

# A Note on a Generalized Elliptic Integral

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An expansion of

$$\Omega_j(k) = \int_0^\pi \frac{d\theta}{(1 - k^2 \cos \theta)^{j+1/2}}$$

in the neighborhood of  $k^2 = 1$  is obtained by a method based on an Abelian theorem.

In a recent paper Epstein and Hubbell have given a short table of the function

$$\Omega_j(k) = \int_0^\pi \frac{d\theta}{(1 - k^2 \cos \theta)^{j+1/2}}, \quad j = 0, 1, 2, \dots \quad (1)$$

which is important in certain problems in radiation physics, [1].<sup>1</sup> It is simple to find a power series expansion for  $\Omega_j(k)$  by expanding the denominator and integrating term by term. It is somewhat more difficult to find an approximation to  $\Omega_j(k)$  which is valid for  $k^2$  close to one, and such an approximation is not found in [1] except in the cases  $j = 0, 1$ . It is the purpose of this note to furnish such an approximation.

Let us use the identity

$$\frac{1}{z^{j+1/2}} = \frac{1}{\Gamma(j+1/2)} \int_0^\infty t^{j-1/2} e^{-zt} dt \quad \text{Re}(z) > 0 \quad (2)$$

to write eq (1) in the form

$$\Omega_j(k) = \frac{1}{\Gamma(j+1/2)} \int_0^\pi d\theta \int_0^\infty t^{j-1/2} e^{-(1-k^2 \cos \theta)t} dt. \quad (3)$$

It is not difficult to justify an interchange of orders of integration. The integration over  $\theta$  can then be performed by making use of the Bessel function formula

$$\frac{1}{\pi} \int_0^\pi e^{k^{2t} \cos \theta} d\theta = I_0(k^2 t)$$

where  $I_0(x)$  is a modified Bessel function.

Hence  $\Omega_j(k)$  is

(4)

$$\begin{aligned} \Omega_j(k) &= \frac{\pi}{\Gamma(j-\frac{1}{2})} \int_0^\infty t^{j-1/2} e^{-t} I_0(k^2 t) dt \\ &= \frac{\pi}{k^{2j+1} \Gamma(j-\frac{1}{2})} \int_0^\infty e^{-x} \left(\frac{1-k^2}{k^2}\right)^x x^{j-1/2} e^{-x} I_0(x) dx. \end{aligned} \quad (5)$$

We have therefore exposed  $\Omega_j(k)$  as a Laplace transform in which the coefficient in the first exponential,  $(1-k^2)/k^2$ , approaches zero as  $k^2$  approaches 1. We can now apply an Abelian theorem for Laplace transforms, [2], to determine the behavior of  $\Omega_j(k)$  in the neighborhood of  $k^2 = 1$ . To do this we note the asymptotic expansion

$$e^{-x} I_0(x) \sim \frac{1}{\sqrt{2\pi x}} \sum_{n=0}^\infty \frac{(-1)^n \Gamma(\frac{1}{2} + n)}{(2x)^n \Gamma(\frac{1}{2} - n)} \frac{1}{n!} \quad (6)$$

as  $x$  tends to infinity. Substituting eq (6) into eq (5) we find for the asymptotic expansion of  $\Omega_j(k)$ :

$$\begin{aligned} \Omega_j(k) &\sim \sqrt{\frac{\pi}{2}} \frac{1}{k^{2j+1} \Gamma(j-\frac{1}{2})} \left\{ \sum_{n=0}^{j-1} \frac{(-1)^n \Gamma(\frac{1}{2} + n)}{2^n \Gamma(\frac{1}{2} - n)} \right. \\ &\quad \left. \frac{(j-n-1)!}{n!} \left(\frac{k^2}{1-k^2}\right)^{j-n} \right. \\ &\quad \left. + \sum_{n=j}^\infty \frac{(-1)^j \Gamma(\frac{1}{2} + n)}{2^n \Gamma(\frac{1}{2} - n)} \frac{1}{(n-j)!} \left(\frac{1-k^2}{k^2}\right)^{n-j} \ln \left(\frac{1-k^2}{k^2}\right) \right\}. \end{aligned} \quad (7)$$

Another useful representation for  $\Omega_j(k)$  can be obtained by noting that the Legendre function  $P_{n-1/2}(\cosh \eta)$  can be written, [3],

$$P_{n-1/2}(\cosh \eta) = \frac{1}{\pi} \int_0^\pi \frac{d\theta}{(\cosh \eta + \sinh \eta \cos \theta)^{n+1/2}} \quad (8)$$

<sup>1</sup> Figures in brackets indicate the literature references at the end of this paper.

from which it follows that

$$\begin{aligned}\Omega_j(k) &= \frac{\pi}{(1-k^2)^{\frac{2j+1}{4}}} P_{j-1/2} \left( \frac{1}{(1-k^4)^{1/2}} \right) \\ &= \frac{\pi}{(1+k^2)^{\frac{2j+1}{2}}} {}_2F_1 \left( \frac{1}{2}, j + \frac{1}{2}; 1; \frac{2k^2}{1+k^2} \right).\end{aligned}\quad (9)$$

However, a derivation of asymptotic behavior starting with eq (9) is not as direct as the proof we have given.

## References

- [1] L. F. Epstein and J. H. Hubbell, Evaluation of a generalized elliptic-type integral, J. Res. **67B** (Math. and Math. Phys.) No. 1, 1 (1963).
- [2] G. Doetsch, Theorie und Anwendung der Laplace Transformation (Dover Reprint, 1943).
- [3] I. M. Ryshik and I. S. Gradshteyn, Tables of series, products and integrals (VEB Deutscher Verlag der Wissenschaften, Berlin, 1957).

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