URSI National Committee Report, XIV General Assembly, Tokyo, September 1963: Commission 7. Radio Electronics

Review of developments occurring within the United States in the field of Radio Electronics during the triennium 1960 through 1962.

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1. Progress in Microwave Power Tubes

I. Feinstein*

The microwave tube field reached maturity during the past three years. This is reflected in the nature of the advances made which consist chiefly of the attainment of higher power and greater gain and bandwidth in both linear beam and crossed-field types.

While the exploration of new principles led to deeper insight into electron-wave interaction, the performance competition from well-established tube types has thus far prevented the emergence of practical new types.

1. Klystrons

The attainment of 100 kw of CW power at 8 Gc/s [McCune et al., 1961] is probably the most spectacular current feat in the tube field. This performance is a result of excellent beam control, good thermal engineering, and new output window fabrication techniques. Similarly, impressive amounts of power have been obtained at other frequency bands in single tubes employing solid beams [Holaday and Petersen, 1961]. Stagger tuning and coupled cavity output circuits [Biss, 1962] have led to electronic bandwidths of 5 percent for 1 Mw peak power klystrons. The principle of distributed klystron interaction employing slow wave structures in lieu of gaps has been verified experimentally [Chodorow and Wesselberg, 1961]. Large signal calculations which take into account the effect of

space charge and of velocity spread have been made [Weber, 1960] and form the basis for optimizing parameters such as drift tube length.

As the limit of power attainable in a single tube is approached, methods of paralleling tubes are receiving increased attention. The multiple beam klystron [Boyd et al., 1962] represents one such method in which corresponding cavities of adjacent klystrons are coupled to form a traveling wave structure.

2. Traveling Wave Tubes

A measure of the engineering progress which has been made in this area during the review period may be gleaned from the fact that $\frac{1}{2}$ w CW at 5 to 6 mm [McDowell et al., 1960] represented the frontier of the millimeter wave amplifier field at the start of the period while 150 w CW in the same band [Nevins, 1962] is the current frontier. The first tube employed a helix while the second operates on the first forward space harmonic of a (inductively) coupled cavity structure.

For satellite repeater applications 10 w amplifier tubes of about 30 percent efficiency and weighing under 1 lb have been developed [Bodmer, 1961; Highstretc et al., 1962].

Considerable effort has been devoted to multiple depressed collectors [Dunn et al., 1960] but only a single stage of depression has appeared in practical tubes.

A backward wave oscillator employing a coupled cavity circuit has generated 2 w CW over the frequency range 50 to 75 Gc/s [Schumacher, 1962].

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3. Crossed-Field Tubes

Magnetron oscillators based on the circular electric mode coupling principle have attained peak power levels of 1.5 Mw at X-band [Cook et al., 1961] and 300 kw at 8.6 mm [La Plante, 1963]. The latter tube is of inverted type leading to a large increase in cathode surface area and extracting the power in a circular electric mode output waveguide. A K_u band carcinatron has been developed with an output power of 100 w [Moats et al., 1962]. A depressed collector has been employed to improve the efficiency of this tube type [Oscpchuk, 1962]. Megawatt peak power, re-entrant stream traveling wave amplifiers have attained saturated gains of 20 to 30 db over 10 percent bandwidths [McDowell et al., 1962]. A drift space is employed in these tubes to debunch the re-entering stream, while a control electrode located in this space functions as a modulating anode.

A 25 w CW 1-lb, S-band amplitron has been designed for space application. This tube has 20 db gain and 50 percent efficiency [Teich et al., 1962]. Development of forward-wave, injected beam, CW amplifiers has reached the 1 kw power level. The large signal theory of this type of interaction has been verified experimentally [Anderson, 1961].

4. Fast Wave Interactions

Theoretical work on d-c pumped cyclotron wave tubes employing for the most part the coupled mode formulation and the kinetic power viewpoint was carried on extensively during this period [Siegmann, 1960; Gordon, 1960]. Experimentally, the fast wave $(\beta=0)$ couplers introduced by Adler in his RF-pumped, low-noise tubes were combined with periodic structures producing d-c electric or magnetic fields to provide gain [Gordon, 1960; Wesselberg et al., 1962]. By the end of this period most of this work had foundered on the problem of beam control in the presence of a gain mechanism which amplified dc perturbations on the beam.

A cyclotron wave backward wave oscillator which has the novel feature of requiring a first-order RF magnetic field force for its operation was demonstrated [Chow et al., 1962].

A fast wave transverse interaction utilizing periodic magnetic deflection was demonstrated. By operating in a circular electric mode waveguide more than 1 Mw of peak power was obtained above X-band [Phillips, 1960].

5. Electron Guns

Supporting work in this category formed the basis for many of the advances recorded above. For solid beams the emphasis was on high perveance (up to five) and high area convergence (up to [Frost et al., 1962]. For more advanced 1000:1)tubes hollow beams were under intensive investigation. These were mainly of the magnetron injection type, and utilized space charge solutions in crossed fields to determine electrode shapes [Kino et al., 1962]. A mathematical approach [Harker, 1960] which made possible convergent computer solutions proved valuable in this area. Methods of focusing hollow beams through long circuits were also explored [Cook, 1961].

Crossed-field gun design methods were also evolved, but RF instabilities in the resultant beam set limits to the range of parameters which could be successfully employed [Midford et al., 1962].

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2. 1960–1963 Advances in the State-of-the-Art of Low-Noise Beam-Type Microwave Tubes

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1.Noise Reduction in Beam-Type Amplifiers— General Aspects

The traveling-wave tube and the electron-beam parametric amplifier have emerged as the important low-noise beam-type amplifiers. During the last three years no work has been reported on low-noise klystrons and triodes and little work is being done on M-type low-noise amplifiers [Anderson, May 1960]. The best noise factors achieved with these latter tube types are higher by a wide margin than those of the traveling-wave tube and the beam-type parametric amplifier, both of which have reached best noise factors of around 1.5 db in the lower microwave frequency ranges. The basic mechanisms of all beam-type tubes, although different in detail, have certain aspects in common. As coupled wave phenomena, they can be treated in a generalized form [Pease, June 1961]. Beam refrigeration using large magnetic fields, experimentally demonstrated in the beam-type parametric amplifier, may well be applicable to the traveling-wave tube also [Adler and Wade, July 1960]. Cooling of the slow spacecharge wave in the traveling-wave tube using wellproven methods of parametric amplifiers has been proposed but not yet experimentally verified [Forster, November 1962].

2. Low-Noise Traveling-Wave Tube Amplifiers

The low-noise traveling-wave tube may be said to have reached the state of a mature product during the past three years. Its noise factor through frequencies up to K-band has been gradually decreasing. This has been achieved by further refining the hollowbeam low-noise gun described earlier [Currie and Forster, April 1958], and by further refining [Bosch and Niclas, March 1962] the solid-beam low-noise gun previously reported [Peter, December 1958]. Both types have been proven to be basically identical, leading to hollow-beam emission under the favorable condition of a strongly divergent electric field at the cathode. The lowest noise factor recently measured with a traveling-wave tube is 1.5 db measured in UHF-band [Vehn and Peter, 1963].

The effect of distributed loss of the helix on the noise factor in traveling-wave tubes [Yariv and Kompfner, May 1961] has been further investigated, [Bloom, June 1961; Israelsen, March 1962]. A noise factor reduction of about 1 db has been dem-

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onstrated in L-band by cooling the helix to liquid nitrogen temperature [Israelsen and Peter, July 1962].

Production refinement of low-noise traveling-wave tubes for specific applications has led to significant advances which include extreme phase and gain linearity of the amplifier, operational life of more than 10,000 hr, permanent-magnet focusing in magnetically shielded packages without loss of noise factor, extremely high tube-to-tube reproducibility, and high limiting capability for large input signals [Mita, Takeya, Hayasi, Kakizaki, Yoshida, October 1962].

3. Noise in Multivelocity Electron Beams

This is the area where the most radical noise factor improvements are still possible. Extensive study of noise-wave propagation through the potential minimum of a space-charge-limited diode or virtual cathode has been carried out both theoretically and experimentally [Whinnery, October 1960; Vivian, June 1960; Berghammer, September 1960; Mueller, March 1961; Yadavalli, June 1961; Pollack, May 1962; Eichenbaum and Hammer, September 1962]. Experimental results indicate that whenever a potential minimum exists, noise reduction takes place. Theory indicates that the electron-plasma frequency in the potential minimum should be in the order of the signal frequency or higher for effective noise reduction.

Several more or less idealized models of a multivelocity electron beam drifting at low voltage have been investigated. All theories indicate that the two uncorrelated noise waves become partially correlated in a low-velocity drift space [Morrison, November 1960; Gray, February 1960; Berghammer and Bloom, March 1960; Mihran, April 1962; Bobroff, June 1962].

Measurements of the basic noise quantities S and Π were made to verify these theories [Zacharias and Smullin, July 1960]. Measurements on a hollow-beam low-noise gun indicated S to be reduced while P/S remained constant.

Partition noise in electron beams resulting from partial current interception has been investigated [Ashkin and White, August 1960; Hart, August 1962] for axial and transverse beam waves.

In actual tubes partition noise can be essentially eliminated. Further, transverse-to-axial noise wave transformation on a beam passing through electric or magnetic lenses can be minimized by suitable design of the gun and by avoiding periodic magnetic beam focusing.

4. Emitters for Low-Noise Electron Guns

All other factors being the same, a cathode emitting lower-temperature electrons will produce a lower noise factor in the beam-type amplifier. For this reason many studies have been carried out on standard as well as novel cathodes for low-noise electron guns.

The fluctuations in the emission of tungsten-base impregnated barium cathodes were investigated [Brodie, October 1961; Hill and Rouze, April 1962]. Considerable progress and success has been achieved in electron cooling by interaction with a cold plasma. β -Eucryptite is a preferred source for the cool plasma in front of the cathode [Johnson, September 1962; Eichenbaum, June 1962; Anderson and Harris, 1960]. Electron temperature reduction from 1100 to 300 °C has been observed.

A new type of cold cathode based upon the tunnel mechanism is being studied. Electrons can tunnel from a core metal into vacuum if a suitable potential configuration exists near the surface inside the emitter. This is achieved by applying a small voltage between surface films [Geppert and Barnes, November 1961; Wade, Briggs, and Lesensky, March 1962].

Currents of several hundred microamps have been drawn from tunnel cathodes. The ratio between required cathode power and emitted current is, however, still much higher than that in a thermionic cathode.

5. Electron-Beam Parametric Amplifier

Most of the work on the electron beam parametric amplifier was concentrated on the degenerate type where the pump frequency is near twice the signal frequency. In this case, the idler frequency is very near the signal frequency and the noise factor measured with a broadband noise source is the so-called double channel noise figure. Best success has been achieved in the UHF, P- and L-frequency bands at an instantaneous bandwith of up to 100 Mc/s; broadband noise figures of about 1 db and gain from 20 to 45 db have been measured. Degenerate amplifiers have also been made for S- and C-band. X-band appears to be a practical upper limit due to the linear increase of magnetic field with frequency, 3500 gauss at X-band.

The presence of the idler near the signal frequency is a disadvantage in the/degenerate amplifier. Recent work has resulted in the successful operation of nondegenerate amplifiers in the region of 400 Mc/s using a 2100 Mc/s pump. A single-channel noise factor of less than 2 db was obtained. Development of similar nondegenerate tubes for L-band, pumped at C band, is now in process. For frequencies higher than that it is not yet known how the degenerate amplifier will perform.

Power output, and therefore dynamic range, limitations may be imposed on the electron-beam parametric amplifier by the fact that this tube operates best if the beam potential is below the ionization potential of the residual gas molecules left in the tube.

Considerable progress has been made in advancing

the theory of the electron-beam parametric or Adlertype amplifier. [Cook, Louisdell and Quate, January 1960; Lea-Wilson, Adler, Hrbek, and Wade, February 1960; Johnson, February 1960; Heffner and Wade, December 1960; Bridges and Ashkin, March 1960; Gordon, November 1960; Gould and Johnson, February 1961; Adler and Wade, April 1961; Fredericks, May 1961; Gordon, July 1961; Ashkin, June 1961; Chang, March 1961; Bløtekjaer, August 1962; Wessel-Berg and Bløtekjaer, September 1962; Adler and Hrbek, January, 1963.]

6. M-Type Low-Noise Amplifiers

M-type amplifiers have long been known as noisy amplifiers. Various instabilities of the electron plasma between cathode or sole and anode lead to considerable amounts of noise. These noise producing mechanisms, however, do not necessarily have to be present in the M-type. Considerable improvements in this type of amplifier can, therefore, be expected since present noise performance leaves much to be desired. Results achieved by Kluver on a Kompfnerdip coupler for M-type tubes look particularly promising [Anderson, May 1960; Van Duzer, Whinnery, June 1960; Kluver, June 1960].

7. Standards and State-of-the-Art Summaries

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3. Developments in Semiconductor Devices in the United States for the Period 1960–1963

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During the period 1960–1963 many existing devices were perfected. Early devices such as fieldeffect transistors became attractive for the first time. Devices such as tunnel diodes received less attention than for the previous 3-year period, whereas varactor diodes gained prominence as reliable frequency multipliers. Considerable progress was made in the field of integrated circuitry.

New devices such as the thin film transistors and incoherent and coherent infrared radiators were developed. It is believed that the many possible forms of laser diodes will have a considerable future.

The list of references contained in this review is only representative for the topics under discussion and is by no means complete.

1. Transistors

Transistors with cutoff frequencies of 1 Gc/s are presently commercially available. Transistors with cutoff frequencies of 2 to 3 Gc/s are in the development stage.

Metal-base transistors, based on the relatively long mean free path for hot electrons in metals have been proposed. In these devices the semiconductor base is replaced by a thin metal layer [Geppert, 1962].

2. Heterojunctions

Abrupt monocrystalline junctions between two different semiconductor materials of nearly equal lattice spacing can be made by epitaxial deposition of the one material on the other substrate [Anderson, 1960]. Ge-GaAs heterojunctions are an example. Such junctions open up possibilities for new types of devices.

3. Solar Cells

It has been found that silicon n on p base cells can tolerate greater integrated radiation fluxes than p on n base cells. GaAs cells appear to be even more resistant to the effects of proton and electron irradiation [Loferski, 1963].

4. Tunnel Diodes

There is a certain amount of interest in tunnel diode amplifiers for UHF and microwave frequencies [Yariv, 1961] and in tunnel diode oscillators [Sterzer, 1961]. The use of tunnel diodes in ultra-high-speed switching circuits is handicapped by the bilateral character of the device but some progress is being made toward their utilization in computer circuitry [Sims, 1961].

The tunnel diode shows considerable promise as a low-noise mixer. The negative resistance region is not beneficial for low-noise operation; the lowest noise figure is obtained when this region is missing [Sterzer, 1962] (the so-called "back diodes").

5. Varactor Diodes

Varactor diodes find at present limited use in low-noise up-converters and in low-noise parametric amplifiers. In the latter case low amplifier noise temperatures are obtained by cooling the idler circuit [Kurokawa, 1961].

Varactor diodes are becoming more and more important in frequency multiplier circuits. They provide crystal-controlled oscillator power in the 3 Gc/s and 9 Gc/s microwave bands comparable with the power available from low-power klystrons. Reasonable efficiencies can be obtained [Johnson, 9960].

A low-noise low-frequency varactor-diode amplifier was developed [Biard, 1963].

6. Field-Effect Transitors

Stable low-noise field-effect transistors have been developed during the last few years and these devices have attracted considerable interest. The noise behavior of these low-noise units is limited by the thermal noise generated in the conducting channel [van der Ziel, 1962; and van der Ziel, 1963].

A CdS polycrystalline field-effect transistor with insulated gate has been developed [Weimer, 1962]. If the stability problem can be overcome, this would be the first semiconductor amplifying device that does not need single crystal material for its fabrication.

7. Integrated Circuits

A considerable development is presently under way in integrated electronics [Lesk, 1960; Swan, 1961; and Wallmark, 1962]. More and more integrated circuits are becoming commercially available. The practical objectives aimed at are the following:

1. Increased capabilities per unit weight and per unit volume.

- 2. Size and weight reduction.
- 3. Increased reliability and reduced maintenance.
- 4. Improved power utilization.
- 5. Reduced cost.

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8. Tunnel Emission Devices

Metal-metal-oxide-metal-semiconductor transistors, have been studied [Spratt, 1961]. It was at first believed that the emission through the oxide layer was of the tunnel emission type and that the hot electrons passed through the metal base layer to be collected by the metal-semiconductor junction. Presently it is held that the emission, at least at room temperature is of the Schottky thermionic type involving thermionic emission from traps [Mead, 1962]. It is also held that in many units the current flow is through pinholes and that the transistor-like device operates more or less as a field-effect transistor [Hall, 1961–1962]. This does not exclude the possibility of constructing a true hot electron transistor, of course.

9. Radiating Diodes, Lasing Diodes

Many junctions, when biased in the forward direction, show (infrared) light emission due to hole-electron recombination radiation. The wavelength of the emitted light is generally somewhat smaller than the energy gap and depends somewhat on the operating temperature and the operating current. GaAs diodes are a good example of such devices [Pankove, 1962a]. Because of the short lifetime of the minority carriers ($\simeq 10^{-10}$ sec) microwave modulation is possible [Pankove, 1962b].

At sufficiently large forward current densities the junctions exhibit laser action, resulting in a coherent infrared radiation of much narrower spectral width than the incoherent radiation. The effect has been observed in GaAs [Hall, 1962] at 8400 Å and in other crystals such as InAs at 31,000 Å. By using compounds such as $\operatorname{Ge}_{x}\operatorname{In}_{1-x}\operatorname{As}$ it should be possible to obtain coherent radiation at any wavelength between 8400 Å and 31,000 Å. Other mixtures have similar interesting possibilities.

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4. Parametric Devices R. Rafuse*

1. Parametric Amplifiers and Other Small-Signal Devices

In the period 1960–1963 parametric circuits have been extensively developed. Perhaps the most important breakthrough was the operation of parametric amplifiers at liquid helium temperatures. Blake [1963] and others at M.I.T. Lincoln Laboratories reported 10 °K excess noise temperature at L-band and Hanson, Fink, and Uenohara [1963] at BTL reported a 2 °K amplifier temperature (exclusive of input losses) at C-band. In both cases the

diodes were point-contact GaAs varactors of the Sharpless [1961] type. Parametric amplifiers of this type are more than competitive with masers. They have wider bandwidth capabilities, higher dynamic range, simpler cryogenics (not temperature sensitive, as are masers), no magnetic field requirements and equivalent noise performance.

The problem of bandwidth in parametric amplifiers and frequency converters has been attacked by Kuh [1962] and Helgesson [1962] with the result that limits on bandwidth and gain-bandwidth have been derived. These limits are fundamental in character and do not necessarily give synthesis schemes for good approximations with finite networks. Continuing experimental work by Gliden and Matthaei

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[1961], Johnson [1961], Little [1961], Schaffner and Voorhaar [1961], and Vincent [1961] and many others has demonstrated that wideband amplifiers and frequency converters can be designed and constructed.

A fairly complete analysis of the various smallsignal devices has been carried out and compiled by Penfield and Rafuse [1962]. They cover the various upconverters, downconverters, and parametric amplifiers for both noise figure and gain. The case of the degenerate parametric amplifier is covered in some detail with particular attention paid to its performance with various detectors. In particular, as de Jager and Robinson [1961] have also shown, the "radiometric" noise figure does not give the proper signal-to-noise ratio for radiometers when written in terms of the post-detector fluctuations; a degenerate parametric amplifier is a factor of $\sqrt{2}$ worse in performance (in a radiometer) than would be predicted from simple noise figure measurements.

The maximum frequencies at which parametric amplification with low noise can be obtained have not risen appreciably during the past three years. Several efforts are now under way to develop 35 Gc/s amplifiers, but the necessary breakthrough in diode cutoff frequency is just beginning. Occasional GaAs point-contact units show cutoff frequencies of 1000 Gc/s, but very recent developments with InSb show great promise. The first InSb varactors to be made have cutoff frequencies as good as the best GaAs units and the next three years may see the advent of millimeter-wave varactor parametric amplifiers.

Interest is now developing in the double-sideband up-converter as a convenient means for achieving ultra-low noise audio- and video-frequency amplification. Analysis of the difficult noise and gain characterization is being carried out at M.I.T. Optimizations of the two-output interface between the upconverter and the following amplifier-synchronous detector combination are being considered with respect to achieving the best noise performance.

There are many other problems under investigation including low-frequency (below the signal) pumping in parametric amplifiers, phase stability in parametric amplifiers and upconverters, and dynamic range. The problem of dynamic range has been severe until very recently. It is now recognized that extremely high dynamic ranges (in excess of 140 db) can be obtained by suitable varactor circuits. A quick analysis of the technique can be found in the work by Penfield and Rafuse [1962] and indicates that amplifier and receiver technology has seen a millionfold decrease in susceptibility to distortion and cross-modulation effects.

2. Frequency Multipliers

Within the past 3 years, all-solid-state microwave power sources have become an operating reality. Solid-state power outputs of 1 w at X-band are now attainable with 5 w at L-band already available. In fact, transistor technology has probably been pressed hardest by the demand for the high power at VHF necessary to drive varactor multiplier chains well up into the microwave region. Reports such as the one by Baldwin, Collins, Johnson, and Priest [1963] of 50 mw at 13.3 Gc/s are common.

Theoretical analyses now exist which make rational design of frequency multipliers possible. Diamond [1961] and Penfield and Rafuse [1962] have derived the equations for charge-controlled frequency multipliers and dividers. The efficiencies, power outputs, and input and load impedances for many multipliers, up to and including the octupler, are given as functions of signal frequency, diode cutoff frequency, breakdown voltage and minimum junction capacitance. It was shown that nonoverdriven abruptjunction multipliers of orders higher than the doubler must have "idling" currents flowing at one or more intermediate harmonics.

The graded-junction doubler was analyzed by Greenspan [1962] and found to be somewhat less efficient and of considerably less ($\frac{1}{2}$) power handling capability. However, very recent results by Davis [1963] indicate that the overdriven varactor multipliers, utilizing the failure of the forward-conduction mechanism at high frequencies, can handle much more power than the nonoverdriven multipliers and with a rise in efficiency. Perhaps the most important result occurs with the graded-junction which, in the overdriven case, can have a higher efficiency (at least as a doubler) than the equivalent abrupt junction.

The problem of bandwidth in frequency multipliers is still essentially unsolved. A first step in the problem, an analysis of the "bandwidth" limits of an abrupt-junction doubler is now under examination at M.I.T. The word "bandwidth" here is of a much more restrictive sense than normally used in linear circuits. It refers to the maximum range of frequencies over which the driving source for the doubler may be *slowly* tuned without appreciable degradation in efficiency. No attempt at any concept of "instantaneous bandwidth" has been made for such nonlinear devices.

Spurious signals in multiplier chains still remain a serious problem. McDade [1962] has carried out a simple analysis of one important mechanism with good results experimentally. However, other mechanisms of more complicated nature still produce spurious outputs and have to be dealt with largely on an experimental basis, circuit by circuit.

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5. Millimeter-Wave Techniques

A. Karp^{*} and J. Allison^{*}

1. Introduction

In the past 3 years, the number of people engaged in millimeter-wave research, development, production and sales, and the number and performance of different devices developed, have been very much greater than in the previous 3-year period. However, the number of essentially new ideas has not increased proportionately. Moreover, workers in the field are aware that the demand for their products remains discouragingly low, perhaps in a vicious cycle in which low demand keeps prices high and vice versa. A recent example is the shelving of plans by the Bell System for a cross-country millimeter-wave "pipeline" when it was found to be cheaper to install low-frequency coaxial cables, albeit in the vast multiplicity needed for the same total bandwidth the "pipeline" would have provided. New, essential applications, where only millimeter waves can do the job, are urgently desired.

The most recent, and broadest, conference in this field was held in January 1963, at Orlando, Fla. (Some of the papers presented are represented in this report.) The most inclusive survey paper available is Coleman's [Coleman, 1963], covering the state of the art up to the opening of this conference.

2. Generation; Amplification

Conventional Tubes

Progress in the development of magnetrons, klystrons, and traveling-wave tubes towards higher performance and frequency has been due to the refinement of mechanical techniques and electron optics, rather than the introduction of essentially new concepts. Increased power output, as well as frequency, is largely a result of new, higher beam voltage levels, as predicted in a millimeter-wave tube survey paper [Karp, 1960] available at the opening of the present report period. A number of simpler tubes have been commercially available for some time [Serchunk, 1962]. A noninclusive sampling of recent developments follows, and is an index of current trends.

A 70 Gc/s reflex klystron with 1 w output was produced at Varian Associates (private communication). Development of millimeter-wave coaxial mag-

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netrons [Olson and von Ohlsen, 1961] is being carried on by Westinghouse (private communication), but the S.F.D. Company is now leaning towards the inverted magnetron (private communication). Ingenious helix support methods contributed to the success of a half-watt CW helix traveling wave amplifier (T.W.A.) in the 50 to 60 Gc/s region [McDowell, 1960; Melroy, 1961]. Workers at Watkins-Johnson have managed to focus beams through 0.005 in. I.D. helices intended for low-noise T.W.A.s at 70 to 85 Gc/s (private communication). At the level of tens of milliwatts, production of backward wave oscillators (BWO's) with the simple two-dimensional slotted-sheet structure—usually photo-etched—continues at several firms [Melroy, D. O., 1961]. The most recent entries are Sperry's with a present upper limit of 110 Gc/s at 4 Kv [Noland et al., 1963].

At higher levels of power and cost, the folded waveguide (or single-hole-coupled cavity) structure, usually fabricated of stacked disks and washers, is a preferred circuit. It is used in a Varian B.W.O. [Schumacher, 1963] (45 to 77 Gc/s, 1.4–8.8 kv; up to 11.5w) in a Hughes B.W.O. using the *third* space harmonic (50 to 60 Gc/s, 5 to 10 kv; 18 w), and in a power amplifier using the second (51 to 56 Gc/s; 150 w at 30 percent efficiency with depressed collector (private communication)), [Forster et al, 1962; Forster et al., 1962–63; Forster, Haney, 1963].

A structure of simple vanes, formed by multiple slotting, with a round beam tunnel, spark eroded, is used in some newer Hughes oscillators (private communication) [Forster et al., 1962; Forster et al., 1962–63; Forster, Haney, 1963] (50 to 56 Gc/s; 5 to 10 kv, 5 to 30 w at 50% efficiency with depressed collector, and 140 Gc/s, 5 w peak at 25% duty factor.)

General Electric has a pulsed (80 kv at 1 a) amplifier at 35 Gc/s with 15 kw peak output [Enderby]. The wave structure consists of a tunnel of rings supported between two metal edges. A 1 kw, CW, coupled-cavity T.W. amplifier at 35 Gc/s is also reported (private communication). At 95 to 101 Gc/s, Watkins-Johnson have demonstrated 100 kw peak and 50 w average B.W.O. powers; at 90 to 200 kv (at 3.5–11a.) only moderate disk loading of the waveguide, as in a linear accelerator, is required [Johnson et al., 1962].

Other Traveling-Wave Tubes

Recently traveling-wave tubes have been studied in which electrons execute periodically structured trajectories inside *smooth* waveguides. Helical trajectories [Chow, Pantell, 1960] (with a longitudinal magnetic field of the cyclotron resonance value) have not yet shown any ability to work in the millimeter band. There may be an inherent limitation [Dickson et al., 1961]. Cycloidal trajectories, in a region of crossed steady fields, though demonstrated to be useful in the millimeter region by Reddish in England and Guematsu in Japan, have not yet been tried in the United States. The use of sinusoidal trajectories established by periodically reversing externally applied magnetic fields, [Phillips, 1960] however, has led to impressive results in the extra high-voltage domain where the pitch of the magnetic structure is reasonable. The preferred arrangement uses a TE_{01} circular-electric waveguide concentric with a hollow beam within which electrons oscillate azimuthally. Over 1 Mw peak power was obtained at 16 Gc/s; at 50 to 75 Gc/s, power output is limited by breakdown of the external, pressurized waveguide [Phillips, 1963].

Development of a d-c pumped quadrupole (transverse wave) amplifier for 70 Gc/s is under way at S.F.D. [Feinstein, 1962].

Other Schemes Using Thermionic Electrons

The following schemes for generating signals in the range 30 to 3000 Gc/s generally differ from the foregoing by existing on a laboratory demonstration rather than a device basis. For brevity, discussion will be confined to the schemes demonstrated in the United States since 1960.

In one general category, tight bunches of electrons are formed and accelerated to megavolt levels so that they may excite a structure through which they pass. The structure is required to have a finite coupling, or radiation, resistance at the desired minimum-wave frequencies, which are also harmonics of the bunching frequency. For example, the bunched beam of a 9.3 Gc/s-3.5 Mev linear accelerator was undulated by an array of magnets so as to be able to couple to the ordinary waveguide surrounding the beam, and harmonics up to 750 Gc/s were observed [Mallory et al., 1963]. (When the waveguide itself is sinuous, it is not necessary to undulate the beam (private communication). Bunched beams have also been used to excite TM_{onm} and other cavities [Sirkis et al., 1961]. Of much interest is the excitation of Cerenkov radiation in various media surrounding the beamusually dielectrics or ferrite [Rosenbaum, Coleman, 1963; Coleman, Enderby, 1960; Hakki, Krumme, 1961; Erteza, Newman, 1962; Jelley, 1962]. When such bunches are annihilated [Hakki, Coleman, 1961] or undergo other transitions, they can be made to radiate some of their energy. Sinusoidal deflection modulation of an unbunched beam of finite diameter causes bunches to appear to an observer watching the beam cross a plane at an angle to the mean velocity. These electrons can couple to the E fields of waves in a dielectric-loaded interferometer oriented at this angle [Baird, Coleman, 1963].

In another category, Weibel's scheme for getting a pulse of radiation by quickly applying a strong magnetic pulse to a trapped electron cloud has been demonstrated [Dressel et al., 1962].

Harmonic Generation and Parametric Devices

The mixing of low-frequency signals, or the generation of their harmonics, in nonlinear media has long been a much studied means of obtaining higherfrequency signals. Extension of results to the millimeter range, however, has been very difficult [Coleman]. A possible exception, at least at low power levels, is the use of a point-contact diode as the nonlinear element, [Ohl et al., 1959; Wentworth et al., 1963]. With a proper choice of base and whisker materials, production of a usable signal at a frequency about 400 Gc/s is not considered exceptional.

Nonlinear-capacitor, semiconductor junctions are perhaps the next contender [Heilmeier, 1960; DeLoach, 1960]. The list of other nonlinear media under study includes gaseous discharges, electron swarms (as in magnetrons), magnetic materials, ferroelectrics, superconductors and field emitters [Coleman]. So far only a few reductions to practice at millimeter wavelengths have been reported [Ayers, 1959; Heller, 1961; Moore, 1962; Knight, Walsh, 1962].

Masers, Quantum Devices, Etc.

In reviewing progress in the generation or amplification of millimeter waves under this heading, mention may be made of the Esaki diode: Burrus and Trambarulo [Burrus, 1960; Burrus, Trambarulo, 1961] demonstrated that with suitable materials signals of frequency higher than 100 Gc/s—albeit at very low power level—could be generated. Shaw's scheme for getting a pulse of radiation by applying a rapidly rising, strong magnetic pulse to a small sphere of single-crystal ferrimagnetic material, (YIG) has been tested. The energy is stored in the precessional motion of electron spins. Operation beyond 50 Gc/s with a peak power of tenths of a watt has been reported (private communication) [Elliott et al., 1961]. Attempts to extend the microwave frequencies observed in the beating of intense optical radiation within, or at the surface of, various media are continuing [Coleman].

Several solid-state masers, based on paramagnetic ions in a suitable host crystal, have been reported operational in the millimeter band [Coleman, Momo, et al., 1960; Foner, 1960; Carter, 1961; Devor, The modes of operation include those 1963]. requiring pulsed and steady magnetic fields, and pumping at lower microwave, higher microwave. and optical frequencies. Stimulated emission at the 88.63 Gc/s transition in hydrogen cyanide gas has also been observed [Marcuse, 1961; Barnes, Maley, 1961]; Marcuse's apparatus, for example, included a Fabry-Perot resonator into which the gas is admitted via an electrostatic state selector.

3. Passive Components

Transmission Systems

The various types of metal-walled waveguide now used at millimeter wave frequencies, for example TE₁₀ mode circular guides and oversized guides, were mostly developed prior to 1960 and more recent work has been concerned with increasing their oper- | response covering infrared as well as millimeter wave-

ating frequencies to above 100 Gc/s [Martin, Karbowiak, 1960; Valenzuela, 1963; King, 1963]. The usefulness of other guiding structures, for example dielectric-loaded, parallel-plane and trough lines, [Cohn, 1960; Cohn et al., 1960; Tischer, 1963] and block-loaded guides [Mims, 1960] have also been studied. Single-wire lines, both coated and uncoated [King, Wiltse, 1962] and double-wire lines [Ishii, 1961 have been used successfully at these frequencies.

Several millimeter-wave optical transmission systems have been studied recently. Sobel et al. [Sobel et al., 1961], use horn antennas in conjunction with large diameter phase-reversing Fresnel zoneplates to obtain transmission losses of only 2 db over 55 ft at 1.4 mm wavelength. Transmission at 4.25 mm wavelengths between parabolic antennas with separations up to 100 ft gives losses about twice those obtained by Sobel [Fellers, 1962]. Goubau [Goubau, Schwering, 1961; Christian, Goubau, 1961] has investigated a system of horns and periodically spaced phase-transforming lenses to guide a beam with theoretically arbitrarily small loss. In practice the attenuation can be less than 1 db per 1000 ft at 8 mm wavelengths.

Application of quasi-optical principles to individual components is represented by an attenuator which couples the evanescent waves near the oblique faces of two prisms with variable separation [Raker, Valenzuela, 1962; Taub et al., 1963]. The use of gratings has also been reported [Mallory et al., 1963].

Resonators

Fabry-Perot resonators, while not new in principle, have been intensively studied recently both as interferometers, and with a view to integrating them into other optical and semioptical systems [Culshaw, 1960; 1961; Fox, Li, 1961; Zimmerer, 1962; Culshaw, 1962]. Confocal resonators, because of their lower diffraction losses for a given size, may become important in the millimeter wave region [Boyd, Gordon, 1961; Boyd, 1961; Boyd, Kogelnik, 1962].

Ferrite Devices

Thanks to new materials (such as hexagonal ferrites, whose anisotropy fields facilitate getting ferromagnetic resonance at the higher frequencies) as well as to the extension of centimeter-wave principles, isolators, circulators, etc., have been developed for frequencies up to 100 Gc/s [Thaxter, Heller, 1960; Barnes, 1961; Rodique, 1963; Harrison et al., 1963]. A representative commercial 4-port device (Watkins-Johnson) may be used as a circulator, isolator, band-pass or -reject filter, etc., at 70 to 85 Gc/s. It uses two 0.014 in. diam. YIG spheres in crossed waveguide. Applied magnetic induction is 30 kilogauss and residual insertion loss is under 2 db.

4. Detectors

Until recently, the only practical detectors with

lengths were thermal types [DeWaard, Wormser, 1959; Goodwin, Jones, 1961; Long, Rivers, 1961; Arnold et al., 1963]. More recently, an InSb photodetector covering this range has been developed in the U.K. [Putley, 1960] and the process believed to be responsible for its operation is discussed by Rollin [Rollin, 1961]. As a low-level mixer and video detector for the lower millimeter wavelengths, the crystal diode is most frequently used. Gallium arsenide has been demonstrated to have superior properties to either silicon or germanium for this purpose [Sharpless, 1961]. Application of d-c bias can increase the gain and sensitivity while reducing the noise [Ishii, Brault, 1962]. Tunnel diodes in the form of gallium-n germanium point contacts, operated on the reverse portion of the I V, curve have yielded good performance up to at least 300 Gc/s [Burrus, 1963].

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6. Status of Gas Lasers

W. R. Bennett, Jr.*

Since the announcement [Javan, Bennett, and Herriott, 1961] of the first successful helium-neon gas laser in February of 1961, developments in this field have continued to move at a rapid pace. A detailed review of material through December of 1962 has been given previously by the author [Bennett, 1962] and more recent compilations are to be contained in the Proceedings of the Third International Conference on Quantum Electronics [Third International Symposium, 1963] and the Symposium on Optical Masers held at Brooklyn Polytechnic Institute [Optical masers Symposium, 1963].

As of the present writing (May 1963), continuous oscillation has been obtained on over 150 transitions, [Javan, Bennett, and Herriott, 1961; Rabinowitz et al., 1962; Patel et al., 1963] ranging from the visible $(0.59 \ \mu)$ [Gordon et al., 1963] to the far infrared $(32.5 \ \mu)$ [McFarlane, 1963]. Oscillation has been obtained in over a dozen gas systems, using at least four different excitation mechanisms. The vast majority of cases has been found in the pure noble gases (or mixtures of noble gases) through use of electron impact as the dominant excitation process. Oscillation has been obtained on some two dozen transitions in the helium-neon system in which near resonant transfer of energy from the $2^{3}S$ and $2^{1}S$ metastable states of helium has been used in collisions with ground state neon atoms to populate the upper neon maser levels. So far, the helium-neon system has been the only one reported in which laser oscilla-

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tion has been produced through this inelastic collision process. However, and somewhat ironically, oscillation has still not been obtained on the transitions initially proposed in the helium-neon case [Javan. 1959 and 1960]. A somewhat more general excitation method using metastable carriers of energy has been found to involve molecular dissociation. The first of these systems is quite representative and makes use of neon metastables to dissociate ground state oxygen molecules, leaving one constituent atom in the upper maser level [Bennett, 1962 a, b, 1963]. Here the existence of three particles in the final state permits obtaining large transfer cross sections (in the order of 10^{-15} cm²) even in cases involving large (1 or 2 ev) energetic resonance discrepancies. Oscillation has since been reported in approximately half a dozen cases involving molecular dissociation [Bennett, 1962a and Patel et al., 1963]. The existence of an anomalous frequency shift in the maser oscillation from the center of the spontaneous emission line is presumably explained by the presence of unidentified molecular absorption bands in the two oxygen masers [Bennett, 1962a and 1963]. Similar difficulties may exist in the other molecular systems, and it seems probable that the noble gas systems are likely to be more suitable for purposes of precise measurement. In contrast to the solid state field, only two cases (both in cesium) have been reported in gas lasers in which optical pumping has been used as an effective excitation process [Rabinowitz et al., 1962 and 1963]. The transitions that have worked through optical pumping have both depended on a fortunate coincidence of the 3889 Å resonance line of helium with one of the resonance lines in cesium. To date, oscillation has only been obtained on atomic transitions and the use of molecular transitions is yet to be explored [Muller et al., 1963].

The gain constants for transitions from analogous excited states vary roughly as the third power of the wavelength in the short wavelength region. Such a rough dependence is to be expected so long as the maser transition is the dominant one from the upper state [Bennett, 1962b], and is at least crudely obeyed through about 3 to 4 $\mu.$ Hence exponential gain constants $\left(\frac{1}{I}\frac{dI}{dx}\right)$ have been reported ranging from 0.1 per meter at 1.15 μ in neon [Bennett, 1962b] to as high as 10 per meter at 3.5 μ in xenon [Paananen et al., 1963]. No gain measurements have been reported in the extreme long wavelength region as yet; however, the gradual decrease in branching ratio in the far infrared will probably result in considerable departure from the λ^3 dependence. In spite of the smaller branching ratios at extremely long wavelengths, however, many likely possibilities for oscillation exist out to nearly 1 mm and it appears probable that the optical-microwave gap in coherent sources of oscillation will be closed in the near future. In contrast to the gain dependence on wavelength, the power available varies roughly as the reciprocal of the wavelength and for

practical dimensions has been limited to some tens of milliwatts in CW systems. Power outputs of several watts have been obtained on a pulsed basis from the helium-neon system at 1 μ [Byerly et al., 1963]. Some hope exists for obtaining larger power outputs through the use of photodissociation in systems currently being studied [Gould, 1963], and through use of power amplifiers in existing systems.

The only precise studies of the spectral characteristics of gas lasers have been conducted with the strong 1.15 μ transition of neon in the helium-neon system. The gain in this system is so small that extremely high Q cavities are required, with the result that the oscillation frequency is determined by the cavity resonance frequencies to first order and hence varies directly with the dimensions of the laser. The high degree of inherent spectral purity therefore loses much of its practical value except in cases where one is primarily interested in the cavity properties or in cases where negative feedback stabilization signals may be derived. In the heliumneon system, internal beat experiments [Javan et al., 1963] have indicated that the inherent spectral width of the oscillation is less than 2 c/s and beats between different lasers have indicated a short term stability good to about 20 c/s under ideal laboratory conditions [Jaseja et al., 1963]. Of more importance for communication purposes, it has been demonstrated [Rabinowitz et al., 1963], that servo mechanisms may be used to lock one gas laser on another to within a few kilocycles on a long term basis. Also, beats between opposite running waves in rotating ring lasers have been used to measure absolute rotation [Macek et al., 1963].

Studies of mode-pulling phenomena in the heliumneon system have shown the existance of a variety of nonlinear frequency-dependent mode-pulling effects which might be used to stabilize the cavity resonances on the center of the Doppler broadened line [Bennett, 1962 and 1963; Gould et al., 1963].

In particular, one such null-balance method should permit absolute stabilization in single mode operation to several kilocycles. [Bennett, 1963; Gould et al., 1963]. The high gain available in the long wavelength region may also permit the construction of gas lasers in which the cavity resonance widths are large compared to the Doppler widths and in which oscillation occurs to first order on the center of the spontaneous emission line see Bennett in TRG final report available from ASTIA.

Theoretical studies of cavity mode properties have been given in references [Boyd and Kogelnik, 1952; Boyd and Gordon, 1961; and Fox and Tingye, 1961] and have been confirmed to at least first approximation by experimental studies in reference [Bennett, 1962; Javan, Bennett, and Herriott, 1961; Javan, Ballik, and Bond, 1962; Herriott, 1962; and Rigrod et al., 1962]. Extensive calculations of the frequency and power dependence on gas laser parameters have been given in references [Bennett, 1962b, 1962c, and 1963; Javan et al., 1962; Lamb, 1963; and Tang and Statz, 1962]. Related experimental studies have been reported in references [Bennett, 1962c; Javan et al., 1961; Javan et al., 1962; Rigrod et al., 1962; Paananen et al., 1963; McFarlane et al., 1963; and Javan, 1963].

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7. Noise in Masers

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The role of spontaneous emission as a source of maser noise had been recognized in the study of maser noise at microwave frequencies. These studies also had revealed that spontaneous emission would be the predominant noise in linear amplification at optical frequencies. Part of the recent work on optical maser noise therefore was concerned with refinements of the analysis of spontaneous emission. Because the spontaneous emission output noise power of a linear amplifier within the frequency band B cannot be reduced to less than $h\nu BG$ (where h is Planck's constant, ν is the frequency at center band, G is the amplifier power gain) it was natural to search for a connection between this minimum noise power and the uncertainty principle. The definition of noise and its relation to coherence of a waveform was also studied. Finally, the first experiments on the frequency stability of optical masers were reported. Preliminary measurements on the amplifier noise have been made.

Spontaneous Emission Noise

The work on spontaneous emission in masers was concerned with refining the physical models of the

active maser medium and the mathematical treatment of the electromagnetic field [Senitzky, 1960, 1961, 1962; Haus and Mullen, 1963]. The complete probability distribution of the signal and noise amplitudes at the output of an optical maser amplifier has been obtained [Louisell, 1963] with the aid of a simple mathematical model for the maser which incorporates the main characteristics of the amplifier. The noise has been shown to be additive Gaussian noise. A study of quantum effects in parametric amplifiers at optical frequencies has arrived at very similar results [Louisell, Yariv, Siegman, 1961; Louisell, Gordon, Walker, 1963]. The similarity of the results is not coincidental. This has been demonstrated by studies which showed that a minimum additive noise is necessary in ideal linear amplifiers [Serber and Townes, 1960; Heffner, 1962; Oliver, 1961; Haus and Townes, 1962] in order to satisfy the uncertainty principle. In this connection the studies on the information-theoretical implications of maser noise are also of interest [Gordon, 1962]. Finally, spontaneous emission noise has been shown to set an ultimate limit to the frequency stability achievable in optical maser oscillators [Townes, 1961].

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Coherence

The concepts of noise and coherence and their quantitative measures have been used for years in electrical engineering on one hand, and optics on the other hand. The advent of the optical maser put these into new perspective and several papers have appeared that deal with these questions, R. J. Glauber [Glauber, 1963; Glauber, to be published in Phys. Rev.] formulated measures of wave coherence that are consistent with quantum field theory. He defines a set of normalized correlation functions which are equal to or less than unity. A waveform is defined to be perfectly coherent if all correlation functions are equal to unity. These new correlation functions introduced by Glauber have been discussed in their classical limit by Wolf [1963] and further interpretations have been attempted [Sudarshan, 1963].

Measurements

Measurements on the frequency stability of the beat-note between two axial modes in a gaseous He-Ne optical maser oscillator have been reported [Javan, Ballik, and Bond, 1961]. The observed frequency variation, referred to the optical frequencies of the modes, corresponds to a stability better than one part in 10^{14} . Since the "microphonic" perturbations shift the frequencies of both modes roughly by the same amount this kind of stability was not observed in the beat-note between modes of two different masers. In a carefully controlled environment, short-time frequency variations of the order of 20 c/s were observed [Jaseja, Javan, and Townes, 1963] corresponding to a stability of 8 parts in 10^{14} . This stability is still far from the ultimate limit on the line-width set by spontaneous emission [Townes, 1963]. The observed stability approaches in order of magnitude the value that was to be expected in these experiments due to the thermal vibrations of the supporting rods.

The measurements on noise in optical maser amplifiers are of a more preliminary character. Measurements 1 on a traveling wave He-Ne maser at a wavelength 3.39 μ indicate that the noise is not much greater than the theoretical minimum 6120 °K (expressed in terms of an effective input noise temperature).

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8. Nonlinear Optical Properties of Solids

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A general discussion of nonlinear optical properties of solids should distinguish between what might be called coherent and incoherent nonlinear processes. The absorption of radiation and subsequent electron emission from a photosensitive surface, for example, depends on the square of the field amplitude and thus could be called a nonlinear effect of the incoherent type. A second effect of this same type but which more clearly deserves the nonlinear designation is the absorption of a photon (i.e., a pump photon) at one frequency and the subsequent emission of a lower energy photon in laser emission. These are called incoherent since the nonlinear optical phenomena is accompanied by a real transition or change of the material system.

The coherent type of nonlinear phenomena is best illustrated by the pioneering work of Franken and coworkers [Franken, Hill, Peters, and Weinreich, 1961]. A large electric field was obtained at an optical frequency by focusing the parallel beam of light emitted from a ruby laser down to a small spot within a piece of quartz. It is well known that materials respond to electric fields by becoming polarized and, if the material lacks a center of inversion symmetry, part of this polarization will be proportional to the square of the electric field. This is just the reason why crystal rectifiers, mixers, and harmonic generators work at lower frequencies and for the very same reason Franken was able to observe the blue light generated as the second harmonic of the red laser light. In this experiment, as in all the coherent nonlinear effects, the material plays a passive role. For every blue photon generated there are exactly two red photons annihilated and electromagnetic energy is exactly conserved. The nonlinear element, the quartz, is unchanged in the process and there is a very definite phase relationship between the incident red beam and the emitted blue beam. One describes this process as a coherent nonlinear phenomena.

These examples are extreme cases and can be assigned the coherent or incoherent designation without any ambiguity. In fact, however, there are phenomena that fall in between and cannot clearly be assigned to either of these two groups. Consider the case of second harmonic generation in a material that has a natural absorption near the second harmonic frequency. If the material should make an upward transition, the question arises as to whether it should be considered as if a second harmonic photon has been created and then absorbed (a coherent nonlinear process followed by linear The first observation of beats between two different light sources followed soon after Franken's original experiment. Bass, Franken, et al. [1962], beat a ruby laser at room temperature and one cooled to liquid nitrogen against one another and observed the second harmonic due to each of them as well as a signal at the sum of their two frequencies.

It is immediately clear to anyone studying these effects that they can be considerably enhanced if the interacting waves all have the same speed. In the first paper on the subject, Franken mentioned that failure to satisfy this condition limited the size of the effect he could detect. Giordmaine [1962] and Maker et al. [1962], simultaneously demonstrated how one could satisfy this condition of phase matching in practice in KDP. In addition, Maker et al. [1962], demonstrated the effects of phase mismatch in quartz. Following these first experimental results, several theoretical papers appeared on the subject. Braunstein [1962] obtained expressions for the nonlinear response of semiconductors and Kleinman [1962] commented on certain phenomenological aspects of the nonlinear tensor. Kroll [1962] suggested how the nonlinearity might be used as a parametric amplifier or oscillator. The most extensive discussion of optical nonlinearities however was given by Armstrong et al. [1962]. They properly incorporate the nonlinearity into the macroscopic Maxwell's equations for dense dielectric media and gave exact solutions to the simultaneous nonlinear differential equations that describe second harmonic generation, the beating of two waves to form either a sum or difference frequency signal, and the general problem of higher harmonics. These solutions are exact in the sense that they properly describe the situation when so much power is generated at the second harmonic, for example, so that the energy of the fundamental frequency wave is significantly depleted. In addition, this paper obtains certain dispersion-symmetry relations between different elements of the nonlinear tensor. There is a fundamental relation between these symmetries and the well-known Manley-Rowe relations.

absorption) or if two photons at the fundamental were absorbed directly (an incoherent nonlinear process) as in the experiments of Kaiser and Garrett [1962]. In fact, it can be either of these two depending on how near the second harmonic is to the atomic resonance frequency. For the remainder of this paper, we will restrict ourselves to phenomena that are clearly of the coherent type. The review articles of Bloembergen [1963] and Franken and Ward [1963] discuss this middle ground at greater length.

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Bloembergen and Pershan [1962] extended the previous solutions for infinite nonlinear media to include effects due to the boundaries of the material. Nonlinear generalization of the Fresnel formulas are obtained. Kleinman [1962] made further theoretical observations and in particular discussed the experimentally significant geometry of harmonic generation at the focus of a lens. Siegman [1962] suggested how the nonlinear effect could, in principle, be used to build a power limiter that would suppress the spikes observed in many of the solid state lasers.

Terhune, Maker, and Savage [1962] observed second harmonic generation in calcite by applying an external d-c electric field. Without this field, the crystal is centrosymmetric and the only second harmonic generation is due to a quadrupole effect which they observed. In this same paper, they reported the first observation of third harmonic generation.

The dispersion-symmetry relations mentioned previously [Armstrong et al., 1962] state that the constant for rectification of light should be directly related to the linear electro-optic constant in the same material. Bass, Franken, et al. [1962], observed the ratio of these constants for ordinary KDP and deuterated KDP and, consistent with theory, found them to be the same. Along the same line, Smith and Braslau [private communication] were able to beat the 3115 Å line of a mercury lamp with a beam from a ruby laser (6943 Å) and observe the 5650 Å beat. To the limits of their accuracy the constant for this process was not different from the constant for beating a 5460 Å mercury line against the ruby to obtain a 3056 Å sum frequency signal [Smith and Braslau, 1962].

Difference frequency beats at a microwave frequency might prove to be of significant value in demodulating light. M. DiDomenico et al. [1962]. observed 1.3 Gc/s beats between adjacent axial modes of a ruby laser when its output is focused on a cadmium selenide crystal that is biased by a large d-c electric field.

The inverse effect of beating a microwave signal with a light in order to modulate the light wave has received much attention [Pershan and Bloembergen. 1961]. The most significant experimental work on this topic has been done by Kaminow and is well documented in his review article [Rigrod and Kaminow, 1963].

A second experiment suggested by the results of Armstrong et al. [1962], was very beautifully done by Terhune, Maker, and Savage [1963]. Using a focused beam from a ruby laser they were able to convert 20 percent of the incident red light into second harmonic blue light and observe the saturation of the effect as further increase in red intensity gave no further increase in blue light.

Further data on harmonic generation in other crystals [Savage and Miller, 1962] and from a different laser, CaWO₄, N_d⁺³ [Miller and Savage, 1962], has been obtained by Miller and Savage.

Most recently Pershan [1963] has given general phenomenological arguments in order to generalize the dispersion-symmetry relations, previously given [Armstrong et al., 1962] for electric dipole nonlinearities to include magnetic and quadrupole nonlinearities as well. Price [1962] has been able to obtain certain integral equations relating real and imaginary parts of the nonlinear tensors that are analagous to the well-known Kramer's-Kronig applicable to linear systems.

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9. Highlights of Progress on Satellite Devices

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All of the significant "real time" satellite communication experiments so far performed have been carried out during slightly more than the past three The first were Project SCORE [Brown, 1960], vears. and [Brown and Senn, 1960] and COURIER [Siglein and Senn. 1962] which were designed for delayed repetition. Next came Echo [Echo, 1961], which used a 100-ft aluminized mylar balloon as a passive reflector for transcontinental communication of speech bandwidth signals. Echo was launched on August 12, 1960 and still returns a usable signal. Plans and preparations for a larger and stiffer balloon satellite are in the making [Reiger, 1963]. There have been sporadic newspaper references to nonoriented passive satellite having gain, but such possibilities seem rather remote [Cutler, 1962].

The second experiment used the active Telstar satellite [Telstar, 1963], which successfully transmitted television across the ocean over a period of four months before radiation effects caused a serious malfunction. The Relay active satellite [Wilmotte, 1962] soon followed and has functioned intermittently in a similar capacity. A fourth communication satellite, Syncom [Meisels, 1962], intended to operate in an approximately synchronous orbit suffered an abortive launch but is expected to be in orbit soon.

An attempt was made to orbit a different sort of passive satellite, Project Westford [Morrow, 1961 and 1962], which will use a multitude of small half wave wire dipoles in a low orbit. This kind of circuit involves multiple scattering; and if broad band signals are to be transmitted, rather special modulation techniques must be used [Berg et al., 1961]. There seems to be some possibility that the presence of such an orbiting belt would interfere with other activities [Liller, 1961], but this allegation has been vocally refuted [Purcell, 1962].

Perhaps the most promising area of progress in communication satellites is in the development of attitude control systems. This problem has been discussed [Kendall et al., 1960; Roberson, 1962; and Haeussermann, 1962]. The problem of damping the oscillations in such a system may be solved by using fluids [Lewis, 1962], resonant mechanical structures [Zajac, 1962] or gyro's [Burt, 1962]. There are fundamental characteristics limiting the rate of damping, which may not be exceeded by any method [Zajac, 1963].

There are interesting possibilities in building active satellites which, although unoriented, have a useful directional radiation characteristic. Most of these are based upon the concept of the Van Atta array [Van Atta, 1959]; Hansen [1961] showed that the linear Van Atta array could profitably be used with amplification in a satellite. Davies [1963] showed that the idea could be extended to cover circular arrays.

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A wide variety of power sources for communication satellites has been discussed from time to time, but only two seem to be serious contenders at this time. [Picquendar, 1961]. Nuclear power sources hold considerable promise, but most systems are being designed around solar cell arrays. The P on N silicon solar cell has been widely used, but it has recently been found to be less resistant to radiation damage from high energy protons [Madelkorn et al., 1960] than are N on P cells. Experience with N on P cells verified the predicted performance and gave about eight times the useful life expectancy of P on N cells [Brown et al., 1963].

For many applications, it is necessary to supplement solar cells with a chemical storage battery. The Nickel-Cadmium battery is most widely used. The kinetic characteristics of the Nickel-Cadmium cell are becoming better known [Thomas, 1962] and has resulted in batteries having improved properties and much greater reliability [Bomberger, Moose, et al., 1963]. The combination of P on N solar cells and a Nickel-Cadmium storage battery was first used in the United States on the Telstar satellite [Bomberger, Feldman, et al., 1963].

At present, the only satellite function which cannot readily be performed with solid state devices is the RF amplification to reasonable power levels (several watts). The traveling wave tube is most suitable for this purpose because it can be made both efficient and long lived. Several tubes are under development. It seems possible to design tubes for 10 to 20 years life expectancy [Bodmer et al., 1963]. By using periodic magnetic focusing or electrostatic focusing. they can be made very light [Nowogrodzki, 1961; Notvest et al., 1960].

The vast bulk of satellite components are semiconductor devices and passive elements; and if a satellite is to have a good chance of surviving for many years, the reliability of these components must be extreme [Ross, 1962; Peck and Wooley, 1963].

Radiation poses a particular hazard to solid state components in satellites. One well-known type of radiation damage is a reduction of carrier lifetime because of bulk damage in the semiconducting material. A more recently discovered effect is a temporary degradation of transistor performance due to the effect of ionization of the gas in the device and resulting ionization of the surface [Peck, Blair, et al., 1963].

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10. Waves in Plasmas^{*}

A. Bers 1 and P. Chorney 2

The continuing and increasing interest in plasmas has brought about a thorough exploration of the characteristic wave properties of a plasma. For the years 1960–1962, to be reviewed here, the U.S. literature on this subject has been so enormous that we shall have to restrict this review to only a few of the many interesting aspects of this field. In the following, we, therefore, shall not review work on: nonlinear properties of plasmas, radiation from plasmas, and work exclusively related to MHD and collisiondominated plasmas.

A brief survey is given of experimental studies on interactions between plasmas, electromagnetic waves, and electron beams. It is noted that experimental activity in this area of plasma physics is on the increase, and satisfactory correlation between theory and experiment has been obtained in many cases.

1. Books and General Reviews

A large amount of the research in plasmas has been made available in various text books, research monographs, conference proceedings, and review articles.

The text books [Glasstone and Loveberg, 1960; Rose and Clark, 1961] cover a wide range of topics on the properties of plasmas, with most of the discussed applications related to the field of fusion. Detailed description of waves in plasmas has ap-

^{*}This work was supported in part by the U.S. Army, the Air Force Office of Scientific Research, and the Office of Naval Research; and in part by the National Science Foundation (Grant G-24073).

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peared in two books that are more in the nature of advanced texts or research monographs. One of these [Allis, Buchsbaum, and Bers, 1962] deals with both free and guided waves in plasmas. The other [Stix, 1962] gives an excellent description of free waves in a hot plasma. Mathematical tables [Fried and Conte, 1961] that aid in the analysis of the hot plasma dispersion equation were also published. The second edition of Spitzer's Monograph [Spitzer, 1962] also contains an expanded treatment of waves in plasmas. The proceedings of several conferences have been

edited in book form. Among these are: The Second Geneva Conference on The Peaceful Uses of Atomic Energy [Allis, 1960]; The International Symposium on Plasma Dynamics at Woods Hole 1958 [Clauser, 1960]: The Stanford Research Institute 1957–1958; Plasma Physics Seminars [Drummond, 1961]; and the Lockheed Fifth Annual MHD Symposium [Mitchner, 1961]. The proceedings of some conferences have also appeared in special issues and as supplements of journals: Proceedings of the 1959 International Plasma Physics Institute [Drummond, 1961]; proceedings of the 1960 International Symposium on Magneto-Fluid Dynamics (1960); and the 1961 Salzburg Conference on Plasma Physics (1962).

Several review articles have been devoted to the linearized description of a plasma and its wave propagation properties [Bernstein and Trehan, 1960; Pai, 1960, and Oster, 1960 and 1961].

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2. Macroscopic Description of a Plasma

Much emphasis has been placed upon determining the proper conductivity tensor for a plasma starting from kinetic theory. The d-c conductivity for a partially ionized gas in the absence of an applied magnetic field has been treated [Dreicer, 1960]; this includes inelastic collisions and the two-body Coulomb interactions. The high-frequency conductivity in a magnetic field [Kelly, 1960] has been treated with a full evaluation of the binary collision integral in the Focker-Planck approximation. A different approach to the solution of the linearized Boltzman equation is to consider the conductivity as the kernel of an integral equation [Drummond et al., 1961]; this approach is readily applicable to an inhomogeneous plasma. For a collisionless plasma the conductivity and damping (Landau) has again received some attention [Buneman 1961]. More recently attention has been directed toward evaluating the effects of collective dynamics [Dawson and Oberman, 1962; Oberman, Ron, and Dawson 1962].

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3. Free Waves

The theory of waves in an unbounded plasma is given in great detail in the books mentioned in section 1 [Allis, Buchsbaum, and Bers, 1962] and [Stix, 1962].

The Landau damping theory has again received a great deal of attention. The details of this have been given for a hot electron-ion plasma [Fried and Gould, 1961]. A very clear picture of the damping has been obtained from energy exchange considerations of multiple beams [Dawson, 1960 and 1961]. Also, the effect of (close range) collisions on Landau damping has been investigated [Platzman and Buchsbaum, 1961]. Attention has also been directed to cyclotron damping (the analog of Landau damping at cyclotron frequencies) in a collisionless plasma [Stix, 1960] and [Scarf, 1962].

Considerable attention has been given to the description of free waves from the Boltzman-Vlasov equations, and the mathematics of their solution [Lewis and Keller, 1962; Backus 1960]; see also Bernstein and Trehau, section 1.

The relativistically hot plasma has also been solved including the effects on damping [Buti, 1962] and [Imre, 1962].

The description of free waves from transport equations is by now fairly complete for isotropic pressures and no heat flow [Tanenbaum, 1961; Pai, 1962]. Extensions to the case of an anisotropic pressure have been started [Jaggi, 1962].

For the cold plasma, the oscillations with multiple ion species has been treated [Buchsbaum, 1960].

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4. Plasma Waveguides

The theory of guided waves in cold plasmas of finite extent is treated in the book mentioned in section 1 [Allis, Buchsbaum, and Bers, 1962]. Many new dispersion characteristics have been found. Detailed computations on the fast waves in an electron plasma waveguide have been made [Bevc and Everhart, 1962]. It also has been shown that backward waves and complex waves exist on an isotropic, plasma slab [Oliner and Tamir, 1962]. Computations on the complex waves in a plasma filled waveguide with a longitudinal magnetic field have also been given [Chorney, 1962].

Perturbation theory as applied to cold, anisotropic plasma waveguides and resonators has been carried out with detailed examples [Buchsbaum, Mower, 1961].

and Brown, 1960] and [Mower and Buchsbaum, 1962].

Very few boundary value problems have been solved for the warm plasma. We mention a few of the problems that have been attempted: Acoustic wave excitation for a semi-infinite plasma [Turcotte and Schubert, 1961]; ion cyclotron waves in a bounded geometry [Engelhardt and Dougal, 1962; and the Green's function approach [Weitzner, 1962].

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5. Electron Beam Waves

Considerable attention has been devoted to the study of the small-signal conservation principles for electron beam systems. Energy and power theorems have been generalized for relativistic irrotational beams and cold plasmas [Bers and Penfield, 1962]. and for nonrelativistic rotational beams [Bobroff, Haus and Kluver, 1962]. Approximate relations for energy and power in thin irrotational beams [Rigrod, 1960] have been derived. Conservation laws for interactions with beams have been put into a formal matrix form [Pease, 1960]. Studies of the conservation principles for small-signal momentum and stress have also been initiated [Sturrock, 1960], [Pierce, 19611

Analyses of wave propagation have been given for the filamentary electron beam [Siegman, 1960], and for the crossed-field, thick, sheet-beam [Hershenov, 1960 and 1961]. Other detailed studies of beam waves were presented at the International Congress on Microwave Tubes [Munich, 1960 and Hague, 1962]. Renewed attention has also been given to the effect of velocity spread in beams and the damping of space-charge waves [Berghammer, 1962].

The instabilities occurring in high density beams as well as the interaction of beam waves with circuits have been reviewed for plasma applications [Pierce, 1961]. The potential minimum instability in diodes has been analyzed in detail and its time behavior has been computed numerically [Birdsall and Bridges,

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6. Stream-Plasma Instabilities

A great deal of attention has been directed toward an understanding of the possible instabilities that may arise in plasmas.

Under the general designation of two-stream plasma systems a great variety of instabilities are These occur in non-Maxwellian plasmas found. [Penrose, 1960], in plasmas with doubly lumped distributions [Ichimaru, 1962], or in distributions with relative drift [Jackson, 1960]. Similarly, instabilities occur in contrastreaming plasmas [Kellogg and Liemohn, 1960 and Ek, Kahalas, and Tidman, 1962]. Related instabilities are the electron and ion runaway [Dreicer, 1960] and ion wave instability [Bernstein, Frieman, Kulsurd, and Rosenbluth, 1960], [Bernstein and Kulsurd, 1960], and the current carrying plasma [Bernstein and Kulsurd, 1961], [Buneman, 1962].

Instabilities can also arise because of inhomogeneities in the plasma [Frieman and Pytee, 1961] or in the magnetic field [Krall and Rosenbluth, 1961 and 1962].

Instabilities in plasmas with an injected electron beam are receiving increasing attention. Analyses for a one-dimensional, cold plasma system, in the absence of a magnetic field have been given [Neufeld and Doyle, 1961; Neufeld, 1962]. Excitation of a plasma by a beam of finite cross section has also been studied [Sturrock, 1960], and used as an amplifying mechanism for microwaves [Boyd, Gould, and Field, 1961] and [Crawford and Kino, 1961]. Analyses of

beam plasma interaction with great theoretical detail is also available [Watson et al., 1960] [Frieman et al., 1962]. Finally, a study has been made of the instabilities in an ion beam crossing a magnetic field as occurs in the Oak Ridge experiments [Burt and Harris, 1961].

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7. Experimental Investigations

Interactions between plasmas and electromagnetic waves have been applied extensively as diagnostic tools. Novel microwave diagnostic techniques have been used to determine high densities such that $\omega_p >> \omega$ [Anderson, 1961], and density profiles of transient and spatially varying plasmas have been determined by a microwave transmission technique [Wharton and Slager, 1960]. The multiple resonances associated with the Dattner-mode method of measuring plasma densities are still being investigated [Hershberger, 1960, 1961; Hershberger and Petroff, 1962]; multiple resonances have also been observed when microwaves are scattered from a plasma column with the electric field parallel to the axis [Willis and Petroff, 1962].

Experimental evidences of Landau damping have been observed. In propagation studies on ion acoustic waves in highly ionized plasmas, nondissipative damping has been seen [Wong, D'Angelo, and Motley, 1962], and rapidly damped signals have been measured on very slowly drifting electron beams [Coulton, Hershenov, and Paschke, 1962].

The wave properties of plasma columns have been utilized in experimental microwave components. These components perform as theoretically predicted. Transmission of signals through a 3000 Mc/s TE₁₁₁ cavity loaded with a coaxial plasma column has been studied; a static magnetic field was applied to the column and maximum transmission was obtained at cyclotron resonance [Olthius, 1961]. A plasma column has been used in a long coupling slot between two UHF rectangular waveguides to give a variable directional coupler; the coupler was variable from a minimum of 3 db with little attendant reflections [Willis, 1962]. The dipole resonance of a cylindrical plasma column has been used to make an electronically tunable band-pass filter in the 3000 Mc/s band [Kaufman and Steier, 1962].

Several beam-plasma interaction experiments have been reported. The interactions are basically the two-stream instability which occur near plasma and cyclotron resonances. Substantial amplification of signals in the 3000 Mc/s band has been obtained by interacting a helix-modulated electron beam with a contact-ionization cesium plasma [Allen and Kino, 1961]. Further details have been reported on earlier interaction experiments with modulated beams [Boyd, Gould, and Field, 1961]. Many workers have studied self-excited oscillations that are sustained by beam-plasma interactions. Very strong oscillations have been investigated in beam-generated plasmas [Targ and Levine, 1961; Crepeau and Keegan, 1962]. Interactions have been observed which are intense enough to produce a microwave avalanche breakdown of the gas, break up of the beam, and x rays [Getty and Smullin, 1962]. Oscillation frequencies could be identified with plasma and cyclotron frequencies. Experiments involving oppositely directed electron beams interacting with an independently generated plasma have resulted in

oscillations near plasma frequency and verification of the Bohm and Gross two-beam theory [Kofoid, 1960, 1962]. The two-stream instability has been identified in a low pressure discharge as being responsible for coherent oscillations; the frequency of oscillation has been ascertained as the plasma frequency [Putnam, Collins, and Oleson, 1961].

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(Paper 68D5-365)