

An Expansion Method for Parabolic Partial Differential Equations¹

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The aim of this paper is to adapt to certain parabolic partial differential equations an expansion method of solution developed by S. Faedo³ for hyperbolic equations. In order to make possible a moderately compact presentation, the equations treated are not the most general to which the method is applicable, but are the simplest nontrivial relatives of the heat equation. Similarly, the boundary values and initial conditions are not the most general, but are assumed to be in a canonical form to which others, if sufficiently smooth, can be reduced. The method of solution not only shows the existence of a solution, but describes a definite procedure for approximating it. Some remarks are made on the possibility of estimating the error.

1. Statement of Problem

In the region S , $0 \leq x \leq \pi$, $0 \leq t < \infty$, we seek a solution of the equation

$$L(u) \equiv u_{xx} - u_t - gu = f \quad (1)$$

subject to the initial conditions

$$u(x, 0) = u(0, t) = u(\pi, t) = 0. \quad (2)$$

The functions f and g are continuous in S , together with their first two partial derivatives. As a compatibility condition, we require that $f(0, 0) = f(\pi, 0) = 0$. It is convenient to assume that $g \geq 0$. This is not at all essential, since a substitution of the form $u = e^{\lambda t} v$ transforms (1) into

$$v_{xx} - v_t - (g + \lambda)v = fe^{-\lambda t}. \quad (1')$$

The new coefficient $g + \lambda$ can be made positive in as large a part of S as desired, and this is sufficient for our purposes.

This method can be applied to more general equations in which u_t has a variable positive coefficient and the u_x term is present. However, the equation selected is of sufficient generality to illustrate the method and results and yet not make the computations unpleasantly complicated. In case $g \equiv \text{constant}$, the method reduces to the method of expanding in eigenfunctions.

2. The Moment Method⁴

We seek an approximation to u in the form^{5,6}

$$u_n(x, t) = \sum_k C_{n,k}(t) \sin kx, \quad (3)$$

subject to

$$\int \{L(u_n) - f\} \sin jx dx = 0, \quad j=1, 2, \dots, n, \quad (4)$$

where

$$C_{n,k}(0) = 0. \quad (5)$$

Equation (4) amounts to

$$-j^2 C_{n,j}(t) - C'_{n,j}(t) - \sum_k C_{n,k}(t) b_{k,j}(t) - a_j(t) = 0, \quad (6)$$

where

$$b_{k,j}(t) = \frac{2}{\pi} \int g(x, t) \sin jx \sin kx dx,$$

$$a_j(t) = \frac{2}{\pi} \int f(x, t) \sin jx dx.$$

The functions $C_{n,j}(t)$ are thus determined by a system of n simultaneous first-order linear differential equations. These equations are already in canonical form, that is, the derivatives are solved for in terms of the other quantities, and so they have a unique solution satisfying (5). Thus the functions $u_n(x, t)$ are determined in unique fashion.

We intend to show that the sequence $u_n(x, t)$ converges to a solution of (1). The method proceeds by showing that the u_n and certain of their derivatives are uniformly equicontinuous. Recall that the functions f of a family are uniformly equicontinuous in a region, provided that ϵ is given greater than 0, there is a $\delta > 0$ such that $|f(p) - f(q)| < \epsilon$ for any f of the family and any two points p and q of the region whose distance does not exceed δ . Now, Ascoli's theorem asserts that out of a uniformly bounded and equicontinuous sequence of functions a uniformly continuous subsequence can be selected. Therefore, by selecting subsequences of subsequences a sufficient number of times, we shall arrive at a subsequence of the u_n which converges uniformly, together with sufficiently many of its derivatives to allow us to make the existence proof. This latter is effected by using a Green's function and showing that the u_n , for $n \rightarrow \infty$, come

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³ S. Faedo, Un nuovo metodo per l'analisi esistenziale e quantitativa dei problemi di propagazioni. Annali della Scuola Normale Superiore di Pisa, [III], 1, Fasc. I-IV (1947).

⁴ Or Galerkin's method. See, I. S. Sokolnikoff, Mathematical theory of elasticity (McGraw-Hill Book Co., Inc., New York, N. Y., 1946).

⁵ As on p. 37 of the reference given in footnote 3, after the existence of the solution has been proved, one can show that the system $\{\sin kx\}$ can be replaced by other suitable orthonormal systems.

⁶ Σ means \sum_1^n and \int means \int_0^π .

successively closer to satisfying an integral equation corresponding to (1), and thus that the limit function does satisfy such an equation.

To prove equicontinuity of our sequences, we make use of a criterion of Tonelli⁷ as follows: A class of continuous functions $v(x,t)$ in the region $S_T: 0 \leq x \leq \pi, 0 \leq t \leq T$, is uniformly equicontinuous in this region if (a) for each \bar{x} and t , $v(\bar{x},t)$ and $v(x,t)$ are absolutely continuous functions of t and x , respectively, (b) the partial derivatives v_x and v_t are summable in S_T , (c) there exists a constant c such that $\int v_x^2 dx < c$ and $\int v_t^2 dx < c$ for all members of the family. Furthermore any accumulation function of a class satisfying (a), (b), and (c) also satisfies (a) and (b).

Now the functions u_n and those partial derivatives in which we are interested are easily seen to be continuous and to satisfy (a) and (b), and the burden of the proof of equicontinuity falls on obtaining inequalities of the form in (c). This is done by making use of devices similar to those frequently employed to show uniqueness, and to recall them, we shall first prove a uniqueness theorem.

3. Uniqueness

Suppose that two solutions of the system (1) (2) exist, all derivatives appearing in (1) being continuous in S . Their difference v will satisfy (1), (2) with f replaced by zero. Thus $L(v)=0$ and so

$$\int L(v)v dx = \int \{v_{xx} - v_t - gv\} v dx = 0. \quad (7)$$

Now

$$\int v_t v dx = \int \frac{1}{2} \frac{d}{dt} v^2 dx = \frac{1}{2} \frac{d}{dt} \int v^2 dx,$$

and

$$\int v_{xx} v dx = v_x v \Big|_0^\pi - \int v_x^2 dx = - \int v_x^2 dx.$$

If we put $K(t) = \int v^2 dx$, (7) becomes

$$\frac{1}{2} K'(t) = - \int (v_x^2 + gv^2) dx \leq 0,$$

since $g \geq 0$. Since by (2), $K(0)=0$, we see that $K(t) \leq 0$. But K is by its nature nonnegative and $K \equiv 0$, which implies that $v \equiv 0$ and that the two solutions are identical.

Theorem 0. There is at most one function, continuous together with all derivatives appearing in (1) which satisfies the system (1), (2).

4. Certain Bounds

This section is devoted to obtaining the inequalities required to carry out the program outlined in section 3.

4.1. *Lemma.* Let $K(x)$ be a nonnegative function satisfying for $x \geq 0$ the inequality $K'(x) \leq \alpha \sqrt{K(x)}$, where $\alpha \geq 0$. Then for $x > 0$, $\sqrt{K(x)} \leq \sqrt{K(0)} + (\alpha/2)x$.

Suppose that for some positive x_0 , $\sqrt{K(x_0)} > \sqrt{K(0)} + (\alpha/2)x_0$. Then there exists x_1 such that $\sqrt{K(x_1)} = \sqrt{K(0)} + (\alpha/2)x_1$, but $\sqrt{K(x)} > \sqrt{K(0)} + (\alpha/2)x$ for $x_1 < x \leq x_0$. In particular, $K(x) > 0$ in $x_1 < x < x_0$, and so in that interval

$$\frac{K'(x)}{\sqrt{K(x)}} \leq \alpha.$$

Integrating from x_1 to x_0 , we have

$$\sqrt{K(x_0)} - \sqrt{K(x_1)} \leq (\alpha/2)(x_0 - x_1),$$

or

$$\sqrt{K(x_0)} \leq \sqrt{K(x_1)} + (\alpha/2)(x_0 - x_1) = \sqrt{K(0)} + (\alpha/2)x_0,$$

a contradiction. The lemma is proved.

Corollary. Let $K(x)$ be a nonnegative function satisfying for $x \geq 0$ the inequality

$$K'(x) \leq \alpha(x) \sqrt{K(x)}, \quad (8)$$

where $\alpha(x)$ is nonnegative and nondecreasing. Then for $x \geq 0$,

$$\sqrt{K(x)} \leq \sqrt{K(0)} + \frac{1}{2} x \alpha(x). \quad (9)$$

Consider any $x_0 > 0$. For $x \leq x_0$, $K'(x) \leq \alpha(x_0) \sqrt{K(x)}$. By the preceding lemma, $\sqrt{K(x)} \leq \sqrt{K(0)} + \frac{1}{2} \alpha(x_0)x$. Setting $x = x_0$ and noting that x_0 is arbitrary, we get the stated result.

4.2. *The bound⁸ for $\int u_n^2 dx$.* Let each of the equations (4) be multiplied by $C_{n,i}(t)$ and the results summed. There comes out the equation

$$\int \{u_{xx} - u_t - gu - f\} u dx = 0. \quad (10)$$

Since

$$\int u_{xx} u dx = - \int u_x^2 dx \text{ and } \int u_t u dx = \frac{1}{2} \frac{d}{dt} \int u^2 dx,$$

(10) may be written

$$\frac{1}{2} \frac{d}{dt} \int u^2 dx = - \int (u_x^2 + gu^2) dx - \int f u dx \leq \int |f u| dx. \quad (11)$$

Set $\int u^2(x,t) dx = K(x)$. From (11) and Schwarz's inequality,

$$\frac{1}{2} K'(x) \leq \int |f u| dx \leq \sqrt{\int f^2 dx} \cdot \sqrt{\int u^2 dx}. \quad (12)$$

Set

$$\max_{\tau \leq t} \sqrt{\int f^2(x,\tau) dx} = F(t). \quad (13)$$

Then (12) gives $K'(t) \leq 2F(t) \sqrt{K(t)}$. From the corollary of 4.1 we deduce that $\sqrt{K(t)} \leq \sqrt{K(0)} + tF(t)$. Now $K(0) = \int u^2(x,0) dx = 0$, and so $\sqrt{K(t)} \leq tF(t)$.

⁷ Tonelli, L'estremo assoluto degli integrali doppi, Annali della Scuola Normale di Pisa (1933).

⁸ In section 4 we shall frequently omit the subscript n from u_n . Whenever u appears in section 4, it will mean u_n .

Theorem 1.

$$\int u_n^2(x,t)dx \leq t^2 F^2(t), \quad (14)$$

where F is given by (13).

4.3. The bound for $\int u_n^2 dx$. Let the equations (4) be differentiated with respect to t , multiplied by $C'_{n,i}(t)$, and added. There results from this

$$\int \{u_{xxt} - u_{tt} - gu_t - g_t u - f_t\} u_i dx = 0. \quad (15)$$

Again,

$$\int u_{xxt} u_i dx = - \int u_{xt}^2 dx \text{ and } \int u_{tt} u_i dt = \frac{1}{2} \frac{d}{dt} \int u_i^2 dx,$$

so putting $K(t) = \int u_i^2 dx$, we get

$$\begin{aligned} \frac{1}{2} K'(t) &= - \int (u_{xt}^2 + gu_t^2) dx - \int (g_t u + f_t) u_i dx \quad (16) \\ &\leq \int (|g_t u| + |f_t|) |u_i| dx \\ &\leq \sqrt{\int (|g_t u| + |f_t|)^2 dx} \cdot \sqrt{K(t)}. \end{aligned}$$

We define

$$\max_{\tau \leq t} |g_t(x,t)| = G_1(t), \quad (17)$$

and

$$\max_{\tau \leq t} \sqrt{\int f_i^2(x,\tau) dx} = F_1(t). \quad (18)$$

Then

$$\begin{aligned} \int (|g_t u| + |f_t|)^2 dx &\leq \int (G_1^2(t) u^2 + 2G_1(t) |u f_t| + f_t^2) dx \\ &\leq G_1^2(t) \int u^2 dx + 2G_1(t) \sqrt{\int u^2 dx \int f_t^2 dx} + \int f_t^2 dx \\ &\leq G_1^2(t) t^2 F^2(t) + 2G_1(t) t F(t) F_1(t) + F_1^2(t) = H^2(t), \end{aligned}$$

this last equation being a definition of $H(t)$.

From (16) we have $K'(t) \leq 2H(t) \sqrt{K(t)}$, and by 4.1,

$$\sqrt{K(t)} \leq \sqrt{K(0)} + tH(t). \quad (19)$$

Now

$$\begin{aligned} K(0) &= \int u_i^2(x,0) = \int (\sum C'_{n,k}(0) \sin kx)^2 dx \\ &= \frac{\pi}{2} \sum C'_{n,k}(0)^2 = \frac{2}{\pi} \sum (\int f(x,0) \sin kx dx)^2 \\ &\leq \int f^2(x,0) dx \end{aligned}$$

by Bessel's inequality. Thus $K(0) \leq F^2(0)$, and from (19) we conclude that $\sqrt{K(t)} \leq F(0) + tH(t)$.

Theorem 2.

$$\int u_n^2(x,t) dx \leq (F(0) + tH(t))^2 = A^2(t), \quad (20)$$

where F is given by (13),

$$H^2(t) = G_1^2(t) t^2 F^2(t) + 2G_1(t) t F(t) F_1(t) + F_1^2(t), \quad (21)$$

and G_1 and F_1 are given by (17) and (18).

4.4. The bound of $\int u_n^2 dx$. Equation 10 may be written

$$\int \{-u_x^2 - uu_t - gu^2 - fu\} dx = 0,$$

or

$$\begin{aligned} \int u_x^2 dx &= - \int \{gu^2 + uu_t + fu\} dx \leq \int (|uu_t| + |fu|) dx \\ &\leq \sqrt{\int u^2 dx} \cdot \sqrt{\int u_t^2 dx} + \sqrt{\int u^2 dx} \cdot \sqrt{\int f^2 dx} \\ &\leq \sqrt{t^2 F^2(t) A^2(t)} + \sqrt{t^2 F^2(t) F^2(t)} \\ &= tF(t)(A(t) + F(t)). \end{aligned}$$

Theorem 3.

$$\int u_n^2(x,t) dx \leq tF(t)(A(t) + F(t)) = B(t), \quad (22)$$

where A is given by (20).

4.5 The bound of $\int u_{n,tt}^2(x,t) dx$. We differentiate (4) twice with respect to t , multiply by $C'_{ni}(t)$ and sum, to get

$$\int \{u_{xxtt} - u_{ttt} - gu_{tt} - 2g_t u_t - g_{tt} u - f_{tt}\} u_{ii} = 0. \quad (23)$$

As before, we write

$$\int u_{xxtt} u_{ii} dx = - \int u_{xt}^2 dx,$$

$$\int u_{ttt} u_{ii} dx = \frac{1}{2} \frac{d}{dt} \int u_{ii}^2 dx,$$

and (23) becomes, on setting $K(t) = \int u_{ii}^2 dx$,

$$\begin{aligned} \frac{1}{2} K'(t) &= - \int \{u_{xt}^2 + gu_{tt}^2\} dx - \int \{2g_t u_t + g_{tt} u + f_{tt}\} u_{ii} dx \\ &\leq \int \{|2g_t u_t| + |g_{tt} u| + |f_{tt}|\} |u_{ii}| dx. \end{aligned}$$

It is clear that by the use of Schwarz's inequality we can obtain an inequality of the form

$$K'(t) \leq J(t) \sqrt{K(t)}, \quad (24)$$

where $J(t)$ is an increasing and continuous function expressible in terms of G_1 , F , F_1 , and the similar functions G_2 and F_2 corresponding to g_{tt} and f_{tt} .

We estimate $K(0)$ as follows. Differentiate (4) with respect to t and put $t=0$, to obtain

$$\int u_{ii}(x,0) \sin jx dx = \int \{u_{xxt}(x,0) - gu_t(x,0) - f_t(x,0)\} \sin jx dx. \quad (25)$$

for $j=1, 2, \dots, n$. Now $u_{ii}(x,0)$ is a linear combination of $\sin jx$, $j=1, 2, \dots, n$, and by (25) its first n Fourier coefficients are the same as those of $u_{xxt}(x,0) - gu_t(x,0) - f_t(x,0)$. Thus by Bessel's inequality,

$$\int u_{ii}^2(x,0) dx \leq \int \{u_{xxt}(x,0) - gu_t(x,0) - f_t(x,0)\}^2 dx. \quad (26)$$

If in (4) one puts $t=0$, there results

$$- \int u_i(x,0) \sin jx dx = \int f(x,0) \sin jx dx.$$

Two integrations by parts gives

$$-\int u_{xx}(x,0)\sin jx dx = \int f_{xx}(x,0)\sin jx dx,$$

and since $u_{xx}(x,0)$ is a linear combination of $\sin jx$, $j=1, 2, \dots, n$, we see as above,

$$\int u_{xx}^2(x,0)dx \leq \int f_{xx}^2(x,0)dx. \quad (27)$$

By using Schwarz's inequality, (27), and the bounds previously obtained, the right hand number of (25) can be dominated by a continuous function depending only on f and g and their derivatives. That is, $K(0)$ is bounded uniformly in n , and from (24) it follows that the same is true of $K(t)$.

Theorem 4. *There exists a continuous function $C(t)$, determined by f, g and their derivatives, such that*

$$\int u_{nxt}^2(x,t)dx \leq C(t). \quad (28)$$

4.6. *The bound for $\int u_{nxt}^2 dx$.* From eq (15) we have on replacing $\int u_{xxt}u_t dx$ by $-\int u_{xt}^2 dx$,

$$\int u_{xt}^2 dx = \int \{u_{tt} + gu_t + g_t u + f_t\} u_t dx.$$

Because of the bounds previously obtained, it is clear that we have the theorem:

Theorem 5. *There exists a continuous function $D(t)$, expressible in terms of f and g and their derivatives such that*

$$\int u_{nxt}^2(x,t)dx \leq D(t). \quad (29)$$

4.7. *The bound for $\int u_{nxx}^2 dx$.* If equations (4) be multiplied by $j^2 \sin jx$ and summed, we have

$$\begin{aligned} \int u_{xx}^2 dx &= \int \{u_t + gu + f\} u_{xx} dx \\ &\leq \sqrt{\int \{u_t + gu + f\}^2 dx} \sqrt{\int u_{xx}^2 dx}. \end{aligned}$$

Now from the previous bounds we can find a continuous function $E(t)$ which depends only on f, g and their derivatives, such that

$$\int \{u_t + gu + f\}^2 dx \leq E(t). \quad (30)$$

Thus we have the theorem:

Theorem 6. *There exists a continuous function $E(t)$ independent of n such that*

$$\int u_{nxx}^2(x,t)dx \leq E(t). \quad (31)$$

4.8. *An integral inequality.* Consider the Green's function $G(x, \xi)$ relative to the equation $y'' + \lambda y = 0 = y(0) = y(\pi)$. Specifically

$$G(x, \xi) = \begin{cases} \frac{x(\pi - \xi)}{\pi} & x \leq \xi, \\ \frac{\xi(\pi - x)}{\pi} & x \geq \xi, \end{cases}$$

and its uniformly convergent Fourier expansion is

$$G(x, \xi) = -\frac{2}{\pi} \sum_1^{\infty} \frac{\sin jx \sin j\xi}{j^2}. \quad (32)$$

For brevity, set

$$u_t + gu + f = \Phi(u). \quad (33)$$

then

$$u_{xx} - \Phi(u) = L(u) - f. \quad (34)$$

Multiply (34) by $G(x, \xi)$ and integrate with respect to x to obtain

$$-u(\xi, t) - \int \Phi(u) G(x, \xi) dx = \int \{L(u) - f\} G(x, \xi) dx. \quad (35)$$

Now

$$\begin{aligned} \int \{L(u) - f\} G(x, \xi) dx &= -\frac{2}{\pi} \sum_1^{\infty} \frac{\sin j\xi}{j^2} \int \{L(u) - f\} \sin jx dx, \\ &= -\frac{2}{\pi} \sum_{n+1}^{\infty} \frac{\sin j\xi}{j^2} \int \{L(u) - f\} \sin jx dx, \end{aligned}$$

because of (4). Now $L(u) - f = u_{xx} - \Phi(u)$, and

$$\begin{aligned} u_{xx} &= -\sum_1^n C_{n,j}(t) j^2 \sin jx; \text{ hence} \\ \int u_{xx} \sin jx dx &= 0 \end{aligned}$$

for $j > n$, giving

$$\int \{L(u) - f\} G(x, \xi) dx = -\frac{2}{\pi} \sum_{n+1}^{\infty} \frac{\sin j\xi}{j^2} \int \Phi(u) \sin jx dx.$$

By Schwarz's inequality,

$$\begin{aligned} |\int \{L(u) - f\} G(x, \xi) dx| &\leq \frac{2}{\pi} \sqrt{\sum_{n+1}^{\infty} \frac{1}{j^4}} \cdot \sqrt{\sum_{n+1}^{\infty} (\int \Phi(u) \sin jx dx)^2}. \quad (36) \end{aligned}$$

Now

$$\sum_{n+1}^{\infty} \frac{1}{j^4} \leq \frac{1}{3n^3}$$

and

$$\begin{aligned} \sum_{n+1}^{\infty} (\int \Phi(u) \sin jx dx)^2 &\leq \sum_1^{\infty} (\int \Phi(u) \sin jx dx)^2 \leq \frac{\pi}{2} \int \Phi^2(u) dx. \end{aligned}$$

By (30), $\int \Phi^2(u) dx \leq E(t)$, and so from (36) we can conclude

$$|\int \{L(u) - f\} G(x, \xi) dx| \leq \frac{2}{\pi} \sqrt{\frac{1}{3n^3}} \cdot \sqrt{\frac{\pi}{2} E(t)}.$$

Referring back to (35) we have a theorem:

Theorem 7.

$$| -u_n(\xi, t) - \int \Phi(u_n) G(x, \xi) dx | \leq \sqrt{\frac{2E(t)}{3\pi n^3}}. \quad (37)$$

5 Existence Theorem

Recall now the discussion of equicontinuity and the criterion of Tonelli given at the end of section 2. Because of the inequalities contained in theorems 1 through 6, we see that the functions u_n , u_{nx} , and u_{nt} are uniformly equicontinuous. Thus by Ascoli's theorem, we can pick out, by a suitable diagonal process, a subsequence $\{v_m\}$ of $\{u_n\}$ such that

$$\begin{aligned} v_m(x,t) &\rightarrow u(x,t), \\ v_{mx}(x,t) &\rightarrow v(x,t), \\ v_{mt}(x,t) &\rightarrow w(x,t), \end{aligned}$$

the approach being uniform in each case. Because of this uniformity, one sees that $v = \partial u / \partial x$ and $w = \partial u / \partial t$. The derivatives u_{xx} , u_{xt} , u_{tt} exist almost everywhere and are summable. We consider the expression

$$-v_m(\xi,t) - \int \Phi(v_m) G(x,\xi) dx. \quad (38)$$

Since $v_m = u_{n_m}$, theorem (7) tells us that the expression (38) tends to zero as $m \rightarrow \infty$. Also $\Phi(v_m) \rightarrow \Phi(u)$ uniformly, and G is bounded, so we have

$$\begin{aligned} \lim_{m \rightarrow \infty} \{ -v_m(\xi,t) - \int \Phi(v_m) G(x,\xi) d\xi \} \\ = -u(\xi,t) - \int \Phi(u) G(x,\xi) dx = 0. \end{aligned}$$

The integral $\int \Phi(u) G(x,\xi) dx$ has a second derivative with respect to x , namely, $-\Phi(u) = -u_t - gu - f$. Hence u_{xx} exists, is continuous, and satisfies

$$-u_{xx}(\xi,t) + \Phi(u) = 0, \quad (39)$$

which is identical with (1). Since u obviously satisfies (2), it is a solution of the system (1), (2) in the region S_T . But T is arbitrary, and a solution u of (1), (2) exists in all S .

We have shown that a subsequence of u_n converges to the solution u , together with first order derivatives. Actually the entire sequence u_n so converges. If this were not the case, there would be a subsequence u_{n_j} such that, for instance, $\lim_{n_j \rightarrow \infty} u_{n_j} \neq u$ at some point x,t . From the sequence u_{n_j} we could pick another subsequence in the manner of this section, such that this subsequence has for limit a solution v of the problem. By the uniqueness theorem, $u = v$, and so $\lim_{n_j} u_{n_j}(x,t) = u(x,t)$, a contradiction. Thus we have

Theorem 8. The sequence $u_n(x,t)$ has for uniform limit in any S_T the unique function $u(x,t)$ satisfying (1) and (2). Furthermore, $u_{nx} \rightarrow u_x$ and $u_{nt} \rightarrow u_t$ uniformly in S_T .

6. Convergence in Mean of u_{nxx}

The solution $u(x,t)$ of the system (1), (2) has a Fourier series

$$u(x,t) = \sum_1^{\infty} \lambda_j(t) \sin jx,$$

which may be differentiated twice with respect to x , and the resulting series converges in the mean to u_{xx} . Set

$$v_n = \sum_1^n \lambda_j(t) \sin jx. \quad (40)$$

If (4) are multiplied by $-j^2(C_{n,j}(t) - \lambda_j(t))$ and summed, we get

$$\int \{ L(u_n) - f \} \frac{\partial^2}{\partial x^2} (u_n - v_n) dx = 0. \quad (41)$$

Since $L(u) = f$, we may write (41) as

$$\int \{ L(u_n - v_n) - L(u - v_n) \} \frac{\partial^2}{\partial x^2} (u_n - v_n) dx = 0. \quad (42)$$

Recalling that

$$L(u) = u_{xx} - u_t - gu - f + f = u_{xx} - \Phi(u) + f,$$

we write (42) in the form

$$\begin{aligned} \int \left\{ \frac{\partial^2}{\partial x^2} (u_n - v_n) \right\}^2 dx \\ = \int \{ -\Phi(u_n - v_n) + f - L(u - v_n) \} \frac{\partial^2}{\partial x^2} (u_n - v_n) dx, \end{aligned}$$

and Schwarz's inequality gives

$$\begin{aligned} \sqrt{\int \left\{ \frac{\partial^2}{\partial x^2} (u_n - v_n) \right\}^2 dx} \\ \leq \sqrt{\int \{ -\Phi(u_n - v_n) + f - L(u - v_n) \}^2 dx}. \quad (43) \end{aligned}$$

Now, $-\Phi(u_n - v_n) + f = -(u_n - v_n)_t - g(u_n - v_n) \rightarrow 0$ uniformly, and $L(u - v_n)$ tends to zero in the mean; thus the right hand member of (43) tends to zero and so

$$\int \left\{ \frac{\partial^2}{\partial x^2} (u_n - v_n) \right\}^2 dx \rightarrow 0.$$

From this and an appropriate triangular inequality it follows that

$$\int \left\{ \frac{\partial^2}{\partial x^2} (u_n - u) \right\}^2 dx \rightarrow 0.$$

Theorem 9. The second derivatives u_{nxx} of the approximating functions tend in the mean to u_{xx} .

It is too much to expect that $u_{nxx} \rightarrow u_{xx}$ uniformly, or even pointwise. For example, $u_{nxx}(0,t) \equiv 0$, while $u_{xx}(0,t) = f(0,t)$, and this latter does not need to vanish identically.

7. Estimating the Error

From theorem 9 it is easy to prove that $L(u_n) \rightarrow L(u)$ in the mean. Set $L(u_n) = f_n$. Then

$$L(u_n - u) = L(u_n) - L(u) = f_n - f \xrightarrow{l.m.} 0.$$

Using the method of 4.2, we prove that

$$\int (u_n - u)^2 dx \leq t^2 F_n^{*2}(t),$$

where

$$F_n^*(t) = \max_{\tau \leq t} \sqrt{\int \{f(x, \tau) - f_n(x, \tau)\}^2 dx}. \quad (44)$$

Theorem 10. The mean square error satisfies the inequality

$$\int (u_n - u)^2 dx \leq t^2 F_n^{*2}(t), \quad (45)$$

where $F_n^*(t)$ is given by (44), and $F_n^*(t) \rightarrow 0$ as $n \rightarrow \infty$.

Theorem 10 unfortunately gives information only about the mean square of the error instead of about the error at a point, which would be desirable. This latter will be mentioned below. Apart from this, however, theorem 10 gives a rather curious sort of information about the m.s.e. It does not give an a priori estimate of this error; that is, we cannot say that after n steps the error will have a certain bound, because we do not know how rapidly F_n tends to zero. On the other hand, we do seem to know more than the mere fact that the m.s.e. tends to zero. We know that if we take n sufficiently large, the m.s.e. will be small and we shall know it to be small. What we do not know is when we shall know it!

In section 5 it was mentioned that u_{xt} is a summable function. Thus from $L(u) = f$, it follows that u_{xxt} is summable. From the equation $L(u_n - u) = f_n - f$, we obtain, using the methods of sections 4.3 and 4.4, the inequality analogous to (22),

$$\int \{ (u_n - u)_x \}^2 dx \leq t F_n^*(t) (A_n^*(t) + F_n^*(t)) = G_n^*(t),$$

where F_n^* and A_n^* are given by (13) and (20), except that $f_n - f$ is used instead of f in their definitions. Now we are unable to say much about the factor A_n^* since it contains terms of the form $\int \{f_{ni} - f_i\}^2 dx$; that is, depends on u_{xxt} and u_{nxx} . However, it appears quite likely in many cases that A_n^* will remain bounded as $n \rightarrow \infty$. In any case the inequality

$$\int \{ (u_n - u)_x \}^2 dx \leq G_n^*(t)$$

is a valid one, and in a given computation $G_n^*(t)$ can actually be computed. Setting $\Delta_n(x, t) = u_n - u$, we have from Schwarz's inequality

$$|\Delta_n(x, t)| = \int_0^x \Delta_{nx}(s, t) ds \leq \sqrt{x \int_0^x \{ \Delta_{nx}(s, t) \}^2 ds} \quad (46)$$

$$\leq \sqrt{x \int \{ \Delta_{nx}(s, t) \}^2 ds} \leq \sqrt{x G_n(t)}.$$

In an analogous fashion we obtain by integrating from π instead of from zero,

$$\Delta_n(x, t) \leq \sqrt{(\pi - x) G_n^*(t)}. \quad (47)$$

From (46) and (47) it follows that

$$\Delta_n(x, t) \leq \sqrt{\frac{2}{\pi} x(\pi - x) G_n^*(t)}.$$

Theorem 11.

$$|u_n(x, t) - u(x, t)| \leq \sqrt{\frac{2}{\pi} x(\pi - x) G_n^*(t)}, \quad (48)$$

where

$$G_n^*(t) = t F_n^*(t) \{ A_n^*(t) + F_n^*(t) \},$$

$$F_n^*(t) = \max_{\tau \leq t} \sqrt{\int \{f_n - f\}^2 dx},$$

$$A_n^*(t) = F_n^*(0) - t H_n^*(t),$$

$$\{ H_n^*(t) \}^2 = G_1(t) t^2 \{ F_n^*(t) \}^2 + 2 G_1(t) t F_n^*(t) + F_{1n}^*(t) + \{ F_{1n}^*(t) \}^2,$$

$$F_{1n}^*(t) = \max_{\tau \leq t} \sqrt{\int \{f_{ni}(x, \tau) - f_i(x, \tau)\}^2 dx}.$$

For any stage n of the solution, the function $G_n^*(t)$ can be explicitly calculated. If the A_n^* do remain bounded, then $G_n^* \rightarrow 0$, and we will get a useful estimate of the error.

LOS ANGELES, April 24, 1952.