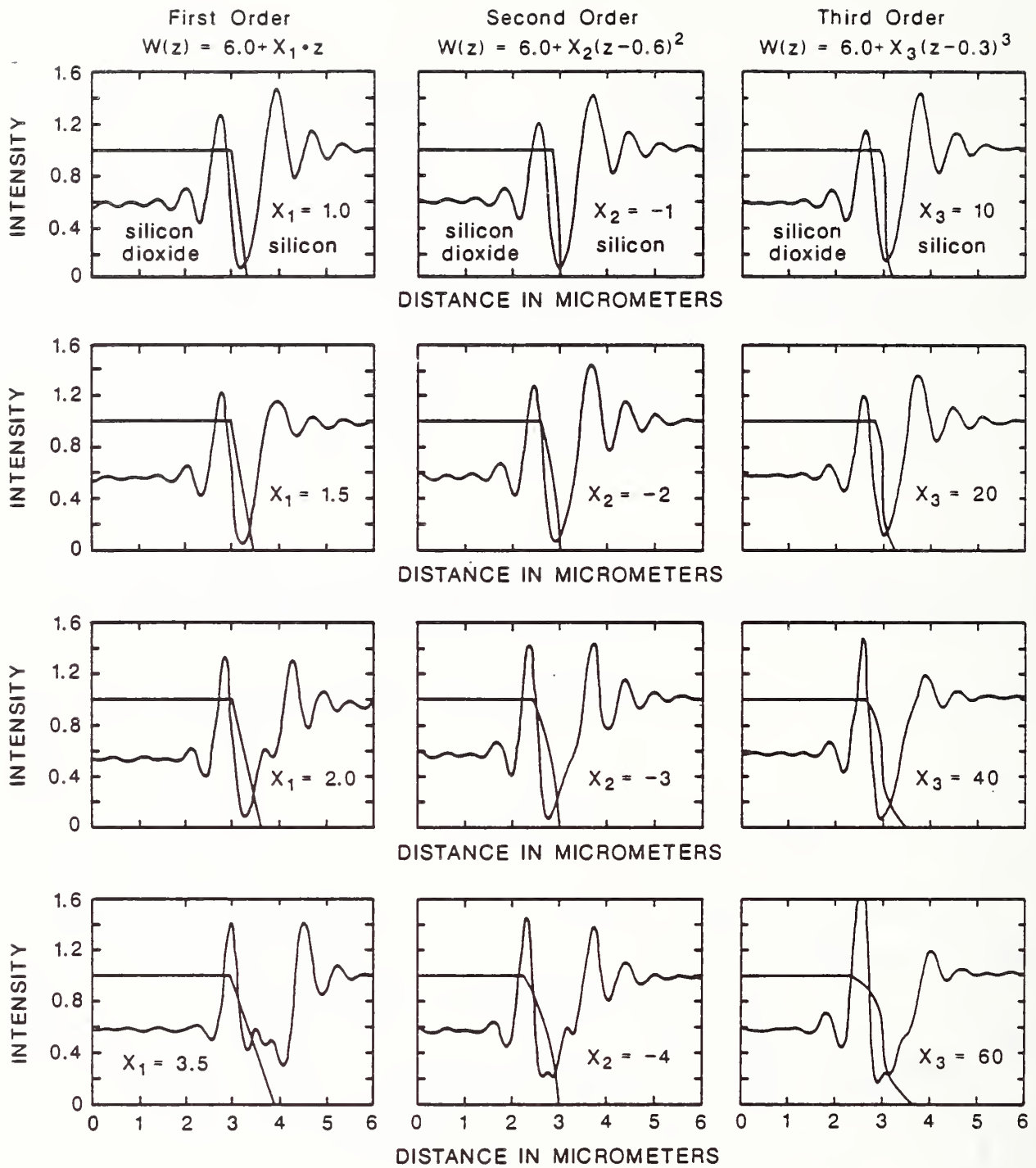


Fig. 5. The effect of edge geometry on the calculated optical image profile. The edge geometry has been superimposed on the image profile for reference.

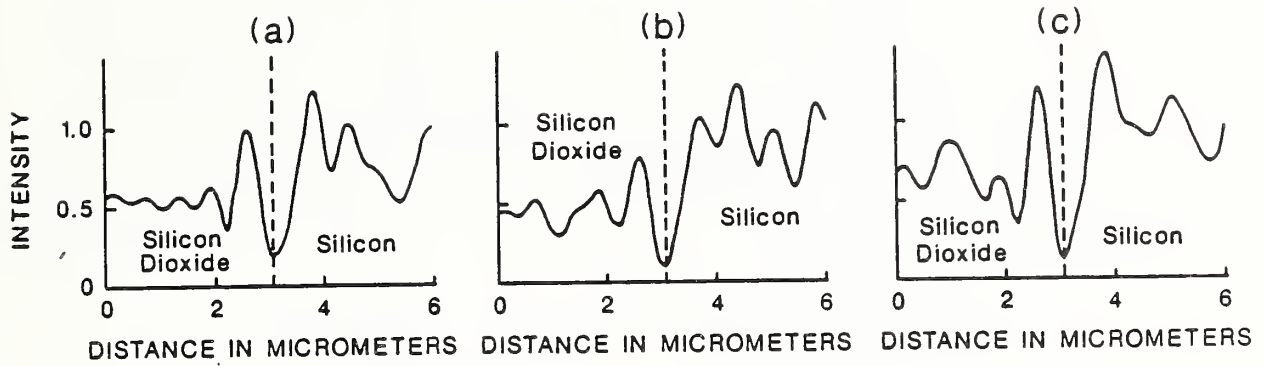


Fig. 6. Theoretical image intensity profiles of a line patterned in a noisy dielectric layer on silicon. A different set of noise data has been used for each curve.

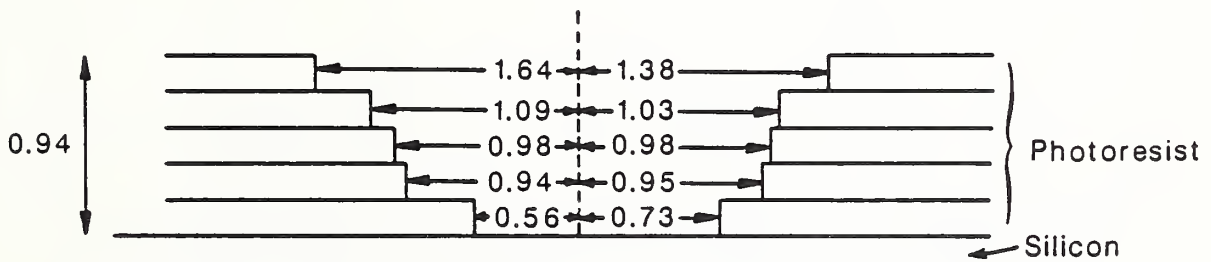


Fig. 7. Physical profile model of the asymmetric window structure in resist used to compute the image profiles in Fig. 8. (Dimensions in micrometers.)

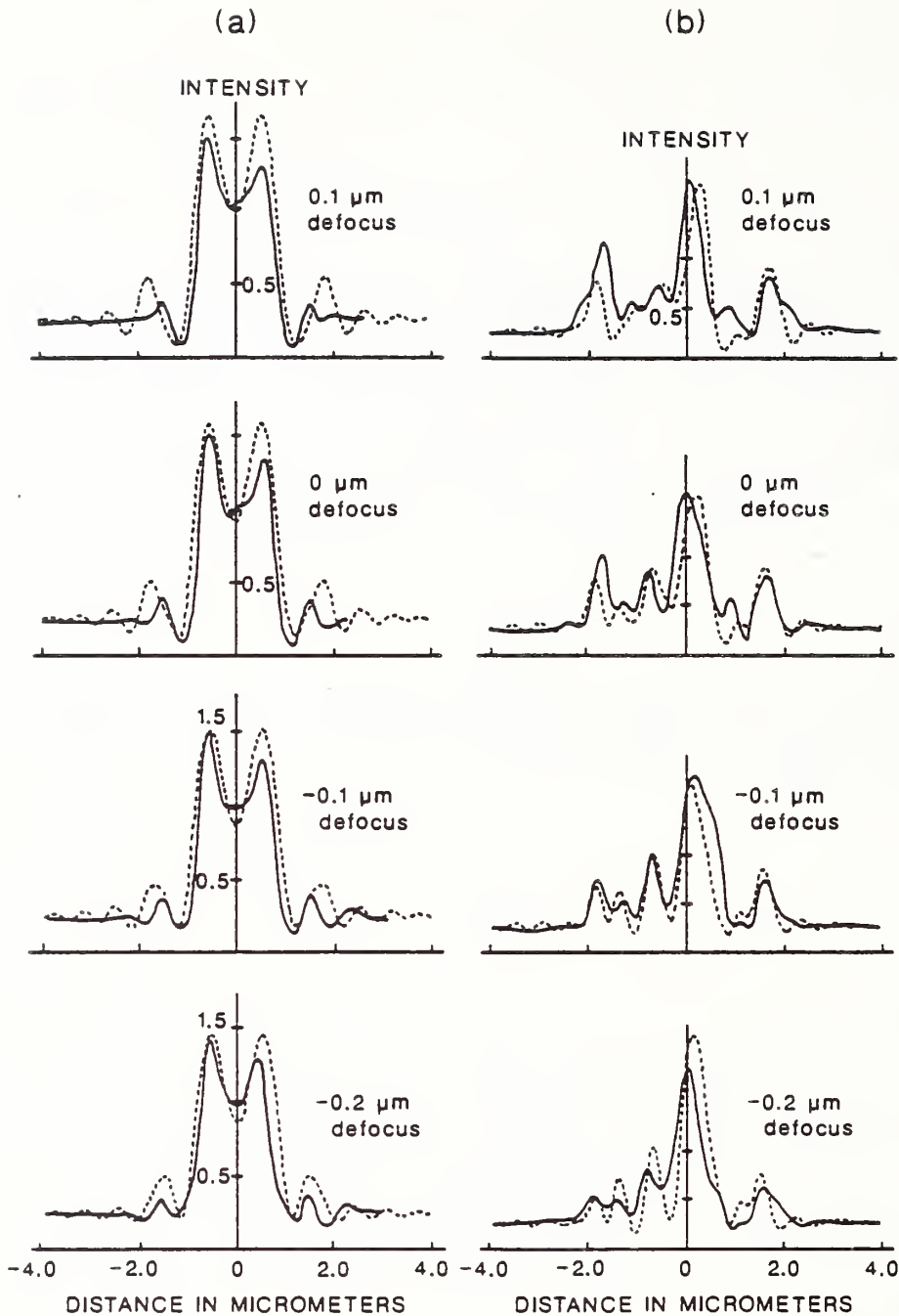


Fig. 8. Comparison of experimental and theoretical image intensity profiles of a window in photoresist on silicon. The window has either vertical edge walls (a) or sloping edges (b). The experimental profiles are shown by the thick line and the theoretical profiles by the thin line. The photoresist was assumed to have a thickness of  $0.94 \mu\text{m}$  and the geometry shown in Fig. 7 was used to model the window in the case of sloping edges.



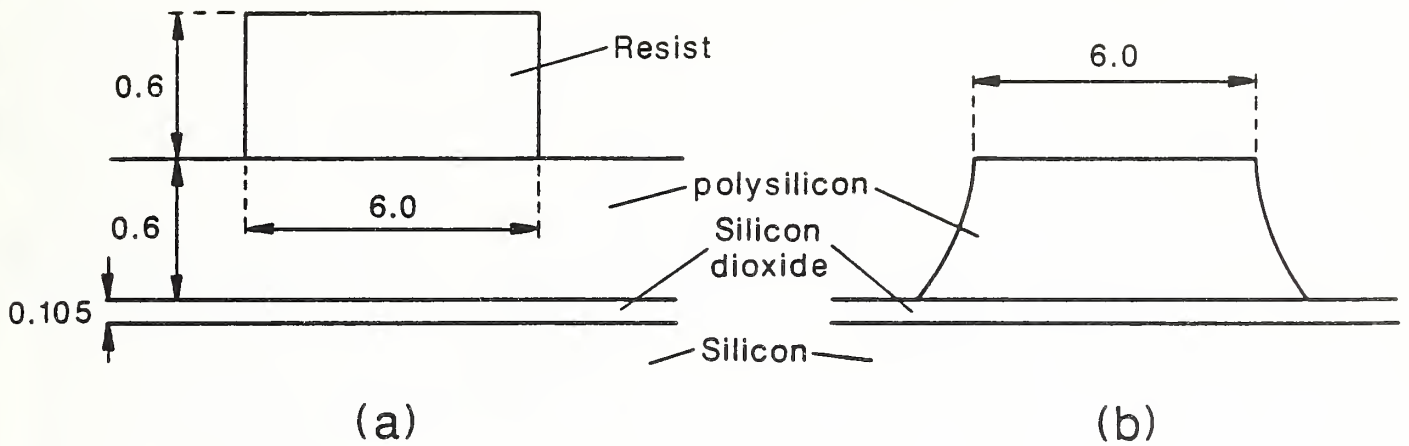


Fig. 9. Cross sections of the shapes used to model the polysilicon patterning stage of making an MOS transistor. The two shapes are patterned resist in unetched polysilicon (a) and etched polysilicon with the resist removed (b). (Dimensions in micrometers.)

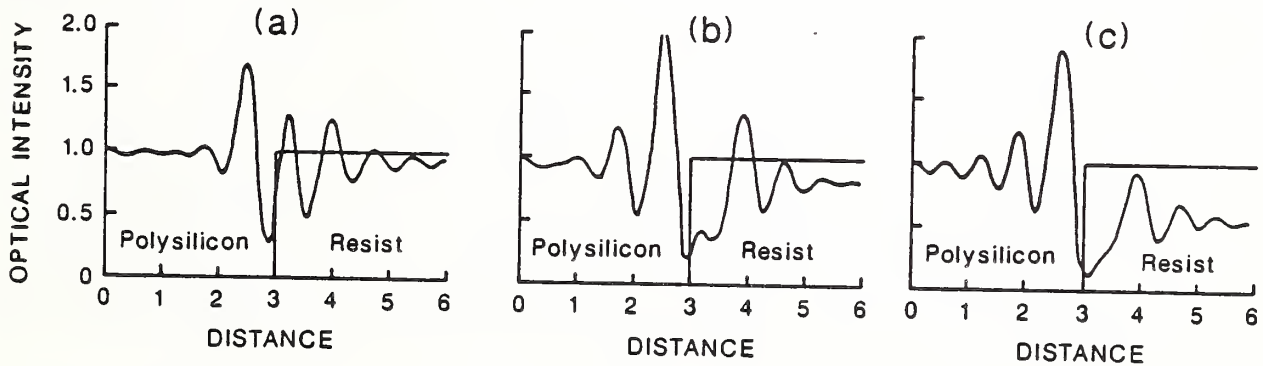


Fig. 10. Theoretical image profiles of a resist line on polysilicon (structure in Fig. 9a) for polysilicon thicknesses of a) 0.5, b) 0.6, and c) 0.7  $\mu\text{m}$ . A vertical edge is assumed.

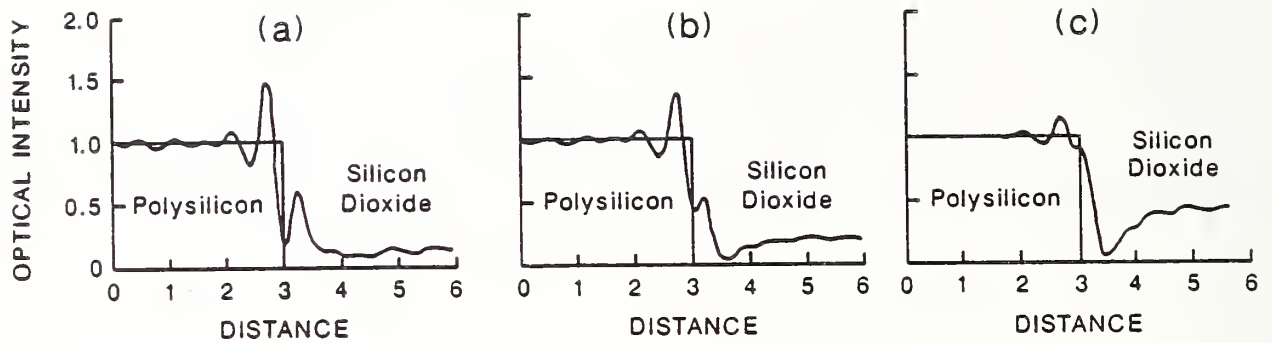


Fig. 11. Theoretical image profiles of a polysilicon line (structure in Fig. 9b) for oxide thicknesses of a) 85, b) 105, and c) 125 nm. A vertical edge is assumed and 0.6  $\mu\text{m}$  polysilicon thickness.

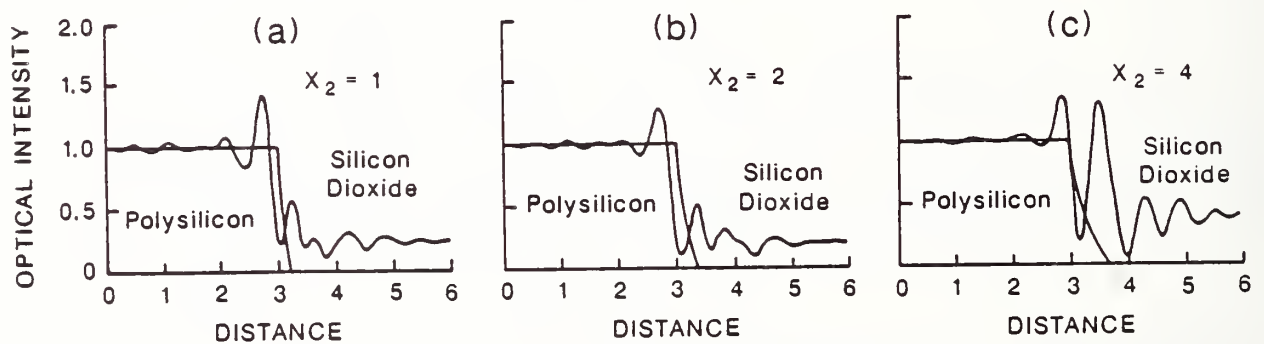


Fig. 12. Theoretical image profiles of a polysilicon line (structure in Fig. 9b) for a range of edge curvatures defined by Eq. 26. Oxide thickness is 105 nm; polysilicon thickness is 0.6  $\mu\text{m}$ .

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8	7	11	05. Software title			<input checked="" type="checkbox"/>	<input type="checkbox"/>	<input type="checkbox"/>						
04. Software date			THKIMAG - Computer software for the computation of optical microscope images of line objects patterned in thick layers viewed in a coherent system.											
Yr.	Mo.	Day	07. Internal Software ID											
8	7	08	06. Short title THKIMAG											
08. Software type			09. Processing mode			10. Application area								
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13. Narrative														
This computer software calculates the optical microscope images of line objects patterned in thick layers (more than one-quarter of the illuminating wavelength thick). The algorithms used are based on a monochromatic, waveguide model which can predict the image of line objects with arbitrary cross-section geometry including multilayer structures with sloped, curved, asymmetric, and undercut edges. Along with the computer software, test cases for help in implementation are given. As written, the program uses subroutines for complex matrix eigenvalues and eigenvectors from the NAG library.														
14. Keywords														
diffraction; dimensional metrology; linewidth; microscopy; optical imaging; optical metrology; waveguide														
15. Computer manufr and model			16. Computer operating system			17. Programing language(s)		18. Number of source program statements						
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19. Computer memory requirements			20. Tape drives			21. Disk/Drum units		22. Terminals						
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