

# Fundamental Energy Levels of Neutral Promethium (Pm I)\*

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The spectrum of atomic promethium has been observed with a variety of light sources and spectrographs. The Zeeman effect has also been recorded. Analysis of the spectrum shows that the ground configuration of the neutral promethium atom is  $4f^36s^2$ . The relative positions (in  $\text{cm}^{-1}$ ) of the low levels of this configuration are:

${}^6\text{H}_{3/2}^{\circ}$	0.00	${}^6\text{H}_{13/2}^{\circ}$	3919.03	${}^6\text{F}_{3/2}^{\circ}$	5872.84
${}^6\text{H}_{7/2}^{\circ}$	803.82	${}^6\text{H}_{15/2}^{\circ}$	5089.79	${}^6\text{F}_{7/2}^{\circ}$	6562.86
${}^6\text{H}_{9/2}^{\circ}$	1748.78	${}^6\text{F}_{1/2}^{\circ}$	5249.48	${}^6\text{F}_{5/2}^{\circ}$	7497.99
${}^6\text{H}_{11/2}^{\circ}$	2797.10	${}^6\text{F}_{3/2}^{\circ}$	5460.50	${}^6\text{F}_{11/2}^{\circ}$	8609.21

This group represents all levels of  $4f^36s^2$  expected below  $14,000 \text{ cm}^{-1}$ . From these results the following values of interaction parameters and their estimated uncertainties have been inferred:

$$\zeta_{4f} = 925 \pm 20 \text{ cm}^{-1} \quad E^3 = 510 \pm 20 \text{ cm}^{-1}$$

Data on 209 upper levels of even parity and 714 classified lines are given.

Key Words: Atomic spectroscopy, electronic energy levels, neutral atom, promethium, rare earth.

## 1. Introduction

In this paper we present the first results of our investigation to determine the electronic structure of neutral promethium.

Promethium was the last lanthanon element to be discovered. Since the early 1900's it was known that an element lying between neodymium and samarium with atomic number 61 remained to be discovered. A large number of attempts were made to find this element in nature without success. Two false claims of discovery resulted in the early names of illinium and florentium for element 61. The long sought-for element was finally identified in 1947, when Marinsky, Glendenin, and Coryell made a chemical separation of a new element from among the fission products of uranium. They chose the name *promethium* for element 61 ". . . after Prometheus, the Titan in Greek mythology who stole fire from heaven for the use of mankind." This name was chosen not only to symbolize the dramatic way in which the element could be produced in quantity as a result of man's harnessing of the energy of nuclear fission, but also to ". . . warn man of the impending danger of punishment by the vulture of war," [1].<sup>1</sup>

The first extensive work on the spectrum of Pm was carried out at the National Bureau of Standards by Meggers, Scribner, and Bozman [2], who published a list of 2249 Pm lines excited in arcs and sparks. Their attempts to separate the lines according to ionization state and to find spectral regularities were not successful. Their paper gives many of the details of the history of the discovery of Pm and of earlier spectroscopic studies.

Promethium is an entirely artificial element. All of its isotopes are radioactive, and none has ever been found in nature. In the several investigations of the promethium spectrum to date, including the present one, the samples have consisted of monoisotopic  $\text{Pm}^{147}$ . This is the only readily available form of Pm. It is produced in nuclear reactors as 2.6 percent of the fission products of uranium.  $\text{Pm}^{147}$  decays by  $\beta^-$  emission (0.22 MeV) into  $\text{Sm}^{147}$  with a half-life of 2.6 years.

The original paper of Meggers, Scribner, and Bozman suggested the existence of appreciable hyperfine structure in a number of  $\text{Pm}^{147}$  spectrum lines. This hfs was investigated by Klinkenberg and Tomkins [3] with a 9-meter grating spectrograph and later by Reader and Davis [4, 5] with Fabry-Perot interferometers. As a result of these investigations and others by the methods of paramagnetic resonance [6] and atomic beam resonance [7], it is known that the  $\text{Pm}^{147}$  nucleus has a spin  $I=7/2$ , a magnetic dipole moment  $\mu_I=2.6 \text{ nm}$ , and an electric quadrupole moment

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<sup>1</sup> Figures in brackets indicate literature references at the end of this paper.

$Q \sim 0.7$  barn. These nuclear moments cause many Pm lines to appear with very broad and complex hyperfine structure which impedes attempts to describe the spectrum.

The ionization energy of the neutral promethium atom has been estimated [8] by an interpolation method to be  $5.55 \pm 0.02$  eV.

## 2. Experimental Procedure

### 2.1. Light Sources

The construction of light sources was undertaken with the greatest care, because Pm, being radioactive, cannot be handled in the ordinary way. The first light source was a cooled hollow cathode, made for the investigation of hyperfine structure. Its use was essential for that work [4], but it was not used extensively in the present investigation.

The more useful sources were electrodeless discharge tubes, constructed especially for us by Earl Worden at the Lawrence Radiation Laboratory, Livermore. The promethium was obtained from Oak Ridge National Laboratory, and purified either there or at the Lawrence Radiation Laboratory, Berkeley. A spectrographic analysis was made prior to its use. A typical tube was constructed of a section of fused silica tubing 2 cm long with 7 mm o.d. and 5 mm i.d., containing 200  $\mu\text{g}$  of  $\text{PmI}_3$ . A long handle was attached. No filling gas was admitted, since an initial heating of the tube released enough iodine to make subsequent starting of the discharge relatively easy.

Altogether, six of these tubes were used. Each tube was operated inside a microwave cavity, supplied with rf power by a magnetron at a frequency of 2450 MHz. The temperature of the tube (and hence intensity of the spectrum) was regulated by adjustment of power to the cavity and the amount of air-cooling. During operation, the discharge was a brilliant blue color.

### 2.2. Spectrographs

Our experimental investigation of the spectrum of promethium has extended over several years, and has included the taking of hundreds of spectrograms on several instruments, under many different conditions. As we gained experience and familiarity with the spectrum, each set of spectrograms was taken to provide specific information.

The 6.4-meter concave grating at Berkeley was used for preliminary testing of the light sources and identification of spectrum lines. The instrument has a plate factor of 1.25  $\text{\AA}/\text{mm}$  at 5000  $\text{\AA}$  in the first order.

The 3-meter Czerny-Turner plane-grating spectrograph at Berkeley was used for accurate wavelength measurements. This instrument has a grating of width 12.5 cm, ruled with 300 grooves/mm. It is used at angles of incidence and diffraction of approximately 64 deg. At 5000  $\text{\AA}$  the instrument is used in the 12th order, and the plate factor is 0.36  $\text{\AA}/\text{mm}$ .

A 3.4-meter Ebert plane-grating spectrograph at the Lawrence Radiation Laboratory, Berkeley, was used

for some of the Zeeman spectrograms. The angles of use, orders, and plate factor of this instrument are about the same as for the 3-meter Czerny-Turner spectrograph just described.

The 9-meter concave grating spectrograph at the Argonne National Laboratory [9, 10] was used for the early spectrograms taken for wavelength measurements, temperature classification, and Zeeman effects. Spectrum lines were observed in orders as high as the eighth, although we generally confined our measurements to the third and fourth, in order to avoid overlapping orders. At these lower orders, overlapping could be prevented through the use of filters. The plate factor is about 0.45  $\text{\AA}/\text{mm}$  at 4000  $\text{\AA}$  in the fourth order.

### 2.3. Exposures

Three different magnets were used for the Zeeman spectrograms. Most of them were taken at Argonne, by using an electromagnet with iron pole-pieces, producing a field of 24,000 G. The discharge tube was excited in the cavity placed between the poles, perpendicular to the field.

The electromagnet used at the Radiation Laboratory produced a field of 29,000 G. A second magnet used there was a superconducting solenoid [11] operated at a field of 41,000 G. In this case, the source was operated in the cavity, but aligned parallel to the magnetic field. A more stable operation of the discharge resulted with this orientation. A small mirror placed inside the solenoid permitted observation of the light emitted in a direction perpendicular to the field.

### 2.4. Exposures

The exposures made at Argonne were taken on  $5 \times 45$  cm plates, with nine separate tracks on each plate. Thorium standard lines were placed on the top and bottom, with seven promethium exposures in-between. Each Pm exposure was taken at a different lamp temperature to help separate the spectra [12]. The intensities differed by as much as a factor of 80 from the weakest to the strongest exposures. These plates proved to be very valuable for spectrum separation, but the presence of small unexplained shifts made them less useful for absolute wavelength measurements. Since the wavelengths of most of the classified lines given in this paper are based on these exposures, it is possible that the present wavelengths may eventually have to be revised by a few thousandths of an angstrom. The Zeeman exposures included pi, sigma, and no-field lines. The field was calibrated by means of patterns belonging to several lines of Ca, Ag, and Cu which appeared on the plates.

The spectrograms taken at Berkeley for wavelength measurements were made on  $10 \times 25$  cm plates, with the thorium and promethium lines overlapping, to eliminate any errors of measurement. A samarium spectrum was also included on these plates to identify impurity lines due to small amounts of Sm formed by the decay of Pm. Overlapping grating orders were

TABLE 1. *Theoretical predictions for the low-lying levels of the  $4f^5 6s^2$  configuration of Pm I.*<sup>a</sup>

Level	Calculated energy (cm <sup>-1</sup> )	Interval (cm <sup>-1</sup> )	Calculated <i>g</i> -value	Percentage composition
<sup>6</sup> H <sub>5/2</sub> <sup>o</sup>	0	816	0.297	96% <sup>6</sup> H
<sup>6</sup> H <sub>7/2</sub> <sup>o</sup>	816	953	0.829	97% <sup>6</sup> H
<sup>6</sup> H <sub>9/2</sub> <sup>o</sup>	1769	1046	1.071	98% <sup>6</sup> H
<sup>6</sup> H <sub>11/2</sub> <sup>o</sup>	2815	1109	1.201	99% <sup>6</sup> H
<sup>6</sup> H <sub>13/2</sub> <sup>o</sup>	3924	1146	1.278	99% <sup>6</sup> H
<sup>6</sup> H <sub>15/2</sub> <sup>o</sup>	5070		1.327	95% <sup>6</sup> H
<sup>6</sup> F <sub>1/2</sub> <sup>o</sup>	4915	182	-0.649	98% <sup>6</sup> F
<sup>6</sup> F <sub>3/2</sub> <sup>o</sup>	5097	381	1.057	
<sup>6</sup> F <sub>5/2</sub> <sup>o</sup>	5478	675	1.303	
<sup>6</sup> F <sub>7/2</sub> <sup>o</sup>	6153	924	1.389	
<sup>6</sup> F <sub>9/2</sub> <sup>o</sup>	7077	1094	1.429	
<sup>6</sup> F <sub>11/2</sub> <sup>o</sup>	8171		1.451	

<sup>a</sup> J. G. Conway and B. G. Wybourne, Phys. Rev. **130**, 2325 (1963).

separated by use of an external prism predisperser [13], which prevented light from the unwanted orders from entering the spectrograph. Each grating order was photographed separately.

Eastman Kodak spectroscopic plates were used for all exposures, processed in the recommended manner. Types 103a-O, 103a-F, 103-O, 103-F, and I-N were utilized. All but the first two were prefogged by exposures to weak light, prior to loading in the spectrograph.

### 2.5. Measurements

The plates were measured on two semiautomatic scanning comparators, one at the University of California and one at the National Bureau of Standards. Wavelengths for the thorium standards were taken from the work of Meggers and Stanley [14]. The estimated uncertainty of the wavelength measurements is about  $\pm 0.005$  Å. At the present time some of our plates from Argonne and most of the plane-grating plates remain to be measured. When these measurements are complete, we expect to provide a full description of the Pm spectrum.

### 3. Analysis

In the neutral rare earths the ground configurations are generally of the type  $4f^N 6s^2$  or  $4f^{N-1} 5d 6s^2$ . Since the ground configurations in Nd I and Sm I are  $4f^4 6s^2$  and  $4f^6 6s^2$ , respectively, it was expected that the ground

configuration in neutral Pm would be  $4f^5 6s^2$ . The lowest level of  $4f^4 5d 6s^2$  was not expected to be less than 10,000 cm<sup>-1</sup> above the lowest level of  $4f^5 6s^2$ . The levels of  $4f^5 5d 6s$  were expected to start at about 8500 cm<sup>-1</sup> above the lowest level of  $4f^5 6s^2$ . Prior to our beginning the Pm I analysis, Conway and Wybourne [15] had published theoretical predictions for the relative energies and *g*-values of the low levels of  $4f^5 6s^2$  by using hydrogenic ratios for the Slater parameters and interpolating values of  $\zeta/F_2$ . These predictions proved to be very useful to us in carrying out the analysis. Their results for Pm I are summarized in table 1.

The first part of the analysis was carried out through the use of the Zeeman data. A large number of self-reversed lines showed resolved patterns involving levels having  $J=5/2$  and  $7/2$  and *g*-values of about 0.30 and 0.83, respectively. It was clear that these lines were transitions to the  $4f^5 6s^2$  <sup>6</sup>H<sub>5/2</sub><sup>o</sup> and <sup>6</sup>H<sub>7/2</sub><sup>o</sup> levels. After a certain amount of trial and error in searching for repeating differences involving these lines, the <sup>6</sup>H<sub>5/2</sub><sup>o</sup> - <sup>6</sup>H<sub>7/2</sub><sup>o</sup> interval was found to be 803.82 cm<sup>-1</sup>. About 15 upper even levels were established in this process. The good agreement between this result and the prediction of Conway and Wybourne showed that the predictions would be useful in extending the analysis.

By continuing to use the Zeeman data and by using an electronic computer to search for constant differences between groups of strong lines, two sets of classified lines were built up. One group represented transitions to the  $4f^5 6s^2$  <sup>6</sup>H<sub>5/2-13/2</sub><sup>o</sup> group of levels; the other represented transitions to the  $4f^5 6s^2$  <sup>6</sup>F<sub>1/2-7/2</sub><sup>o</sup> levels. However the connection between these two groups was not known. After a considerable amount of searching for this connection, a few weak, resolved

TABLE 2. *Low levels of the  $4f^5 6s^2$  configuration of Pm I*

Level	Energy (cm <sup>-1</sup> )	Interval (cm <sup>-1</sup> )	<i>g</i> -value
<sup>6</sup> H <sub>5/2</sub> <sup>o</sup>	0.00		0.305
<sup>6</sup> H <sub>7/2</sub> <sup>o</sup>	803.82	803.82	0.831
<sup>6</sup> H <sub>9/2</sub> <sup>o</sup>	1748.78	944.96	1.079
<sup>6</sup> H <sub>11/2</sub> <sup>o</sup>	2797.10	1048.32	1.205
<sup>6</sup> H <sub>13/2</sub> <sup>o</sup>	3919.03	1121.93	1.307
<sup>6</sup> H <sub>15/2</sub> <sup>o</sup>	5089.79	1170.76	1.33
<sup>6</sup> F <sub>1/2</sub> <sup>o</sup>	5249.48		-0.68
<sup>6</sup> F <sub>3/2</sub> <sup>o</sup>	5460.50	211.02	1.051
<sup>6</sup> F <sub>5/2</sub> <sup>o</sup>	5872.84	412.34	1.293
<sup>6</sup> F <sub>7/2</sub> <sup>o</sup>	6562.86	690.02	1.385
<sup>6</sup> F <sub>9/2</sub> <sup>o</sup>	7497.99	935.13	1.440
<sup>6</sup> F <sub>11/2</sub> <sup>o</sup>	8609.21	1111.22	1.458

Zeeman patterns were found in the ultraviolet which proved to be transitions from upper levels known from combinations with  ${}^6F^\circ$  levels down to the low  ${}^6H^\circ$  levels. This established the energies of the  ${}^6F^\circ$  levels relative to the ground state. The  ${}^6H_{15/2}^\circ$ ,  ${}^6F_{9/2}^\circ$ , and  ${}^6F_{11/2}^\circ$  levels were later established, thus bringing our knowledge of the low levels to its present state. According to a recent diagonalization by Conway [16] this group represents all levels of  $4f^56s^2$  expected below  $14000\text{ cm}^{-1}$ .

The results are given in tables 2, 3, and 4. Table 2 gives the energies and  $g$ -values of the odd levels. Table 3 gives the energies,  $J$ -values and  $g$ -values for the even levels. The classified lines are given in table 4. The estimated uncertainty in the values of the low levels given in table 2 is  $\pm 0.01\text{ cm}^{-1}$ . The estimated uncertainty in the values of the high levels given in table 3 is about  $\pm 0.03\text{ cm}^{-1}$ .

#### 4. Discussion

The total number of self-reversed lines in Pm I is 122. Of these 120 have been classified as transitions to  $4f^56s^2$ . This makes it certain that the ground configuration of Pm I is  $4f^56s^2$ .

The eigenvectors given by Conway and Wybourne as a result of their diagonalization of  $4f^5$  show the  ${}^6H^\circ$  and  ${}^6F^\circ$  terms to be nearly pure in  $LS$  coupling. We therefore would have expected a somewhat regular variation in intensity of the lines making transitions from a single upper level to several lower levels. However, according to our observations this is not the case. Figure 1 gives some of the more striking examples of the irregular intensities. Of special note are the lines from 26725.52(3/2). In this case the transition to  ${}^6F_{3/2}^\circ$  is just barely visible on the plates. The complete absence of the line from 28186.31(7/2) to  ${}^6H_{7/2}^\circ$  is also very striking. A similar set of puzzling intensities has been observed by Shenstone [17] in the  $3d^64s-4p$  transitions of Co III. Here the anomalous intensities were found in groups of lines connecting terms of different multiplicity. In this connection Shenstone noted "Especially difficult to understand is the not

uncommon habit of intersystem combinations of missing the central of three levels of successive  $J$ ." Although this phenomenon has not yet been investigated theoretically, it is clear that it stems from the lack of pure  $LS$  coupling in the upper configurations. It is likely that a theoretical study of the transition probabilities for the  $3d^64s-3d^64p$  array in Co III would shed more light on this problem.

No real effort has been made yet to understand the origin of the known upper levels in Pm I. However we note the following points. There are only two configurations which can make transitions to  $4f^56s^2$  with appreciable intensity, namely  $4f^56s6p$  and  $4f^45d6s^2$ . For the  $4f^56s6p$  configuration, the strongest transitions to the  $4f^56s^2\ {}^6H^\circ$  and  ${}^6F^\circ$  levels will originate from levels of the type  $4f^5({}^6H)6s6p$  and  $4f^5({}^6F)6s6p$ . If we consider the levels of the type  $4f^5({}^6H)6s6p(J=5/2)$ , we would expect them to fall into two groups: six of the type  $4f^5({}^6H)+6s6p({}^3P)$  and two of the type  $4f^5({}^6H)+6s6p({}^1P)$ . This type of coupling (first described by Shenstone [18] in the case of the  $3d^64s4p$  configuration of Cu I) will hold approximately here because the parameter which determines the  ${}^1P-{}^3P$  splitting,  $G_1(sp)$ , is expected by interpolation from other rare earths to be about  $2640\text{ cm}^{-1}$ , whereas  $\zeta_p$  and  $\zeta_f$  are only about  $1000\text{ cm}^{-1}$ . The  $f$ - $s$  and  $f$ - $p$  interactions are much smaller (see Smith and Wybourne's treatment of the  $4f^7({}^8S)6s6p$  configuration in Eu I [19]) and for our purpose may be neglected. If one interpolates a value for the  $4f^5({}^6H)6s^2-4f^5({}^6H)6s6p$  energy difference and uses the above parameters to estimate the level positions, one finds that the six  ${}^3P$  type levels with  $J=5/2$  will lie in the region  $14000-16000\text{ cm}^{-1}$  and the two  ${}^1P$  type levels with  $J=5/2$  will lie at about  $20000-21000\text{ cm}^{-1}$  [20]. Of the 27 observed upper levels with  $J=5/2$ , 22 make strong transitions to the  ${}^6H$  group of lower levels. These 22 levels are distributed as follows: 14 between  $20250$  and  $23550\text{ cm}^{-1}$ , 8 between  $26600$  and  $28350\text{ cm}^{-1}$ . Thus of the observed levels with  $J=5/2$ , only 2 would be expected to belong to  $4f^5({}^6H)6s6p$ . It should also be noted that none of the lines classified so far in Pm I shows the appreciable hyperfine structure which would be expected if one of the configurations contained a single  $6s$  electron. Judd [21] has shown that it is possible for lines from certain levels of  $4f^N6s6p$  configurations to  $4f^N6s^2$  to show no hyperfine structure. However, there are too many levels here whose transitions show no  $hfs$  to believe that the theory is applicable here. The absence of  $hfs$  in these lines more likely is evidence of a closed  $6s$  shell in both upper and lower configurations. In this case the hyperfine structures due to the  $4f$  electron in the upper and lower levels have the same sign and about the same magnitude, so that the observed line shows no resolved structure. For these reasons we believe that most of the known upper levels belong to the  $4f^45d6s^2$  configuration.

Comparison of the observed positions of the  $4f^56s^2\ {}^6H^\circ$  and  ${}^6F^\circ$  levels with the calculations of Conway and Wybourne shows that their predictions for the intervals within each term are very good, generally within  $\pm 30\text{ cm}^{-1}$  of the observed levels. However, the predicted positions of the  ${}^6F^\circ$  levels are too low

EVEN \ ODD	${}^6H^\circ$						${}^6F^\circ$				
	5/2	7/2	9/2	11/2	13/2	15/2	1/2	3/2	5/2	7/2	9/2
21 348.22 (7/2)	700	100	400								
22 294.96 (9/2)		X	500	400							
26 181.98 (13/2)				3	300	200					
26 725.52 (3/2)	300						500	1	50		
28 169.71 (7/2)	X	300	100						X	X	X
28 186.31 (7/2)	200	X	200						X	500	50
29 002.94 (3/2)	X						35	8	150		

FIGURE 1. Anomalous intensities in Pm I.  
No observed transition = X.



by nearly 7 percent in every case. This discrepancy results from the fact that the calculated intervals within the individual terms are very sensitive to the value of the spin-orbit parameter, but rather insensitive to the electrostatic parameters. On the other hand the separation between the barycenters of terms is governed primarily by the electrostatic parameters. At the time of Conway and Wybourne's work there was no neutral rare earth with more than one known term of  $4f^N 6s^2$ . Therefore, information about the electrostatic parameters could be inferred from the known levels only through second order effects. Thus, a difference of only 7 percent between the predicted and observed  ${}^6\text{H}^\circ$ - ${}^6\text{F}^\circ$  separation in Pm I can be considered to be fairly good agreement.

It is not possible to obtain reliable values for the electrostatic parameters  $E^1$ ,  $E^2$ , and  $E^3$  from the known levels, because only one term separation is available. However, we note that since the  ${}^6\text{H}^\circ$  and  ${}^6\text{F}^\circ$  terms are nearly pure in  $LS$  coupling, to a first approximation the energy difference between the  ${}^6\text{H}^\circ$  and  ${}^6\text{F}^\circ$  barycenters  $\Delta E({}^6\text{F}, {}^6\text{H})$  will be equal to  $9E^3$  [22]. If we include the Trees  $\alpha L(L+1)$  correction [23], to a very good first approximation we then have:

$$\Delta E({}^6\text{F}, {}^6\text{H}) = 9E^3 - 18\alpha.$$

If we use the known positions of the  ${}^6\text{F}^\circ$  and  ${}^6\text{H}^\circ$  levels to determine  $\Delta E({}^6\text{F}, {}^6\text{H})$ , we find  $\Delta E({}^6\text{F}, {}^6\text{H}) = 4167.92 \text{ cm}^{-1}$ . If we set  $\alpha = 30 \text{ cm}^{-1}$  as indicated by the theoretical interpretation of the spectra of Ce III [24] and Pr III [25] we find  $E^3 \sim 520 \text{ cm}^{-1}$ .

A preliminary value of  $\zeta_{4f}$  can be obtained by considering the total widths of the  ${}^6\text{H}^\circ$  and  ${}^6\text{F}^\circ$  terms. To a first approximation the sum of these two widths is  $91\zeta_{4f}$ , which gives  $\zeta_{4f} \sim 930 \text{ cm}^{-1}$ .

Crosswhite [26] has made a least squares fit of the  ${}^6\text{F}^\circ$  and  ${}^6\text{H}^\circ$  levels to the theoretical energy formulas by using hydrogenic ratios for the electrostatic parameters and a fixed value of  $20 \text{ cm}^{-1}$  for  $\alpha$ . The parameters  $E^3$  and  $\zeta_{4f}$  were allowed to vary. A diagonalization with  $E^3 = 510 \text{ cm}^{-1}$  and  $\zeta_{4f} = 914 \text{ cm}^{-1}$  gave a mean error of  $25 \text{ cm}^{-1}$ . This could be reduced to about  $8 \text{ cm}^{-1}$  if slightly different values of  $\zeta_{4f}$  were used for the two terms:  $912 \text{ cm}^{-1}$  for  ${}^6\text{H}$  and  $938 \text{ cm}^{-1}$  for  ${}^6\text{F}$ . Crosswhite notes that this is probably caused by a spin-other-orbit interaction. The  $J$ -dependence of this interaction is the same as that of the spin-orbit interaction. However, the spin-other-orbit interaction constant varies from term to term. Thus, this interaction will cause the spin-orbit constants derived from different terms of a configuration to appear to be slightly different. This effect was first treated for  $l^N$  configurations by Horie [27]; the principal aspects of the theory have been summarized by Wybourne [28]. Since spin-other-orbit effects cannot be observed by studying only one term of a configuration, the present results for Pm I provide a first opportunity to view their magnitude for the neutral rare earths.

The residual errors in the above calculation have a form very close to that expected from neglect of a spin-spin interaction. If an estimate of the spin-spin

interaction energy is made by using Judd's [29] matrix elements and interpolated values of the radial integrals from the calculations of Blume, Freeman, and Watson [30], the mean error can be further reduced to about  $2 \text{ cm}^{-1}$ . When this is done the values of  $\zeta_{4f}$  are changed to  $910 \text{ cm}^{-1}$  for  ${}^6\text{H}^\circ$  and  $940 \text{ cm}^{-1}$  for  ${}^6\text{F}^\circ$  [26].

In view of the uncertainties in the ratios of the electrostatic parameters, the value to be used for  $\alpha$ , and the Hamiltonian needed to describe the levels, we give the values of  $E^3$  and  $\zeta_{4f}$  for the  $4f^5 6s^2$  configuration of Pm I as:

$$\zeta_{4f} = 925 \pm 20 \text{ cm}^{-1}$$

$$E^3 = 510 \pm 20 \text{ cm}^{-1}$$

That this value of  $\zeta_{4f}$  fits in well with other values of  $\zeta_{4f}$  in the rare earths is shown by the plot in figure 2.

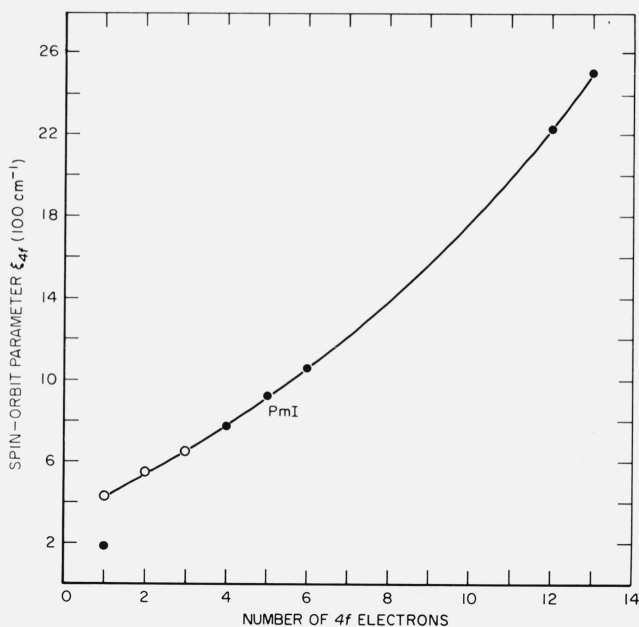


FIGURE 2. Values of  $\zeta_{4f}$  in the rare earths.

The solid circles represent values derived from  $4f^N 6s^2$  configurations of neutral atoms; the open circles are derived from  $4f^N$  configurations of doubly ionized atoms. References:  $N=1$ , La I, H. N. Russell and W. F. Meggers, J. Res. NBS **9**, 625 (1932);  $N=1$ , La III, J. Sugar and V. Kaufman, J. Opt. Soc. Am. **55**, 1283 (1965);  $N=2$ , Ce III, N. Spector, J. Opt. Soc. Am. **55**, 492 (1965);  $N=3$ , Pr III, R. Trees, J. Opt. Soc. Am. **54**, 651 (1964);  $N=4$ , Nd I, J. G. Conway and B. G. Wybourne, Phys. Rev. **130**, 2325 (1963);  $N=5$ , Pm I, this paper;  $N=6$ , Sm I, J. G. Conway and B. G. Wybourne, Phys. Rev. **130**, 2325 (1963);  $N=12$ , Er I, J. Reader, unpublished calculations based on data of L. C. Marquet and S. P. Davis, J. Opt. Soc. Am. **55**, 471 (1965). This value of  $\zeta_{4f}$  of  $2237 \text{ cm}^{-1}$  is nearly identical to the  $2236 \text{ cm}^{-1}$  value of  $\zeta_{4f}$  derived from the  $4f^{12} 6s$  configuration of Er II by Z. Goldschmidt, J. Opt. Soc. Am. **53**, 594 (1963);  $N=13$ , Tm I, W. F. Meggers, Rev. Mod. Phys. **14**, 96 (1942). For simplicity, several values of  $\zeta$  which have been published for  $4f^N 6s$  configurations of singly ionized rare earths have not been included.

TABLE 3. Even levels of Pm I. Levels with asterisk are uncertain.

Energy (cm <sup>-1</sup> )	<i>J</i>	<i>g</i>	Energy (cm <sup>-1</sup> )	<i>J</i>	<i>g</i>	Energy (cm <sup>-1</sup> )	<i>J</i>	<i>g</i>	Energy (cm <sup>-1</sup> )	<i>J</i>	<i>g</i>
17104.72	7/2	0.885	23443.79	5/2	0.784	26065.63	1/2	-0.36	28169.71	7/2	0.9
20006.04	3/2	0.068	23480.63	9/2	1.16	26080.99	17/2	1.29	28186.31	7/2	1.98
20157.85	7/2	0.503	23501.57	11/2	1.283	26096.75	17/2	1.25	28196.56	1/2	2.44
20265.98	5/2	0.527	23538.86	5/2	0.780	26101.28	13/2		28273.52	7/2	0.764
20517.96	5/2	0.659	23550.60	11/2	1.170	26103.56	11/2	1.2	28274.21	9/2	1.055
20567.76	5/2	0.910	23571.27	9/2	1.26	26181.98	13/2	1.55	28325.13	7/2	1.38
20660.00	7/2	1.114	23584.31	7/2	1.123	26211.44	7/2	0.715	28338.98	5/2	0.963
20675.81	5/2	1.075	23629.06	13/2	1.09	26237.84	9/2	1.43	28467.52	7/2	0.87
20909.00	7/2	0.929	23712.56	7/2	1.181	26282.20	13/2	0.95	28490.35	9/2	1.123
21100.10	7/2	1.319	23732.57	1/2	3.24	26285.02	7/2	0.877	28565.66	9/2	1.0
21143.06	7/2	0.977	23740.42	13/2		26300.30	1/2	0.955	28607.33	9/2	1.079
21237.49	9/2	0.841	23743.96	7/2	1.111	26456.26	13/2		28608.57	3/2	1.740
21348.22	7/2	0.815	23760.57	11/2		26468.80	1/2	-0.45	28657.02	1/2	0.29
21371.05	7/2	0.927	23926.91	3/2	1.754	26479.61	3/2	0.731	28680.26	7/2	0.88
21590.60	3/2	0.135	23938.76	13/2	1.044	26522.35	3/2	1.135	28994.90	11/2	1.30
21625.45	9/2	1.117	24013.29*	13/2	1.29	26545.85	17/2	1.14	29002.94	3/2	0.99
21657.89*	5/2	1.01	24038.82	11/2	1.08	26555.44	13/2	1.07	29074.03	9/2	
21666.80	7/2	0.696	24071.03	9/2	1.12	26591.40	13/2	1.10	29129.60	9/2	1.162
21732.93	9/2	1.137	24091.39	3/2	1.395	26609.39	5/2	0.56	29161.96	9/2	1.0
21920.49	9/2	0.986	24122.41	13/2		26630.56	5/2	0.32	29242.64	11/2	1.259
21946.12	3/2	-0.01	24180.80	13/2	1.28	26694.38	11/2	1.095	29585.21	11/2	1.2
21976.26	7/2	1.218	24204.37	7/2	0.670	26695.79	5/2		29595.58	9/2	1.04
22013.40	3/2	0.887	24234.42	9/2		26703.97	13/2	1.30	29648.42*	9/2	
22080.08	5/2	0.571	24245.66	13/2	1.14	26725.52	3/2	0.65	29705.77	11/2	
22084.65	7/2	0.858	24338.33	9/2	1.17	26830.74	3/2	0.794	29757.69	11/2	
22205.44	9/2	0.974	24418.44	7/2	1.124	26841.36	5/2	1.38	29784.08	9/2	
22259.21	1/2	-0.32	24443.15	13/2	1.17	26955.22	15/2		29856.72	11/2	1.10
22294.96	9/2	1.245	24443.57	9/2	1.037	27036.66	5/2	0.931	29883.87	7/2	1.199
22301.24	5/2	0.976	24471.10	7/2	0.83	27042.18	15/2		29908.90	9/2	
22309.94	7/2	0.850	24503.45	11/2	1.22	27109.75	15/2		29960.42*	11/2	1.18
22355.68	9/2	1.374	24520.23	13/2	1.34	27245.99	5/2	0.761	30008.40	11/2	1.28
22388.06	3/2	0.84	24533.27	11/2	1.183	27272.46	7/2	0.8	30063.62	9/2	1.224
22414.17	11/2	1.12	24558.56	9/2	1.39	27304.15*	9/2		30251.50*	11/2	
22425.58	5/2	0.83	24627.53	9/2	0.961	27319.28	15/2	1.44	30281.98	13/2	1.2
22446.20	11/2	1.531	24681.68	11/2	0.895	27334.48	7/2	1.274	30374.95	13/2	
22456.72	9/2	0.936	24705.25	15/2	1.102	27351.42	5/2	0.947	30457.44	11/2	1.226
22522.90	5/2	0.735	24754.58	9/2	0.888	27383.92	7/2	0.92	30541.28	13/2	1.226
22586.77	9/2	1.283	24770.04	13/2	1.17	27468.45	15/2		30726.26	11/2	
22654.34	7/2	0.84	24789.86	9/2	1.08	27476.28	3/2	0.21	30785.03	9/2	
22656.68	5/2	0.936	24884.90	11/2		27512.95	7/2	1.077	31103.24	9/2	
22761.33	11/2	1.296	24912.34	11/2	1.29	27596.27	5/2	0.913	31846.70	11/2	
22817.13	11/2	1.134	25104.27	11/2	1.15	27621.74	13/2	1.29	32022.32	9/2	
22905.24	5/2	1.16	25306.07	15/2	1.21	27685.89*	7/2		32435.06	7/2	
22934.70	7/2	1.237	25351.46	3/2	0.58	27829.89	3/2	2.049	33180.50	13/2	
23006.35	11/2		25357.24	11/2		27919.29	9/2	1.024		15/2	1.285
23033.95	11/2	1.0	25402.61	11/2	1.211	27923.37	3/2	1.31	33246.65	9/2	
23178.13	7/2	1.150	25405.29	13/2	1.034	27939.87	5/2	0.869			
23188.54	5/2	1.48	25448.28	9/2	1.14	28008.09	9/2	1.025			
23198.33	11/2	1.12	25474.46	15/2		28030.99	7/2				
23276.10	9/2	0.83	25521.55*	9/2	0.910	28075.94	13/2	1.39			
23278.90	7/2	1.0	25537.36	9/2	1.39	28084.28	9/2	1.155			
23334.10	5/2	0.571	25618.77	13/2	1.13	28086.21	5/2	0.904			
23337.53	7/2	1.257	25755.17	15/2	1.016	28150.73	11/2	1.150			
23345.07	11/2	1.323	25919.50	11/2	1.26	28153.69	7/2	1.071			
23435.40	9/2	1.00	26015.94	13/2			5/2	0.632			

TABLE 4. *Classified lines of Pm I*

The wavelengths are in air. Even levels are designated by the energy in  $\text{cm}^{-1}$  followed by the  $J$ -value in parenthesis. Intensities are visual estimates on a scale of 1 to 1000.

*C* - complex  
*I* - shaded to longer wavelengths  
*S* - shaded to shorter wavelengths  
*W* - wide  
*P* - perturbed by close line, but resolved  
*U* - unperturbed by close line, but unresolved  
*D* - double  
*B* - blend  
*H* - hazy  
*R1* - very widely reversed  
*R2* - widely reversed  
*R3* - moderately reversed  
*R4* - slightly reversed  
*R5* - barely detectably reversed  
 \* - classification in doubt

$\lambda$ (Å)	<i>I</i>	$\sigma$ ( $\text{cm}^{-1}$ )	Classification	$\lambda$ (Å)	<i>I</i>	$\sigma$ ( $\text{cm}^{-1}$ )	Classification
6420.171	10	15571.61	$6F_{11/2}^{\circ}$ - 24180.80 (13/2)*	5818.750	3 <i>P</i>	17181.06	$6F_{5/2}^{\circ}$ - 23743.96 (7/2)
6355.910	100 <i>P</i>	15729.04	$6F_{9/2}^{\circ}$ - 24338.33 (9/2)	5776.992	200	17305.25	$6F_{5/2}^{\circ}$ - 23178.13 (7/2)
6348.654	8	15747.02	$6F_{7/2}^{\circ}$ - 22309.94 (7/2)	5730.809	200	17444.70	$6F_{5/2}^{\circ}$ - 22905.24 (5/2)
6335.048	100	15780.84	$6F_{9/2}^{\circ}$ - 23278.90 (7/2)	5724.260	60	17464.66	$6F_{5/2}^{\circ}$ - 23337.53 (7/2)
6313.796	15	15833.95	$6F_{11/2}^{\circ}$ - 24443.15 (13/2)*	5710.075	20	17508.05	$6F_{7/2}^{\circ}$ - 24071.03 (9/2)*
6313.663	10	15834.29	$6F_{11/2}^{\circ}$ - 24443.57 (9/2)	5671.018	500	17628.62	$6F_{11/2}^{\circ}$ - 26237.84 (9/2)
6311.586	85	15839.50	$6F_{9/2}^{\circ}$ - 23337.53 (7/2)	5657.260	300	17671.50	$6F_{7/2}^{\circ}$ - 24234.42 (9/2)
6308.577	20	15847.05	$6F_{9/2}^{\circ}$ - 23345.07 (11/2)	5644.515	10	17711.40	$6F_{5/2}^{\circ}$ - 23584.31 (7/2)
6302.377	60	15862.64	$6F_{7/2}^{\circ}$ - 22425.58 (5/2)	5639.234	50	17727.98	$6F_{3/2}^{\circ}$ - 23188.54 (5/2)
6283.239	70	15910.96	$6F_{11/2}^{\circ}$ - 24520.23 (13/2)	5624.182	20	17775.43	$6F_{7/2}^{\circ}$ - 24338.33 (9/2)
6278.104	2	15923.97	$6F_{11/2}^{\circ}$ - 24533.27 (11/2)	5603.922	125 <i>L</i>	17839.69	$6F_{5/2}^{\circ}$ - 23712.56 (7/2)
6268.138	150	15949.29	$6F_{9/2}^{\circ}$ - 24558.56 (9/2)	5598.944	2	17855.55	$6F_{7/2}^{\circ}$ - 24418.44 (7/2)
6263.942	10	15959.97	$6F_{7/2}^{\circ}$ - 22522.90 (5/2)	5597.807	10 <i>H</i>	17859.18	$6F_{9/2}^{\circ}$ - 25357.24 (11/2)
6255.078	10	15982.59	$6F_{9/2}^{\circ}$ - 23480.63 (9/2)	5594.077	100	17871.09	$6F_{5/2}^{\circ}$ - 23743.96 (7/2)
6246.909	50	16003.49	$6F_{9/2}^{\circ}$ - 23501.57 (11/2)	5591.080	35 <i>C</i>	17880.67	$6F_{7/2}^{\circ}$ - 24443.57 (9/2)*
6220.118	10	16072.42	$6F_{11/2}^{\circ}$ - 24681.68 (11/2)	5583.619	15	17904.56	$6F_{9/2}^{\circ}$ - 25402.61 (11/2)*
6219.809	50	16073.22	$6F_{9/2}^{\circ}$ - 23571.27 (9/2)	5559.525	150 <i>L</i>	17982.15	$6F_{11/2}^{\circ}$ - 26591.40 (13/2)
6214.768	75	16086.25	$6F_{9/2}^{\circ}$ - 23584.31 (7/2)	5559.188	5	17983.24	$6F_{3/2}^{\circ}$ - 23443.79 (5/2)
6208.146	80	16103.41	$6F_{5/2}^{\circ}$ - 21976.26 (7/2)	5555.354	100	17995.66	$6F_{7/2}^{\circ}$ - 24558.56 (9/2)
6193.898	200	16140.46	$6F_{5/2}^{\circ}$ - 22013.40 (3/2)	5546.769	10 <i>U</i>	18023.51	$6F_{9/2}^{\circ}$ - 25521.55 (9/2)
6186.089	2	16160.83	$6F_{11/2}^{\circ}$ - 24770.04 (13/2)	5537.379	200 <i>C</i>	18054.07	$6F_{5/2}^{\circ}$ - 23926.91 (3/2)
6168.403	5	16207.17	$6F_{5/2}^{\circ}$ - 22080.08 (5/2)	5529.938	3	18078.36	$6F_{3/2}^{\circ}$ - 23538.86 (5/2)
6165.603	70	16214.53	$6F_{5/2}^{\circ}$ - 23712.56 (7/2)	5527.883	5	18085.09	$6F_{11/2}^{\circ}$ - 26694.38 (11/2)*
6153.685	80	16245.93	$6F_{9/2}^{\circ}$ - 23743.96 (7/2)	5524.945	150	18094.70	$6F_{7/2}^{\circ}$ - 26703.97 (13/2)
6142.459	100	16275.62	$6F_{11/2}^{\circ}$ - 24884.90 (11/2)	5487.381	20	18218.57	$6F_{5/2}^{\circ}$ - 24091.39 (3/2)
6132.934	40	16300.90	$6H_{7/2}^{\circ}$ - 17104.72 (7/2)	5484.872	10 <i>U</i>	18226.90	$6F_{7/2}^{\circ}$ - 24789.86 (9/2)*
6132.125	200	16303.05	$6F_{11/2}^{\circ}$ - 24912.34 (11/2)	5471.313	150 <i>C</i>	18272.07	$6F_{3/2}^{\circ}$ - 23732.57 (1/2)
6117.382	100	16342.34	$6F_{7/2}^{\circ}$ - 22905.24 (5/2)	5430.600	25	18409.06	$6H_{3/2}^{\circ}$ - 20157.85 (7/2)
6085.359	150 <i>P</i>	16428.34	$6F_{5/2}^{\circ}$ - 22301.24 (5/2)	5421.374	10	18440.38	$6H_{11/2}^{\circ}$ - 21237.49 (9/2)
6082.102	30 <i>S</i>	16437.13	$6F_{5/2}^{\circ}$ - 22309.94 (7/2)	5413.736	100 <i>C</i>	18466.40	$6F_{3/2}^{\circ}$ - 23926.91 (3/2)
6053.356	150	16515.19	$6F_{5/2}^{\circ}$ - 22388.06 (3/2)	5408.846	100 <i>C</i>	18483.09	$6F_{11/2}^{\circ}$ - 23732.57 (1/2)
6039.646	30	16552.68	$6F_{5/2}^{\circ}$ - 22425.58 (5/2)	5405.328	2	18495.12	$6H_{13/2}^{\circ}$ - 22414.17 (11/2)
6032.252	90	16572.97	$6F_{5/2}^{\circ}$ - 24071.03 (9/2)	5392.454	10	18539.28	$6H_{15/2}^{\circ}$ - 23629.06 (13/2)
6015.361	5	16619.50	$6F_{3/2}^{\circ}$ - 22080.08 (5/2)	5390.622	20	18545.58	$6F_{5/2}^{\circ}$ - 24418.44 (7/2)
6013.146	40	16625.63	$6F_{7/2}^{\circ}$ - 23188.54 (5/2)	5365.939	50	18630.89	$6F_{3/2}^{\circ}$ - 24091.39 (3/2)
5980.648	100 <i>D</i>	16715.97	$6F_{7/2}^{\circ}$ - 23278.90 (7/2)	5352.567	25 <i>C</i>	18677.43	$6F_{11/2}^{\circ}$ - 23926.91 (3/2)
5973.373	10	16736.32	$6F_{9/2}^{\circ}$ - 24234.42 (9/2)	5334.737	200	18739.85	$6F_{9/2}^{\circ}$ - 26237.84 (9/2)
5969.214	20	16747.98	$6F_{5/2}^{\circ}$ - 25357.24 (11/2)	5305.840	75 <i>S</i>	18841.91	$6F_{11/2}^{\circ}$ - 24091.39 (3/2)
5960.957	3	16771.18	$6F_{7/2}^{\circ}$ - 23334.10 (5/2)	5305.732	50 <i>B</i>	18842.30	$6H_{3/2}^{\circ}$ - 22761.33 (11/2)
5959.738	125	16774.61	$6F_{7/2}^{\circ}$ - 23337.53 (7/2)	5303.854	7	18848.97	$6H_{15/2}^{\circ}$ - 23938.76 (13/2)
5953.095	100	16793.33	$6F_{11/2}^{\circ}$ - 25402.61 (11/2)	5290.064	15	18898.10	$6H_{13/2}^{\circ}$ - 22817.13 (11/2)
5951.191	100	16798.71	$6F_{5/2}^{\circ}$ - 22259.21 (1/2)	5286.398	2	18911.21	$6H_{15/2}^{\circ}$ - 20660.00 (7/2)
5936.504	200	16840.26	$6F_{9/2}^{\circ}$ - 24338.33 (9/2)	5252.679	25	19032.61	$6H_{15/2}^{\circ}$ - 24122.41 (13/2)
5922.221	3	16880.88	$6F_{7/2}^{\circ}$ - 23443.79 (5/2)	5230.058	10	19114.93	$6H_{13/2}^{\circ}$ - 23033.95 (11/2)
5909.318	5	16917.74	$6F_{7/2}^{\circ}$ - 23480.63 (9/2)	5227.747	20	19123.38	$6H_{11/2}^{\circ}$ - 21920.49 (9/2)
5905.899	100	16927.53	$6F_{5/2}^{\circ}$ - 22388.06 (3/2)	5218.879	100	19155.87	$6H_{15/2}^{\circ}$ - 24245.66 (13/2)
5899.630	20	16945.52	$6F_{9/2}^{\circ}$ - 24443.57 (9/2)	5217.699	30 <i>P</i>	19160.20	$6H_{9/2}^{\circ}$ - 20909.00 (7/2)
5892.838	10 <i>P</i>	16965.05	$6F_{3/2}^{\circ}$ - 22425.58 (5/2)	5185.467	100	19279.30	$6H_{13/2}^{\circ}$ - 23198.33 (11/2)
5877.354	30 <i>C</i>	17009.74	$6F_{5/2}^{\circ}$ - 22259.21 (1/2)	5171.689	10	19330.66	$6F_{11/2}^{\circ}$ - 27939.87 (9/2)
5878.337	50	17021.38	$6F_{7/2}^{\circ}$ - 23584.31 (7/2)	5166.176	10 <i>P</i>	19351.29	$6H_{9/2}^{\circ}$ - 21100.10 (7/2)
5869.549	30	17032.36	$6F_{5/2}^{\circ}$ - 22905.24 (5/2)	5165.623	150	19353.36	$6H_{15/2}^{\circ}$ - 24443.15 (13/2)
5859.865	50	17060.51	$6F_{9/2}^{\circ}$ - 24558.56 (9/2)	5165.443	100	19354.03	$6H_{7/2}^{\circ}$ - 20157.85 (7/2)
5844.718	200 <i>L</i>	17104.72	$6F_{5/2}^{\circ}$ - 17104.72 (7/2)	5150.995	150	19408.32	$6H_{11/2}^{\circ}$ - 22205.44 (9/2)
5833.183	125	17138.55	$6F_{11/2}^{\circ}$ - 22388.06 (3/2)	5146.298	500 <i>R3</i>	19426.03	$6H_{13/2}^{\circ}$ - 23345.07 (11/2)
5829.407	50	17149.65	$6F_{7/2}^{\circ}$ - 23712.56 (7/2)	5145.126	400 <i>R4</i>	19430.46	$6H_{15/2}^{\circ}$ - 24520.23 (13/2)

TABLE 4. Classified lines of Pm I—Continued

$\lambda$ (Å)	$I$	$\sigma$ (cm <sup>-1</sup> )	Classification	$\lambda$ (Å)	$I$	$\sigma$ (cm <sup>-1</sup> )	Classification
5136.752	5	19462.13	${}^6\text{H}_{7/2}^{\circ}$ — 20265.98 (5/2)	4892.516	700R2	20433.68	${}^6\text{H}_{7/2}^{\circ}$ — 21237.49 (9/2)
5132.836	10C	19476.98	${}^6\text{F}_{9/2}^{\circ}$ — 28086.21 (11/2)	4890.549	10	20441.89	${}^6\text{F}_{9/2}^{\circ}$ — 27939.87 (9/2)*
5129.749	200	19488.70	${}^6\text{H}_{9/2}^{\circ}$ — 21237.49 (9/2)	4889.191	50C	20447.57	${}^6\text{H}_{15/2}^{\circ}$ — 25537.36 (13/2)
5127.342	400R3	19497.85	${}^6\text{H}_{11/2}^{\circ}$ — 22294.96 (9/2)	4887.018	500R3	20456.66	${}^6\text{H}_{15/2}^{\circ}$ — 22205.44 (9/2)
5111.431	40	19558.54	${}^6\text{H}_{11/2}^{\circ}$ — 22355.68 (9/2)	4885.070	25	20464.82	${}^6\text{F}_{11/2}^{\circ}$ — 29074.03 (9/2)
5105.170	50	19582.53	${}^6\text{H}_{3/2}^{\circ}$ — 23501.57 (11/2)	4882.946	1	20473.72	${}^6\text{F}_{7/2}^{\circ}$ — 27036.66 (5/2)
5100.766	400R3	19599.44	${}^6\text{H}_{9/2}^{\circ}$ — 21348.22 (7/2)	4881.674	2C	20479.06	${}^6\text{H}_{11/2}^{\circ}$ — 23276.10 (9/2)*
5096.601	50	19615.45	${}^6\text{H}_{5/2}^{\circ}$ — 24705.25 (15/2)	4872.416	700R2	20517.97	${}^6\text{H}_{5/2}^{\circ}$ — 20517.96 (5/2)
5096.181	200	19617.07	${}^6\text{H}_{11/2}^{\circ}$ — 22414.17 (11/2)	4871.845	30	20520.37	${}^6\text{F}_{11/2}^{\circ}$ — 29129.60 (9/2)
5094.831	400R3	19622.27	${}^6\text{H}_{9/2}^{\circ}$ — 21371.05 (7/2)	4870.957	50	20524.11	${}^6\text{H}_{13/2}^{\circ}$ — 24443.15 (13/2)
5092.418	200	19631.56	${}^6\text{H}_{3/2}^{\circ}$ — 23550.60 (11/2)	4869.801	400R4	20528.98	${}^6\text{H}_{5/2}^{\circ}$ — 25618.77 (15/2)
5087.872	10C	19649.10	${}^6\text{H}_{11/2}^{\circ}$ — 22446.20 (11/2)	4866.147	100	20544.40	${}^6\text{H}_{7/2}^{\circ}$ — 21348.22 (7/2)
5085.154	25	19659.61	${}^6\text{H}_{11/2}^{\circ}$ — 22456.72 (9/2)	4865.724	500R2	20546.19	${}^6\text{H}_{5/2}^{\circ}$ — 22294.96 (9/2)
5081.182	50	19674.98	${}^6\text{F}_{7/2}^{\circ}$ — 26237.84 (9/2)	4865.302	300R3	20547.97	${}^6\text{H}_{11/2}^{\circ}$ — 23345.07 (11/2)
5079.822	100	19680.24	${}^6\text{H}_{5/2}^{\circ}$ — 24770.04 (13/2)	4862.183	100	20561.15	${}^6\text{H}_{9/2}^{\circ}$ — 22309.94 (7/2)
5072.149	150	19710.02	${}^6\text{H}_{13/2}^{\circ}$ — 23629.06 (13/2)	4860.745	700R1	20567.23	${}^6\text{H}_{7/2}^{\circ}$ — 21371.05 (7/2)
5071.094	25	19714.12	${}^6\text{H}_{5/2}^{\circ}$ — 20517.96 (5/2)	4860.619	400R3	20567.77	${}^6\text{H}_{5/2}^{\circ}$ — 20567.76 (5/2)
5058.311	300R2	19763.93	${}^6\text{H}_{7/2}^{\circ}$ — 20567.76 (5/2)	4858.210	25	20577.97	${}^6\text{F}_{9/2}^{\circ}$ — 28075.94 (9/2)
5051.739	60	19789.65	${}^6\text{H}_{11/2}^{\circ}$ — 22586.77 (9/2)	4856.687	100	20584.42	${}^6\text{H}_{13/2}^{\circ}$ — 24503.45 (11/2)
5043.672	15	19821.30	${}^6\text{F}_{9/2}^{\circ}$ — 27319.28 (7/2)	4855.787	75	20588.23	${}^6\text{F}_{9/2}^{\circ}$ — 28086.21 (11/2)
5034.813	50L	19856.18	${}^6\text{H}_{7/2}^{\circ}$ — 20660.00 (7/2)	4852.727	350R3	20601.21	${}^6\text{H}_{13/2}^{\circ}$ — 24520.23 (13/2)
5030.805	200R3	19871.99	${}^6\text{H}_{7/2}^{\circ}$ — 20675.81 (5/2)	4851.808	50	20605.12	${}^6\text{F}_{9/2}^{\circ}$ — 26065.63 (1/2)
5029.624	50C	19876.66	${}^6\text{H}_{9/2}^{\circ}$ — 21625.45 (9/2)	4851.379	35	20606.94	${}^6\text{H}_{9/2}^{\circ}$ — 22355.68 (9/2)
5026.019	30	19890.91	${}^6\text{F}_{9/2}^{\circ}$ — 25351.46 (3/2)	4849.663	20	20614.23	${}^6\text{H}_{3/2}^{\circ}$ — 24533.27 (11/2)
5019.185	100	19918.00	${}^6\text{H}_{9/2}^{\circ}$ — 21666.80 (7/2)	4844.012	200	20638.28	${}^6\text{H}_{11/2}^{\circ}$ — 23435.40 (9/2)
5007.559	15	19964.24	${}^6\text{H}_{11/2}^{\circ}$ — 22761.33 (11/2)	4841.379	15	20649.50	${}^6\text{F}_{9/2}^{\circ}$ — 26522.35 (3/2)
5002.571	10	19984.15	${}^6\text{H}_{9/2}^{\circ}$ — 21732.93 (9/2)	4840.626	100	20652.71	${}^6\text{F}_{9/2}^{\circ}$ — 28150.73 (7/2)
4997.095	500R1	20006.04	${}^6\text{H}_{5/2}^{\circ}$ — 20006.04 (3/2)	4838.919	400R4	20660.00	${}^6\text{H}_{5/2}^{\circ}$ — 20660.00 (7/2)
4993.676	50	20019.74	${}^6\text{H}_{13/2}^{\circ}$ — 23938.76 (13/2)	4837.655	800R1	20665.40	${}^6\text{H}_{9/2}^{\circ}$ — 22414.17 (11/2)
4993.601	75	20020.04	${}^6\text{H}_{11/2}^{\circ}$ — 22817.13 (11/2)	4835.220	10S	20675.80	${}^6\text{H}_{5/2}^{\circ}$ — 20675.81 (5/2)
4981.727	5L	20067.76	${}^6\text{F}_{7/2}^{\circ}$ — 26630.56 (5/2)	4833.506	50	20683.14	${}^6\text{F}_{7/2}^{\circ}$ — 27245.99 (5/2)
4975.162	5	20094.24	${}^6\text{H}_{3/2}^{\circ}$ — 24013.29 (13/2)*	4833.417	75	20683.52	${}^6\text{H}_{11/2}^{\circ}$ — 23480.63 (9/2)
4973.250	25	20101.97	${}^6\text{F}_{7/2}^{\circ}$ — 25351.46 (3/2)	4832.297	50	20688.31	${}^6\text{F}_{9/2}^{\circ}$ — 28186.31 (7/2)
4972.448	150R5	20105.21	${}^6\text{H}_{7/2}^{\circ}$ — 20909.00 (7/2)	4830.170	125C	20697.42	${}^6\text{H}_{5/2}^{\circ}$ — 22446.20 (11/2)
4968.843	35S	20119.79	${}^6\text{H}_{13/2}^{\circ}$ — 24038.82 (11/2)	4828.524	20	20704.47	${}^6\text{H}_{11/2}^{\circ}$ — 23501.57 (11/2)
4965.600	5	20132.94	${}^6\text{F}_{7/2}^{\circ}$ — 26695.79 (5/2)	4827.716	400R3	20707.94	${}^6\text{H}_{9/2}^{\circ}$ — 22456.72 (9/2)
4959.461	700R1	20157.86	${}^6\text{H}_{5/2}^{\circ}$ — 20157.85 (7/2)	4821.054	40	20736.56	${}^6\text{F}_{5/2}^{\circ}$ — 26609.39 (5/2)
4956.053	100	20171.72	${}^6\text{H}_{9/2}^{\circ}$ — 21920.49 (9/2)	4817.116	400R4	20753.51	${}^6\text{H}_{11/2}^{\circ}$ — 23550.60 (11/2)
4948.286	125	20203.38	${}^6\text{H}_{13/2}^{\circ}$ — 24122.41 (13/2)	4816.438	75	20756.43	${}^6\text{F}_{7/2}^{\circ}$ — 27319.28 (7/2)
4946.851	100	20209.24	${}^6\text{H}_{11/2}^{\circ}$ — 23006.35 (11/2)	4816.131	7P	20757.75	${}^6\text{F}_{5/2}^{\circ}$ — 26630.56 (5/2)
4945.127	125	20216.28	${}^6\text{H}_{5/2}^{\circ}$ — 25306.07 (15/2)	4815.000	50C	20762.63	${}^6\text{H}_{13/2}^{\circ}$ — 24681.68 (11/2)
4942.390	40S	20227.48	${}^6\text{H}_{9/2}^{\circ}$ — 21976.26 (7/2)	4812.914	100	20771.63	${}^6\text{F}_{7/2}^{\circ}$ — 27334.48 (5/2)
4940.101	30	20236.85	${}^6\text{H}_{11/2}^{\circ}$ — 23033.95 (11/2)	4812.323	100	20774.18	${}^6\text{H}_{11/2}^{\circ}$ — 23571.27 (9/2)
4934.027	50	20261.77	${}^6\text{H}_{13/2}^{\circ}$ — 24180.80 (13/2)	4811.850	50P	20776.22	${}^6\text{F}_{9/2}^{\circ}$ — 28274.21 (9/2)
4932.994	600R1	20266.01	${}^6\text{H}_{5/2}^{\circ}$ — 20265.98 (5/2)	4809.536	700R2	20786.22	${}^6\text{H}_{13/2}^{\circ}$ — 24705.25 (15/2)
4929.960	3	20278.48	${}^6\text{F}_{7/2}^{\circ}$ — 26841.36 (5/2)	4808.994	30	20788.56	${}^6\text{F}_{7/2}^{\circ}$ — 27351.42 (7/2)
4925.637	10C	20296.28	${}^6\text{H}_{5/2}^{\circ}$ — 21100.10 (7/2)	4802.618	30C	20816.16	${}^6\text{F}_{1/2}^{\circ}$ — 26065.63 (1/2)
4918.283	400R4	20326.62	${}^6\text{H}_{13/2}^{\circ}$ — 24245.66 (13/2)	4801.356	900R1	20821.63	${}^6\text{H}_{11/2}^{\circ}$ — 21625.45 (9/2)
4917.008	2C	20331.90	${}^6\text{F}_{9/2}^{\circ}$ — 27829.89 (9/2)	4801.051	100	20822.95	${}^6\text{F}_{5/2}^{\circ}$ — 26695.79 (5/2)
4916.047	3C	20335.87	${}^6\text{H}_{9/2}^{\circ}$ — 22084.65 (7/2)	4800.083	3	20827.15	${}^6\text{F}_{9/2}^{\circ}$ — 28325.13 (7/2)
4915.234	125	20339.23	${}^6\text{H}_{7/2}^{\circ}$ — 21143.06 (7/2)	4799.491	150	20829.72	${}^6\text{H}_{15/2}^{\circ}$ — 25919.50 (13/2)
4904.278	300R4	20384.67	${}^6\text{H}_{15/2}^{\circ}$ — 25474.46 (15/2)*	4798.977	700R1	20831.95	${}^6\text{H}_{11/2}^{\circ}$ — 23629.06 (13/2)
4904.035	15S	20385.68	${}^6\text{F}_{11/2}^{\circ}$ — 28994.90 (11/2)	4797.586	100	20837.99	${}^6\text{H}_{11/2}^{\circ}$ — 22586.77 (9/2)
4900.296	400R3	20401.23	${}^6\text{H}_{11/2}^{\circ}$ — 23198.33 (11/2)	4797.171	15	20839.79	${}^6\text{F}_{3/2}^{\circ}$ — 26300.30 (1/2)
4897.663	3P	20412.20	${}^6\text{F}_{5/2}^{\circ}$ — 26285.02 (7/2)*	4794.588	250	20851.02	${}^6\text{H}_{13/2}^{\circ}$ — 24770.04 (13/2)

TABLE 4. Classified lines of Pm I—Continued

$\lambda$ (Å)	<i>I</i>	$\sigma$ (cm <sup>-1</sup> )	Classification	$\lambda$ (Å)	<i>I</i>	$\sigma$ (cm <sup>-1</sup> )	Classification
4794.207	50	20852.68	<sup>6</sup> F <sub>5/2</sub> <sup>o</sup> —26725.52 (3/2)	4678.918	60	21366.48	<sup>6</sup> H <sub>15/2</sub> <sup>o</sup> —26456.26 (13/2)
4791.840	35C	20862.98	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> —21666.80 (7/2)	4678.093	400R5	21370.25	<sup>6</sup> F <sub>3/2</sub> <sup>o</sup> —26830.74 (3/2)
4782.081	150R4	20905.55	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —22654.34 (7/2)	4677.916	500R2	21371.06	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —21371.05 (7/2)
4781.292	900R1	20909.00	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —20909.00 (7/2)	4677.456	200	21373.16	<sup>6</sup> F <sub>5/2</sub> <sup>o</sup> —27245.99 (5/2)
4780.285	150	20913.41	<sup>6</sup> F <sub>5/2</sub> <sup>o</sup> —27476.28 (7/2)	4675.764	5C	21380.89	<sup>6</sup> F <sub>3/2</sub> <sup>o</sup> —26841.36 (5/2)
4776.699	100R5	20929.11	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> —21732.93 (9/2)	4675.148	35	21383.71	<sup>6</sup> H <sub>11/2</sub> <sup>o</sup> —24180.80 (13/2)
4773.458	700R1	20943.32	<sup>6</sup> H <sub>9/2</sub> <sup>o</sup> —23740.42 (13/2)*	4674.420	500R4	21387.04	<sup>6</sup> H <sub>13/2</sub> <sup>o</sup> —25306.07 (15/2)
4771.916	150	20950.09	<sup>6</sup> F <sub>7/2</sub> <sup>o</sup> —27512.95 (5/2)	4671.759	400	21399.22	<sup>6</sup> F <sub>11/2</sub> <sup>o</sup> —30008.40 (11/2)
4770.142	150	20957.88	<sup>6</sup> F <sub>5/2</sub> <sup>o</sup> —26830.74 (3/2)	4671.234	500R2	21401.63	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> —22205.44 (9/2)
4768.866	25	20963.48	<sup>6</sup> H <sub>11/2</sub> <sup>o</sup> —23760.57 (11/2)	4665.188	400R2	21429.36	<sup>6</sup> H <sub>9/2</sub> <sup>o</sup> —23178.13 (7/2)
4767.719	60	20968.53	<sup>6</sup> F <sub>5/2</sub> <sup>o</sup> —26841.36 (5/2)	4663.455	600R2	21437.33	<sup>6</sup> H <sub>9/2</sub> <sup>o</sup> —24234.42 (9/2)
4766.024	5C	20975.99	<sup>6</sup> F <sub>11/2</sub> <sup>o</sup> —29585.21 (11/2)	4663.261	300R4	21438.22	<sup>6</sup> H <sub>3/2</sub> <sup>o</sup> —25357.24 (11/2)
4763.670	2	20986.35	<sup>6</sup> F <sub>11/2</sub> <sup>o</sup> —29595.58 (9/2)	4661.729	25U	21445.26	<sup>6</sup> F <sub>7/2</sub> <sup>o</sup> —28008.09 (7/2)*
4762.569	700R1	20991.20	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —26080.99 (17/2)	4661.476	40C	21446.43	<sup>6</sup> F <sub>5/2</sub> <sup>o</sup> —27319.28 (7/2)
4762.309	5C	20992.35	<sup>6</sup> F <sub>9/2</sub> <sup>o</sup> —28490.35 (9/2)	4661.012	60	21448.56	<sup>6</sup> H <sub>11/2</sub> <sup>o</sup> —24245.66 (13/2)
4762.095	75U	20993.29	<sup>6</sup> H <sub>3/2</sub> <sup>o</sup> —24912.34 (11/2)	4660.794	500R3	21449.56	<sup>6</sup> H <sub>9/2</sub> <sup>o</sup> —23198.33 (11/2)
4758.996	800R1	21006.96	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —26096.75 (17/2)	4659.745	2	21454.40	<sup>6</sup> F <sub>11/2</sub> <sup>o</sup> —30063.62 (9/2)
4758.694	150	21008.29	<sup>6</sup> F <sub>3/2</sub> <sup>o</sup> —26468.80 (1/2)	4659.383	300	21456.06	<sup>6</sup> H <sub>13/2</sub> <sup>o</sup> —26545.85 (17/2)
4757.968	5	21011.50	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —26101.28 (13/2)	4658.169	75	21461.65	<sup>6</sup> F <sub>5/2</sub> <sup>o</sup> —27334.48 (5/2)
4757.732	500R2	21012.54	<sup>6</sup> H <sub>11/2</sub> <sup>o</sup> —22761.33 (9/2)	4657.301	50	21465.66	<sup>6</sup> H <sub>15/2</sub> <sup>o</sup> —26555.44 (13/2)
4756.247	90D	21019.10	<sup>6</sup> F <sub>3/2</sub> <sup>o</sup> —26479.61 (3/2)	4655.046	500R5	21476.05	<sup>6</sup> F <sub>1/2</sub> <sup>o</sup> —26725.52 (3/2)
4751.701	20	21039.21	<sup>6</sup> F <sub>11/2</sub> <sup>o</sup> —29648.42 (9/2)*	4654.496	400	21478.59	<sup>6</sup> F <sub>3/2</sub> <sup>o</sup> —27351.42 (7/2)
4749.076	70D	21050.84	<sup>6</sup> F <sub>1/2</sub> <sup>o</sup> —26300.30 (1/2)	4653.413	400R3	21483.59	<sup>6</sup> H <sub>13/2</sub> <sup>o</sup> —25402.61 (11/2)
4747.268	3C	21058.86	<sup>6</sup> F <sub>7/2</sub> <sup>o</sup> —27621.74 (7/2)	4652.834	5	21486.26	<sup>6</sup> H <sub>13/2</sub> <sup>o</sup> —25405.29 (13/2)
4746.595	30	21061.84	<sup>6</sup> F <sub>3/2</sub> <sup>o</sup> —26522.35 (3/2)	4650.525	500	21496.93	<sup>6</sup> F <sub>9/2</sub> <sup>o</sup> —28994.90 (11/2)
4745.282	10P	21067.67	<sup>6</sup> F <sub>9/2</sub> <sup>o</sup> —28565.66 (9/2)	4650.421	600R5	21497.41	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> —22301.24 (5/2)
4745.128	350R2	21068.35	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —22817.13 (11/2)	4649.508	600	21501.63	<sup>6</sup> H <sub>15/2</sub> <sup>o</sup> —26591.40 (13/2)
4739.776	200C	21092.15	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —26181.98 (13/2)	4648.537	50	21506.12	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> —22309.94 (7/2)
4738.779	2	21096.58	<sup>6</sup> F <sub>11/2</sub> <sup>o</sup> —29705.77 (11/2)	4647.028	600R3	21513.11	<sup>6</sup> F <sub>7/2</sub> <sup>o</sup> —28075.94 (9/2)
4737.987	500R2	21100.11	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —21100.10 (7/2)	4645.234	40	21521.41	<sup>6</sup> F <sub>7/2</sub> <sup>o</sup> —28084.28 (5/2)
4735.915	75C	21109.34	<sup>6</sup> F <sub>9/2</sub> <sup>o</sup> —28607.33 (9/2)	4643.959	200	21527.32	<sup>6</sup> H <sub>9/2</sub> <sup>o</sup> —23276.10 (9/2)
4734.274	800R1	21116.66	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> —21920.49 (9/2)	4643.355	700R2	21530.13	<sup>6</sup> H <sub>9/2</sub> <sup>o</sup> —23278.90 (7/2)
4728.678	400R3	21141.65	<sup>6</sup> H <sub>11/2</sub> <sup>o</sup> —23938.76 (13/2)	4640.961	400R3	21541.23	<sup>6</sup> H <sub>11/2</sub> <sup>o</sup> —24338.33 (9/2)
4728.361	700R1	21143.06	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —21143.06 (7/2)	4638.672	100	21551.86	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> —22355.68 (9/2)
4727.144	10	21148.50	<sup>6</sup> F <sub>11/2</sub> <sup>o</sup> —29757.69 (11/2)	4633.473	300B	21576.04	<sup>6</sup> F <sub>9/2</sub> <sup>o</sup> —29074.03 (9/2)
4727.062	300	21148.87	<sup>6</sup> F <sub>3/2</sub> <sup>o</sup> —26609.39 (5/2)	4633.452	600R3	21576.14	<sup>6</sup> F <sub>3/2</sub> <sup>o</sup> —27036.66 (5/2)
4723.722	125	21163.83	<sup>6</sup> F <sub>3/2</sub> <sup>o</sup> —27036.66 (5/2)	4632.353	10P	21581.26	<sup>6</sup> F <sub>1/2</sub> <sup>o</sup> —26830.74 (3/2)
4722.332	40D	21170.06	<sup>6</sup> F <sub>3/2</sub> <sup>o</sup> —26630.56 (5/2)	4630.930	200	21587.89	<sup>6</sup> F <sub>7/2</sub> <sup>o</sup> —28150.73 (7/2)
4721.752	30B	21172.66	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> —21976.26 (7/2)	4630.742	15	21588.77	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —23337.53 (7/2)
4719.607	2	21182.28	<sup>6</sup> F <sub>9/2</sub> <sup>o</sup> —28680.26 (7/2)	4630.349	125	21590.60	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —21590.60 (3/2)
4718.800	100	21185.90	<sup>6</sup> H <sub>9/2</sub> <sup>o</sup> —22934.70 (7/2)	4629.127	150	21596.30	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —23345.07 (11/2)
4717.351	75	21192.41	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —26282.20 (13/2)	4627.595	400R3	21603.45	<sup>6</sup> F <sub>5/2</sub> <sup>o</sup> —27476.28 (7/2)
4712.058	50	21216.21	<sup>6</sup> H <sub>11/2</sub> <sup>o</sup> —24013.29 (13/2)*	4625.289	500R2	21614.22	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —26703.97 (13/2)
4711.368	125C	21219.32	<sup>6</sup> F <sub>1/2</sub> <sup>o</sup> —26468.80 (1/2)	4624.409	900W	21618.33	<sup>6</sup> H <sub>3/2</sub> <sup>o</sup> —25537.36 (13/2)
4708.968	150	21230.14	<sup>6</sup> F <sub>1/2</sub> <sup>o</sup> —26479.61 (3/2)	4623.675	700R1	21621.76	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> —22425.58 (5/2)
4706.401	100	21241.72	<sup>6</sup> H <sub>11/2</sub> <sup>o</sup> —24038.82 (11/2)	4623.310	500	21623.47	<sup>6</sup> F <sub>7/2</sub> <sup>o</sup> —28186.31 (7/2)
4705.114	100	21247.52	<sup>6</sup> F <sub>11/2</sub> <sup>o</sup> —29856.72 (11/2)	4621.571	500	21631.61	<sup>6</sup> F <sub>9/2</sub> <sup>o</sup> —29129.60 (9/2)
4702.886	10W	21257.59	<sup>6</sup> H <sub>9/2</sub> <sup>o</sup> —23006.35 (11/2)	4619.750	500R4	21640.13	<sup>6</sup> F <sub>5/2</sub> <sup>o</sup> —27512.95 (5/2)
4701.242	1	21265.03	<sup>6</sup> F <sub>3/2</sub> <sup>o</sup> —26725.52 (3/2)	4618.487	400R4	21646.05	<sup>6</sup> H <sub>11/2</sub> <sup>o</sup> —24443.15 (13/2)
4700.803	100C	21267.01	<sup>6</sup> F <sub>7/2</sub> <sup>o</sup> —27829.89 (9/2)	4618.398	200	21646.47	<sup>6</sup> H <sub>11/2</sub> <sup>o</sup> —24443.57 (9/2)
4699.508	200	21272.87	<sup>6</sup> F <sub>1/2</sub> <sup>o</sup> —26522.35 (3/2)	4617.023	600R1	21652.91	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> —22456.72 (9/2)
4699.276	150	21273.92	<sup>6</sup> H <sub>9/2</sub> <sup>o</sup> —24071.03 (9/2)	4615.961	50SP	21657.89	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —21657.89 (5/2)
4698.761	150	21276.25	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> —22080.08 (5/2)	4614.670	2	21663.96	<sup>6</sup> F <sub>9/2</sub> <sup>o</sup> —29161.96 (9/2)
4697.749	30C	21280.83	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> —22084.65 (7/2)	4614.059	150	21666.82	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —21666.80 (7/2)
4696.796	500R2	21285.16	<sup>6</sup> H <sub>9/2</sub> <sup>o</sup> —23033.95 (11/2)	4612.787	10	21672.80	<sup>6</sup> F <sub>11/2</sub> <sup>o</sup> —30281.98 (13/2)
4693.589	40L	21299.70	<sup>6</sup> F <sub>11/2</sub> <sup>o</sup> —29908.90 (9/2)	4609.846	500R3	21686.62	<sup>6</sup> H <sub>9/2</sub> <sup>o</sup> —23435.40 (9/2)
4687.952	20	21325.31	<sup>6</sup> H <sub>9/2</sub> <sup>o</sup> —24122.41 (13/2)	4607.062	150	21699.73	<sup>6</sup> H <sub>3/2</sub> <sup>o</sup> —25618.77 (15/2)
4682.920	700R1	21348.22	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> —21348.22 (7/2)	4605.657	600R2	21706.35	<sup>6</sup> H <sub>11/2</sub> <sup>o</sup> —24503.45 (11/2)
4682.266	150	21351.20	<sup>6</sup> F <sub>11/2</sub> <sup>o</sup> —29960.42 (11/2)*	4604.739	400C	21710.67	<sup>6</sup> F <sub>7/2</sub> <sup>o</sup> —28273.52 (7/2)
4680.223	2	21360.52	<sup>6</sup> F <sub>7/2</sub> <sup>o</sup> —27923.37 (5/2)	4604.593	400	21711.36	<sup>6</sup> F <sub>7/2</sub> <sup>o</sup> —28274.21 (9/2)









TABLE 4. Classified lines of Pm I—Continued

$\lambda$ (Å)	$I$	$\sigma$ (cm <sup>-1</sup> )	Classification	$\lambda$ (Å)	$I$	$\sigma$ (cm <sup>-1</sup> )	Classification
3442.984	50	29036.25	<sup>6</sup> H <sub>9/2</sub> <sup>o</sup> — 30785.03 ( 9/2)	3345.327	20	29883.85	<sup>6</sup> H <sub>5/2</sub> <sup>o</sup> — 29883.87 ( 7/2)
3437.803	15	29080.01	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> — 29883.87 ( 7/2)	3341.011	50	29922.46	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> — 30726.26 ( 9/2)
3420.721	75	29225.22	<sup>6</sup> H <sub>11/2</sub> <sup>o</sup> — 32022.32 ( 9/2)	3334.464	75	29981.21	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> — 30785.03 ( 9/2)
3416.680	40	29259.79	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> — 30063.62 ( 9/2)				
3416.484	50	29261.47	<sup>6</sup> H <sub>13/2</sub> <sup>o</sup> — 33180.50 (15/2)	3319.548	25	30115.92	<sup>6</sup> H <sub>9/2</sub> <sup>o</sup> — 31846.70 (11/2)
				3302.264	25	30273.53	<sup>6</sup> H <sub>9/2</sub> <sup>o</sup> — 32022.32 ( 9/2)
3405.661	3	29354.45	<sup>6</sup> H <sub>9/2</sub> <sup>o</sup> — 31103.24 (11/2)	3283.174	25	30449.56	<sup>6</sup> H <sub>11/2</sub> <sup>o</sup> — 33246.65 ( 9/2)
3373.083	75	29637.95	<sup>6</sup> H <sub>11/2</sub> <sup>o</sup> — 32435.06 (13/2)	3202.305	25	31218.48	<sup>6</sup> H <sub>7/2</sub> <sup>o</sup> — 32022.32 ( 9/2)

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