# Infrared Spectrum of Hydrogen Sulfide in the 6,290-cm<sup>-1</sup> Region<sup>1</sup>

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The absorption of hydrogen sulfide in the region 6,100 to 6,500 cm<sup>-1</sup> has been measured under high resolution. The rotational fine structure has been analyzed through the use of published energy tables for the rigid rotor. A classical centrifugal-distortion correction was applied to the rigid energy levels. It is found that there are two overlapping bands in was applied to the rigid energy levels. It is found that there are two overlapping bands in this region, a normally strong, A-type band, the  $(n_1, n_2, n_3) = (1,1,1)$ , and a normally weak, B-type band, which becomes strong enough to be observed by borrowing intensity from the A-type band through a Coriolis interaction. The B-type band is the lower component of the resonating pair (2,1,0) and (0,1,2). The excited state inertial parameters giving the best fits are (1,1,1) band, A=10.398, B=8.935, C=4.548,  $\mu_0=6289.26$  cm<sup>-1</sup>; (0,1,2)band, A=10.394, B=8.918, C=4.547,  $\mu_0=6288.28$  cm<sup>-1</sup>.

## 1. Introduction

The absorption of hydrogen sulfide in the 1.6  $\mu$ region has been measured previously under low resolution and the rotational fine structure of the band envelope analyzed  $[1]^3$  As it is now possible to increase the resolution 10 times, it was thought advisable to remeasure this region. The analysis of the new data confirms the main conclusions of the previous analysis, i. e., the absorption is due largely to the A-type band (1,1,1). However, the improved data show a second and weaker band overlapping the (1,1,1) band. This second band is a B-type band and is assigned as the lower component of the resonating pair (0,1,2) and (2,1,0).

## 2. Experimental Procedure

The measurements of the  $1.6-\mu$  band of hydrogen sulfide were made on an infrared grating spectrometer constructed in the Bureau's Radiometry Section. The grating has 15,000 lines/in., and the rulings extend over 7½ in. A lead-sulfide cell is used as the detecting element. Other details of the instrument are given in a previous paper [2]. In the measurement of the hydrogen-sulfide spectra at  $1.6^{-}\mu$  it is possible to use narrow slits (0.03 mm), and sharp lines separated by 0.15 cm<sup>1</sup> could be resolved. Two cells were used in the measurements, one 60 cm and the other 600 cm in length. The long cell was so constructed that there was no loss in aperture. This cell transmits about 85 percent of incident A detailed description of this cell is given in energy. another publication [3]. Various pressures were used in the absorption cells so that both weak and strong lines could be observed. Figure 1 shows a recorder tracing obtained on unruled paper for the 600-cm



FIGURE 1. Absorption bands (1, 1, 1) and (2, 0, 1) for hydrogen sulfide in the region 6,100 to 6,500 cm<sup>-1</sup>. The calculated positions of the lines are shown above the spectra.

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FIGURE 2. Center of (1,1,1) band at slow scanning speed. The calculated positions and their intensities are placed above the spectra.

cell, with hydrogen sulfide at a pressure of 40 cm Hg. Figure 2 represents the center of the same region obtained by using a slower scanning speed.

The frequencies of the absorption peaks were obtained by using a comparison spectrum. Higher orders of the atomic lines of mercury, krypton, argon, and neon were superimposed on the absorption band. An arrangement of shutters was used so that either the absorption lines or the emission line could be recorded on the chart, and thus the shifting of position of the atomic line by the overlapping of the absorption line was avoided. Several tracings were measured, and the frequency given for each line is a result of the average of several determinations. In these measurements the precision of most lines is  $0.1 \text{ cm}^{-1}$  or better. The experimental determinations of the frequencies are listed in table 1.

TABLE 1. Comparison of observed and calculated spectra in 1.6-µ region

Line No.	<b>T</b>	Observed	rved ve ber Transition $\frac{\text{Calculated wave number}}{(111)}$	vave number	Inten- Line		Intensity	Observed	Transition	Calculated wave number		Inten-	
	Intensity	wave number		(111)	(012)	sity	No.	Intensity	number	1 ransition	(111)	(012)	sity
31	3.8	$cm^{-1}$ 6148, 77	11-10-12-11	$cm^{-1}$ 6148, 86		0.6	59	3. 3	$cm^{-1}$ 6175.87	70-8-1	$cm^{-1}$ 6175.78		0.9
		00000000	$11_{-11} - 12_{-12}$			. 2				$5_3 - 6_3$		6176.09	. 6
32	. 5	6149.51					60	. 4					
33	. 2	6150.19					61	.4	6178.15				
34	.7	6150.65					62	1.5	6178.95	$8_{-5} - 9_{-7}$		6178.92	. 6
35	. 3	6151.99								$8_{-6} - 9_{-6}$			. 2
36	1.7	6153.21				· · · · ·	63	1.9	6179.43	$6_6 - 7_5$	6179.08		. 4
37	.1	6154.45					64	4.6	6180.19	$7_{-1} - 8_{-2}$	6180.07		.3
38	3.2	$6155.\ 60$	$10_{-7} - 11_{-8}$	6155.52		. 2				$8_{-5} - 9_{-6}$	6180.18		. 5
			$10_{-8} - 11_{-9}$			. 6				$8_{-6} - 9_{-7}$			1.5
39	1.3	6156.67					65	1.2	6180.71	$4_0 - 5_3$	6180.79		. 6
40	.1	6157.63					66	. 6	6181.49				
41	. 6	6158.79	$7_5 - 8_4$	6158.72		.1	67	1.3	6182.19	$6_3 - 7_2$	6182.13		.3
			$5_1 - 6_4$	6158.82		.1	68	.1	6182.83	$7_{-3}$ $- 8_{-3}$		6182.60	. 6
42	. 6	6159.55								$7_{-2} - 8_{-4}$			.2
43	2.5	6160.71	$9_{-4} - 10_{-5}$	6160.71		.8	69	4.9	6184.17	41-53		6183.66	. 6
			$9_{-5} - 10_{-6}$			. 3				$7_{-2} - 8_{-3}$	6183.97		1.4
			$10_{-9} - 11_{-11}$		6160.93	. 5	-	0.7		7-3-8-4			.5
			$10_{-10} - 11_{-10}$			. 2	70	3.7	6184.95	$0_0 - 7_{-1}$		6185.02	1.4
44	3.7	6162.05	$10_{-10} - 11_{-11}$	6162.20		1.1	71	1.3	6186.21	8-7-9-9		0180.37	1.2
			10-9-11-10			.4				8-8-9-8		6106 00	.4
45	. 2	6163.05		0100 41		0	70	7.9	6107 40	43-05 8	6197 49	0180.82	9.5
46	.3	6163.58	71-80	0103.41		.2	12	1.0	0107.40	8-8-9-9 8-7-0	0107.40		2.0
47	3.0	6164, 63	8-2-8-3	6164.41		.0	72	1	6180.00	0-7 0-8			.0
40	1.0	0105 00	7. 8.	0104.40		. 2	74	1.0	6180.09	5 6.	6190.05		6
48	1.8	0105.23	74-83	6165 62		.4	75	1.5	6100 86	7 -8	0130.05	6190 81	1.0
49	1.0	6166 80	0 = -10 =	0100.03	6166 70	1.0	10	1.0	0100.00	7-4-8-6		0100.01	6
50	.1	0100. 80	9 - 7 - 10 - 7		0100.70	1.0	76	1.0	6191 37	9_0-9_7	6191.24		2
51	3 5	6168 13	5=-6=		6167 82	.0	10	110	0101101	10 - 8 - 10 - 5	6191, 11		.1
51	0.0	0100.10	9_6-10_7	6168.01	0101.02	1.0	77	5.7	6192.09	7-4-8-5	6192.00		2.1
			$9_{-7} - 10_{-8}$	0100101		.3				7-5-8-6			.7
52	2.9	6168.73					78	.1	6193.37				
53	. 2	6168.87					79	2.2	6194.23	6-1-7-3		6193.95	.8
54	.1	6170.87								$6_{-2} - 7_{-1}$			.3
55	1.0	6171.79					80	2.9	6194.69				
56	3.4	6172.63	8-4-9-5	6172.59		.9	81	.5	6195.26	$6_{-2}$ -7-3	6195.00		1.9
			8-3-9-4			.3				$6_{-1}$ 7 $_{-2}$			.6
57	3.2	6173.83	$4_2 - 5_5$	6173.75		.2	82	2.6	6197.53	$5_4 - 6_3$	6197.54		.9
			9-9-10-9		6173.81	.8	83	2.2	6198.76	7-7-8-7		6198.55	1.7
			9-8-10-10			. 3				7-8-8-7			. 6
			64-73	6174.00		. 9	84	8.3	6199.62	$7_{-6} - 8_{-7}$	6199.63	13	3.4
58	6.7	6174.94	$6_2 - 7_1$	6174.64		1.2				7-7-8-8			1.1
			9-9-10-10	6174.90		. 6				$5_2 - 6_1$	6199.83	,	1.7
			9-8-10-9			1.7	85	. 3	6201.15				

TABLE 1.	Comparison of observe	d and calculated spectra	in 1.6-µ region—Continued
	1	· · · · · · · · · · · · · · · · · · ·	, ,

				1			11	1			1		
		Observed		Calculated v	vave number				Observed		Calculated v	wave number	-
Line No.	Intensity	wave	Transition		1	Inten-	Line	Intensity	wave	Transition		ī	Inten-
1101	~	number		(111)	(012)	bitty	110.		number		(111)	(012)	SIUY
	2	$cm^{-1}$		cm-1					$cm^{-1}$		cm-1		
86	1.5	6202.45	$6_{-3} - 7_{-5}$		6202.42	1.2	124	10.0	6244.71	$3_{-2} - 4_{-3}$	6244.83		6.1
			6-4-7-4			.4				$3_{-3} - 4_{-4}$			2.0
			$5_2 - 6_0$		6202.49	.7				$2_0 - 3_{-1}$	6244.94		4.2
87	6.8	6203.57	$6_{-4} - 7_{-5}$	6202.42		2.8	125	. 7					
1			$6_{-3} - 7_{-4}$			. 9	126	. 6	6246.48				
88	1.0	6204.17	$5_{-1}-6_1$		6203, 92		127	2.5	6247.97	$2_1 - 3_0$	6247.96		1.0
89	2.2	6205.01	F 0	0000.00			100	1.0	0040 10		6248.34		. 9
90	5.9	6206.29	$3_0 - 0_{-1}$	6205, 28	6906 51	2.4	128	1.0	6249.13	4. 4.		6950 09	0
01	5.6	6207 30	32 - 44 $4_2 - 5_1$	6207 32	0200. 31	28	129	1. (	0230.12	$4_3 - 4_{-1}$		6250.02	.8
92	3.8	6208.70	$4_2 - 5_2$	6208.80		2.0	130	2.8	6250 90	0-1-01		0250.20	
93	2.5	6210.11	$3_3 - 4_3$		6209.72	1.6	131	3.7	6251.35	$6_{-2} - 6_1$	6251.29		.9
1.1.1.2			$6_{-5} - 7_{-7}$		6210.43	2.1	132	4.4	6251.98	$5_{-3} - 5_{-1}$		6251.69	.8
			$6_{-6} - 7_{-6}$			.7				$4_{-4} - 4_{-1}$	6252.01		1.1
94	7.9	6211.46	$6_{-6} - 7_{-7}$	6211.46		4.4	133	3.1	6252.33	$1_1 - 2_1$		6252.41	1.7
			$6_{-5} - 7_{-6}$			1.5	134	3.3	6252.95	$5_{-2} - 5_{-1}$	6252.85		1.2
95	2.5	6213.67	$5_{-2} - 6_{-4}$		6213.49		135	4.6	6254.32	2 - 2 - 3 - 2		6254.03	.8
			$5_{-3} - 6_{-3}$		6213.50					$7_{-1} - 7_1$		6254.35	. 5
			$4_3 - 5_2$	6213.89		. 5				$2_{-1} - 3_{-3}$		6254.40	2.5
96	9.0	6214.69	$5_{-2}-6_{-3}$	6214.63		3.5	136	8.5	6255.17	$2_{-2} - 3_{-3}$	6254.97		5.4
			5-3-6-4			1.2				$2_{-1} - 3_{-2}$			1.8
07	1	6015 50	40-0-1	6214.82		3.0	107	0.1	0050 40	$0_0 - 0_3$	6255.29		.5
97	.1	6215.79	8 8 .	6916 59		2	137	2.1	6255.40	$7_0 - 7_1$	6250.60		.0
30	2.1	0210.90	8-6-0-3	6216.53		.0	138	4.9	6258 27	$1_1 - 2_0$	0207.04		1.2
			$7_{-7} - 7_{-4}$	6216.64		.1	140	1.0	6259 80				
			7-6-7-5	6216.66		.4	141	1.4	6260.73	4-1-41		6260.74	1.1
99	3.2	6217.93	$4_{1}-5_{0}$	6217.80		.9	142	4.0	6261.66	1-1 11		0200.11	1. 1
			$4_1 - 5_{-1}$		6217.68	.8	143	1.7	6262.70	$4_0 - 4_3$	6262.36		. 5
100	. 7	6219.11								$3_{-3} - 3_{-1}$		6262.97	. 9
101	. 5	6220.37					144	1.7	6263.60	$1_{-1} - 2_{-1}$		6263.50	1.8
102	. 3	6221.06	$3_1 - 4_1$		6220.86					$7_3 - 7_5$		6263.54	. 5
103	2.5	6221.89	$5_{-5} - 6_5$		6221.90	2.5	145	4.1	6264.01	$3_{-2} - 3_{-1}$	6264.01		1.4
			$5_{-4} - 6_{-6}$			. 9	146	3.9	6264.65	$1_{-1}$ - $2_{-2}$	6264.75		1.4
104	8.7	6222.98	$2_0 - 3_3$	6222.46		.7	147	8.5	6265.85	$4_1 - 4_3$		6265.23	. 9
			5-4-6-5	6222.94		5.3				$1_0 - 2_{-2}$		6265, 25	. 5
105		0004.00	5-5-6-6			1.8				$4_{-1} - 4_0$	6265.51		. 6
105	2.8	0224.03	33-42 4 - 5 -	6224.10	6994 11	.8				5-1-51	2005 01	6265.60	.0
106	2.9	6994 46	4-1-5-2	6994 44	0224.11	.0				$1_0 - 2_{-1}$ $7_1 - 7_2$	0205.91	6966 04	0. 1 7
107	8.1	6225.49	4-1-5-2	6225 22		4.0	148	4.5	6267 68	71 - 73 52 - 55		6267 20	. 1
201	011	0	4-1-5-2	0220.22		1.0	149	1.4	6268.71	$2_{-1} - 2_{9}$	6268 42	0201.20	4
			$2_1 - 3_3$		6225, 40	1.3	150	. 9	6269, 48	2-1 22	0200.12		
			$3_1 - 4_0$	6225.67		1.2	151	1.5	6270.57	$2_{-1} - 2_1$		6270.62	. 9
108	1.5	6227.97					152	2.7	6271.38	$2_{-2} - 2_1$	6271.29		1.2
109	4.3	6229.07	$7_{-5}$ 7_2	6228.79			153	4.0	6272.94	$0_0 - 1_0$		6273.23	. 5
			$6_{-6} - 6_{-3}$	6228.82		.6	154	1.3	6273.73	$7_2 - 7_3$	6270.40		.8
			$6_{-5} - 6_{-4}$	6228.85		. 2				$5_1 - 5_3$		6237.57	1.6
			$7_{-4}$ $-7_{-3}$	6228.81		. 5	155	4.0	6274.96	$4_0 - 4_2$		6274.55	. 7
110	. 2	6229.80	12.11.12	1						$3_{-1} - 3_1$		6275.00	2.1
111	5.0	6230.40	$3_2 - 4_1$	6230.62		2.3	156	5.7	6275.47	$0_0 - 1_{-1}$	6275.54		2.8
112	1.3	6231.08	$2_2 - 3_2$		6231.13	.6		0.0		$2_{-1}-2_0$	6275.71		0.6
113	3.1	6232.85	4-3-0-5		6232.86	2.8	157	0.3	6277.39	$31 - 3_3$	0077 50	6277.36	2.0
114	0.0	6922 02	4-4-0-4		6922 60	.9	150	0.9	6970 59	$3_0 - 3_1$	6277.52		2.8
114	0. 0	0235. 95	31-4-2		6233 83	.0	150	2.0	6280 70	$4_1 - 4_2$ $2_2 - 2_2$	02/9.71	6380 56	1.0
1.02			4_4-5_r	6234 06	0205.00	6.0	160	4 7	6281 81	$2_0 - 2_2$ $5_0 - 5_0$	6281 00	0280.00	.0
			4-3-5-4	0201.00		2.0	161	2.7	6282.43	1-1-11	0201,00	6282 43	2.0
115	3.5	6234, 81	3-1-4-2	6235, 15		1.4	162	1.7	6283.72	1-1 11		0202.10	2.0
116	7.0	6236.35	$3_0 - 4_{-1}$	6236.35		4.0	163	7.1	6284.89	$10 - 1_1$	6284.88		4.1
117	. 4	6238.18								74-75	6285.15		1.6
118	3.8	6239.73	$6_{-3} - 6_{-1}$		6239.55	. 6	164	5.0	6285.93	$2_1 - 2_2$	6285.98		2.1
111			$5_{-5} - 5_{-3}$		6239.63	. 6	165	8.4	6287.17	$3_2 - 3_3$	6287.12		7.5
119	6.8	6240.28	$2_2 - 3_1$	6240.32		3.3	166	6.1	6288.15	$4_3 - 4_4$	6288.10		2.4
120	5.8	6240.66	$5_{-5}-5_{-2}$	$6240.\ 62$		. 3	167	8.2	6288.80	$5_4 - 5_5$	6288.71		5.7
S. C. F.	1. 1. 1.	2 2 2 2 2 2	$5_{-4} - 5_{-3}$	6240.67		. 8	168	9.8	6289.62	$6_5 - 6_6$	6289.11		1.3
	1.113.1		$6_{-4} - 6_{-1}$	6240. 81		.8		1.1		$7_6 - 7_7$	6289.07		2.2
121	2.1	6241.57	7-2-7-1	6241.29		. 6				$5_5 - 5_4$	6289.67		1.9
122	.7	6242.67	$2_0 - 3_0$		6242. 57	.5		S		66 - 65	6289.82		3.9
123	3.0	6243.76	3-3-4-3		6243.79	2.8							
	1		3-2-4-4		0243.84	1.0	1 Sec. 1	1	8 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1			1	1

	Colculated wave number							Calculated wave number					
Line	Intensity	Observed wave	Transition			Inten-	Line	Intensity	Observed wave	Transition			Inten-
N0.		number		(111)	(012)	sity	10.		number		(111)	(012)	sity
		cm-1		cm-1					$cm^{-1}$		<i>cm</i> <sup>-1</sup>		
169	8.6	6290, 44	$4_4 - 4_3$	6290. 36		7.3	211	0.2	6334.50		0.00	λ	
			88-87	6290.49		1.2	212	5.7	6335.65	5 - 5 - 4 - 3		6335.71	3.7
170	5.3	6291.33	$3_3 - 3_2$	6291.22		2.5				$5_{-4} - 4_{-4}$			1.2
171	5.4	6291.80	$6_4 - 6_3$	6292.09		2.4	213	10.0	6336.75	$8_{-2}$ $- 8_{-3}$	6336.59		. 5
172	8.4	6292.48	$2_2 - 2_2$	6292.37		6.4				5-4-4-3	6336, 75		7.8
173	4.0	6293, 64	$1_1 - 1_0$	6293.49	6202 01	1.4				5-5-4-4		6336 05	2.6
174	3.8	6204 31	$1_1 - 1_{-1}$ $5_2 - 5_2$	6204 44	0295. 91	2.1	214	4.3	6337 78	$7_{-2} - 7_{-5}$	6337.47	0550. 55	. 5
174	.7	6295.50	$2_{2}-2_{0}$	6295. 53		.8	211	1.0	0001.10	$3_1 - 2_2$	6337.57		1.4
176	5.8	6297.14	$4_2 - 4_1$	6297.07		3.3				$4_{-2}$ -30		6338.05	. 7
177	1.0	6298.45	$3_3 - 3_1$		6298.37	2.0				$6_{-4} - 6_{-5}$	6338.22		1.1
178	. 7	6298.79					215	1.0	6338.85	$7_{-3}$ $-7_{-4}$	6338.59		
179	3.2	6299.69	$4_2 - 4_0$		6299.59	.8	216	3, 3	6339.37	$4_{-1}$ -30	6339.41		1.7
			$3_1 - 3_0$	6299.60		1.0		-		$4_{-1}$ - $3_{-1}$		6339.71	2.1
100	_	4000.05	31-3-1		6299.84	2.2	217	7.6	6340.65	$4_{-2}$ - $3_{-1}$	6340.48		5.4
180	•7	6300.85	10-00		6300.26	.0	218	.2	0341.03 6241.07		2		
			$3_3 - 3_1$ $2_0 - 2_0$		6300.81	1.7	219	4.8	6342 93	6-5-5-5		6343.10	3.5
181	4.8	6301.81	20 - 2 - 2 20 - 2 - 1	6301.65	0000.00	2.1		1.0	0012.00	$6_{-6} - 5_{-4}$			1.2
182	4.1	6303.17	$1_1 - 0_0$	6302.75		. 9	221	9.3	6344.12	6 - 6 - 5 - 5	6344.17		7.3
			$6_2 - 6_1$	6303.28		. 9				$6_{-5} - 5_{-4}$			2.4
183	. 2	6304.72					222	2.5	6345.57	$3_2 - 2_2$		6345.22	. 8
184	1.5	6306.42	$5_1 - 5_{-1}$		6306.58	1.4				8-4-8-5	6345.15		. 5
185	1.9	6306.94	$5_1 - 5_0$		6306.91	. 6	200		0010.00	7-4-7-7	6345.71		. 5
186	.6	6307.65	0.9	6200 62		1.1	223	3.9	6346.86	53-4-1		6247.06	2.4
187	0.4	6309.77	$3_2 - 3_{-1}$	6309.62		1.1	994	8.0	6347 98	5 - 2 - 4 - 2	6347 92	0347.00	5.5
188	7	6310 44	$4_0 - 4_{-1}$ $4_2 - 4_1$	0505.77	6310 12	1.0	224	0.0	0517.30	$5_{-2} - 4_{-2}$	6348.15		1.8
100		0010.11	$2_{-2}-1_0$		6310. 38	. 6	225	4.5	6350.01	7-7-6-5		6350.09	3.1
			$3_{-1} - 3_{-3}$	2	6310.82	1.1				7-6-6-6			1.0
189	4.3	6311.73	$2_{-1} - 1_0$	6311.67		1.4				$3_3 - 2_1$	6350.42		1.9
	~		$3_{-1} - 3_{-2}$	6311.83		. 5	226	9.4	6351.32	$4_6 - 3_1$	6351.04		5.1
190	8.0	6312.83	$2_{-1} - 1_{-1}$		$6312 \ 10$	2.1				$7_{-6}$ - $6_{-5}$	6351, 19		6.3
			$4_1 - 4_{-1}$		6312.74	1.3	007		0250 52	7-7-6-6			2, 1
101	4.7	6919 05	2-2-1-1	6312.77		4.7	227	.2	6352.53	 99 -			. 3
191	4.7	6314 07	$5_0 - 5_{-3}$ $7_1 - 7_{-1}$	0313, 34	6313 94	1.5	220		0555.04	8-6-8-7	6353_03		. 1
152	2.0	0014.07	41-4.2	6314, 12	0010.01	.5	229	5, 6	6353.79	$4_2 - 3_3$	6353, 57		3.4
193	1.2	6314, 43	1 - 2	0011112			230	5.7	6354.21	$4_1 - 3_1$		6353.91	1.7
194	4.0	6315.73	$6_3 - 6_1$		6315.57	. 7	231	4.7	6354.48	$6_{-4} - 5_{-2}$		6354.64	.8
			$5_2 - 5_{-1}$	6316.01		1.1				$6_{-3} - 5_{-3}$		6354.68	. 23
195	3.1	6317.62	$6_0 - 6_{-1}$	6318.29		1.1	232	8.8	6355.69	$5_{-1} - 4_1$		6355, 35	1.2
196	4.8	6319.66	$5_{-1} - 5_{-3}$		6319.07	1.0				$6_{-4} - 5_{-3}$		6355.94	5.1
			3-3-2-1		6319,68	2.9	022	5.0	6256 45	$b_{-3} - b_{-2}$		6356 71	1.7
			$0_1 - 0_{-1}$		6319.98	1.0	200	5.0	0330, 43	8-7-7-7		0550, 71	2.4
197	10.0	6320.75	$4_{-1} - 4_{-3}$		6320.52	.9	234	9.3	6357.79	$5_0 - 4_1$	6357.65		3.9
			$20 - 1_1$	6320.57		4.1				8-8-7-7	6357.91		5.1
			$3_{-2} - 2_{-1}$	6320.71		6.4				$8_{-7}$ ${7-6}$			1.7
			$3_{-3} - 2_{-2}$	6320.95		2.1	235	2.0					
198	6.7	6321.47	4-2-4-3	6321, 30		1.3	236	1.0	6359.73				
			$5_0 - 5_{-3}$	6321.39		1.4	237	2.7	6360.63	$5_{-1} - 4_0$	6360.29		1.5
199	.1	6322.70	$7_2 - 7_1$	6322.28	4909.00	.6	000		0901 00	$5_0 - 4_0$		6360.18	.5
200	4.7	6324.00	$2_1 - 1_1$		6323.99	2.0	208	2.9	0301.88	7-5-0-3 7-6-4		6361.85	2.0
201	2.2	6325 30					239	7.9	6363 07	9-0-8-7		6362.96	1.9
202	4 2	6327 23	$3_{-1} - 2_1$		6327_04	1.2	200	1.0	0000.01	9_s-8_s		0002.00	. 6
204	6.0	6327.98	4-4-3-2		6327.97	1.2				$7_{-4} - 6_{-3}$	6 <b>3</b> 63.07		4.2
			4-3-3-3		6328.01	3.5				$7_{-5} - 6_{-4}$			1.4
205	1.00	6329, 01	$5_{-3} - 5_{-5}$		6328.79	.8	240	7.2	6364.14	$9_{-8} - 8_{-7}$	6364.08		3.8
			$7_0 - 7_{-3}$	6328.92		.7				$9_{-9} - 8_{-8}$			1.3
			$4_{-4}$ $3_{-3}$	6329.01		7.5	241	2.1	6364.52				
		40.00	$4_{-3}-3_{-2}$			2.5	242	1.0	6365.31	$6_{-2} - 5_0$		6365.20	.5
206	7.0	6329, 51	6-2-6-3	6329.25		1.3	243	5.0	6366.60	6-1-5-1	6260 60	6366, 44	1.4
207	1.0	6331 46	$3_0 - 2_1$	0329.56		3.8	244	5.1	6366 80	0 - 1 - 30	0908.09		1. 2
207	1. 9	6332.00	3-1-20	6332.23		1.7	245	7.2	6367.44	$6_{-2} - 5_{-1}$	6367.48		3.3
209	3.7	6332. 85	30-20		6332.68	.6							
210	2.7	6333, 89	$2_2 - 1_{-1}$	6333. 85		. 9							

TABLE 1.	Comparison of	f observed an	d calculated	spectra in	1.6-µ region	-Continued
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Line		Observed		Calculated v	vave number	Inton	Tine		Observed		Calculated v	Calculated wave number	
No.	Intensity	wave number	Transition	(111)	(012)	sity	No.	Intensity	wave number	Transition	(111)	(012)	Inten- sity
246	6.0	cm <sup>-1</sup>	87 .	<i>cm</i> <sup>-1</sup>	6369 75	0.5	964	1.0	cm-1		cm-1		
210	0.0	0000.00	8-5-7-5		0508.75	1.5	204	4.0	0587.15	5. 4.		6997 70	0.0
			$5_1 - 4_2$	6368, 83		1.0	265	2.5	6388_03	04-44		0387.79	0.9
			10-9-9-9	6368, 85		1. 1	266	6.7	6388 44	9-4-8-2	6388 79		1.6
			10 - 8 - 9 - 8	0000,00			200	0.1	0000.11	$9_{-5} - 8_{-4}$	0388.15		1.0
247	9.5	6370.00	$10_{-9} - 9_{-8}$	6370.00		2.7				$11_{-8} - 10_{-7}$	6388 79		11
			$10_{-10} - 9_{-9}$			.9				$11_{-9} - 10_{-8}$	0000.10		4
			8 - 6 - 7 - 5	6370.02		3.3	267	4.7	6389.55	55 - 43		6389.52	2.2
			$8_{-5} - 7_{-4}$			1.1	268	1.3	6390.47	$5_3 - 4_1$		6390, 49	. 9
248	1.5	6371.06					269	2.0	6391.65				
249	. 2	6372.34					270	2.0	6392.46	$7_2 - 6_3$	6392.41		1.5
250	2.0	6373.39	7 - 3 - 6 - 1		6373.24	1.2	271	5.0	6393.33				
			7 - 2 - 6 - 2			.4	272	3.3	6393.98	$8_{-2} - 7_{-1}$	6393.82		1.7
251	3.0	6373.96					<b>2</b> 73	5.0	6394.97	$10_{-6} - 9_{-5}$	6395.07		1.0
252	6.7	6374.54	$11_{-11} - 10_{-9}$		6374.33	. 9			-	$10_{-5} - 9_{-4}$			. 3
			$11_{-10} - 10_{-10}$			. 3	274	4.4	6397.55	$6_3 - 5_3$		6397.73	1.1
			7-2-6-1	6374.61		3.0	275	.1	6398.78	$7_1 - 6_2$	6398.59		.7
			7-3-6-2			1.0	276	1, 3	6399.06				
253	6.8	6375.49	$9_{-7} - 8_{-5}$		6375.31	.7	277	3.9	6399.88	$9_{-2}$ -8_{-1}	6400.36		1.1
			$9_{-6} - 8_{-6}$			. 2	278	2.7	6401.20				
			$6_1 - 5_2$	6375.39		. 9	279	.7	6401.80				
			$11_{-10} - 10_{-9}$	6375.52		1.8	280	2.2	6403.45				
014			$11_{-11} - 10_{-10}$			. 6	281	1.3	6403.69				
254	4.2	6376.17	0 0	0070.00			282	1.2	6405.35				
200	0.7	0370.00	$9_{-6} - 8_{-5}$	6376.62		2.4	283	2.0	6406.07				
955 A	2 7	6970 06	$9_{-7} - 8_{-6}$			.8	284	3.5	6406.67	$8_0 - 7_1$	6406.95		1.4
200A. 956	0.1 9.6	0378.80	4 9	6970 09			285	1.3	6407.67				
200	5.0	0379.07	44-51	0379.03	2270 42	.7	286	3.0	6408.34				
			12-11-11-11 1211		0379.40	0. 2	287	4.1	6408.77	$6_5 - 5_5$		6408.61	1.6
957	5.0	6280 51	12-12-11-10 6- 5	6970 07		.0	288	1.0	6409.49	$0_6 - 0_4$		6409.50	. 5
201	0.0	0380. 31	$6_3 - 5_4$	6380.22		.0	289	2.0	6410.20				
			12 - 11 - 11	6380.60		0.0	290	2.0	5410.00 6411 52				
			12 - 11 - 11 - 10 12 - 10 - 11 - 10	0580, 05		.4 1.9	291	2.2	6412.08				` )
			$\frac{12-12}{8-2-7-2}$		6380 54	1.2	292	2.0	6412.08				
			8-2-7-4		0500.04	.0	293	1.0	6414 21				
258	8.0	6381.72	$6_1 - 5_1$		6381_64	9	204	2.0	6414 48				1.9
			8-4-7-3	6381.96	0001.01	2.3	296	2.0	6415 89	87.	6415 76		1.2
			8-3-7-2			8	297	1.3	6417.14	02 13	0110.70		1.0
259	5.6	6382.85	10-8-9-7	6382.87		1.6	298	2.1	6418 26				
			10-7-9-6			. 5	299	2.7	6419.15				
			$7-1-6_1$		6382.99	. 8	300	. 3	6421.94				
260							301	.7	6423.05				
261			$6_4 - 5_5$	6383.30		1.1	302	1.5	6424.33				
262	4.3	6384.43					303	2.7	6425.52				
263	6.0	6385.16	$7_0 - 6_1$	6385.23		2.2	304	. 4	6427.40				
			$6_2 - 5_0$		6385.47	2.9	305	2.0	6428.99				
							306	2.4	6429.50	77-65		6249.57	1.0

TABLE 1. Comparison of observed and calculated spectra in 1.6-µ region-Continued

#### 3. Analysis of Experimental Results

The previous analysis of the low-resolution data showed that the absorption in this region could be accounted for by an A-type band, i. e., the dipole moment change along the least inertial axis. The previously calculated spectrum was compared to the new data, and the strong transitions were brought into agreement with the observed absorption peaks through the use of derivatives of the energy with respect to the inertial parameters [4]. Then all transitions of intensity sufficient to be observed were calculated. The fit between the observed and calculated spectra was very good, except that quite a few of the weaker peaks were not accounted for. Through the use of known ground-state energies [4] it was possible to determine that many of these unaccounted for peaks resulted from a B-type band centered about 1 cm<sup>-1</sup> lower than the main band. The excited state inertial parameters for this band were determined as described previously [1]. A classical centrifugal distortion correction was applied to the excited-state energy levels of both bands [4].

The excited vibrational state of the A-type band is undoubtedly the (1,1,1) [1]. The excited vibrational state involved in the B-type band is probably the lower component of the resonating pair (0,1,2)and (2,1,0). A calculation based on the quadratic expression for the vibrational-band system of hydrogen sulfide [5] predicts this band to be at  $6,285 \text{ cm}^{-1}$ , compared to the observed  $6288.30 \text{ cm}^{-1}$ .

The parameters giving the best fit of the observed data are (1,1,1) A=10.398, B=8.935, C=4.548,  $\nu_0=6289.26$  cm<sup>-1</sup>; (0,1,2) A=10.394, B=8.918,  $\mathring{C}=4.547, \nu_0=6288.28 \text{ cm}^{-1}.$ 

Many of the listed peak positions cannot be attributed to single transitions. The overlapping of two or more transitions will give an observed peak position that will not be the true frequency of any one transition. In many instances, however, definite assignments can be made. This is generally true for the extremely strong transitions because the observed frequency will not be greatly affected by weaker satellites. In the (1,1,1) band the mean deviation of 76 such assignments is 0.10 cm<sup>-1</sup>. In the (0,1,2) band some 36 peaks were assigned that do not have a dual assignment in the (1,1,1) band. The mean deviation in this case was  $0.09 \text{ cm}^{-1}$ . That so many of the individual deviations are greater than the mean must be attributed to the effect of finite slit on overlapping transitions. This is especially true of the (0,1,2) band, which is very weak compared to the main band in this region. The classical centrifugal-distortion correction may also be responsible for some of the large deviations. However, it should be pointed out that for low J the corrections are small in any event, and for high J the correction should become classical. Another advantage of this method is that this correction can be evaluated level for level and applied to the rigid levels calculated from published energy tables. The agreement between the observed and calculated spectra may be seen in figure 1 and in table 1. The relative intensity scale within a given band is the same as that used in the previous analysis. In reporting the intensities, the line strengths in the (1,0,2) band were divided by two to give a better picture of the relative importance of the two bands. The P branch of the (1,1,1) band is roughly three times as strong as the P branch of the (0,1,2) band, and the R branch of the (1,1,1) band is approximately twice as strong as the R branch of the (0,1,2) band. Some transitions of the two bands appear to be different from the intensities that one would predict from simple theory. These intensity discrepancies no doubt arise from the effect of Coriolis interaction between the two bands.

The new values of the inertial parameters of the (1, 1, 1) band agree with those determined previously within the estimated error. As the (0, 1, 2) band is somewhat weaker, and the transitions somewhat harder to identify, it is felt that the inertial parameters for this band are not as well determined as those for the main band.

#### 4. Discussion

In figure 1 it will be noted that the calculated spectrum does not account for all the observed absorption lines at the high and low limits of the band. The calculated spectrum in these regions is incomplete as calculations were not extended above J=12, the upper limit of published energy tables [6]. Likewise, the other half of the resonating pair,

(0, 1, 2) and (2, 1, 0), is predicted to fall in the highfrequency region of this absorption. It has not been possible to locate any transitions that can be definitely assigned to a band having this upper state; however, it would in all probability be lost in the other two overlapped bands.

In the infrared spectra of water-type molecules, it has been observed that A-type bands are normally very much stronger than B-type bands. In the hydrogen-sulfide spectrum only two B-type bands, (0, 1, 0) and (0, 2, 0), have been observed. It may be possible that other B-type bands are present in the spectrum of hydrogen sulfide that are completely overlapped by strong A-type bands [5, 7, 8, 9]. In each case observed so far the overlapping bands have fitted the criterion for Coriolis-type interaction [10]. In the analysis of the rotational fine structure of all these bands, it has been possible to account for the observed absorption by calculating a theoretical spectrum with effective moments of inertia. There has been no drastic perturbation of individual energy levels. This is in line with what might be expected. As the inertial parameters of the two excited states are nearly the same, the initial separation of the unperturbed energy levels is very nearly constant. Both the interaction and the energy vary nearly as  $K^2$ ; hence it is possible to fit the observed spectra through the use of effective moments of inertia [5], the effect of the Coriolis interaction being absorbed in the inertial constants. The extent of this effect cannot be ascertained until the vibrational system of hydrogen sulfide is better understood. The fact that these B-type bands are observed only when overlapped by a normally strong A-type band seems to indicate that the intensity is interchanged quite readily.

As has been observed in other B-type bands [11], R branch transitions between high  $\tau$  levels of adjacent J values are apparently stronger than would be predicted by the simple rigid theory. There is still no adequate explanation for this.

The value of  $\Delta = I_c - I_a - I_b$  lends additional support to the vibrational assignments proposed here. In both bands its value is large indicating that  $n_2$ equals at least 1 in each of the excited states. The exact value of  $\Delta$  is of little significance because of the presence of the Coriolis interaction.

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